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# RESPONSE OF A TURBULENT BOUNDARY LAYER TO LATERAL ROUGHNESS DISCONTINUITIES

by

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#### ABSTRACT

# RESPONSE OF A TURBULENT BOUNDARY LAYER TO LATERAL ROUGHNESS DISCONTINUITIES

The structure of the turbulent boundary layer on a flat plate consisting of alternate, longitudinal strips of smooth and rough (Sand paper, Grit 4) surfaces parallel to the direction of the flow has been investigated. Measurements of mean velocity, wall shear stress, turbulent intensities, turbulent shear stresses and energy spectra of the streamwise turbulent velocity were obtained.

The flow field can be subdivided into three regions. These are: (a) the "smooth" region along the centerline of each smooth strip; (b) the "rough" region along the centerline of each rough strip; and (c) the "intermediate" region lying between the smooth and the rough regions. In both smooth and rough regions, the flow conditions are found to be nearly analogous to those in two-dimensional boundary layers over smooth and rough walls, respectively. In the intermediate region the wall shear stress was found to adjust itself very rapidly to the local conditions of the wall while going from the smooth to the rough region. The vertical distribution of mean velocity can satisfactorily be expressed by the two-dimensional descriptions for the law of the wall and the velocity-defect law. However, the two laws are not universal but form a family of curves depending on the local wall shear stress. The spanwise change from smooth-wall conditions to rough-wall conditions will induce a weak cross-flow directed from the rough to the smooth region near the wall. A theoretical description of the turbulent

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shear stress  $-\rho v w$  which is responsible for the generation of this crossflow has been presented. The spanwise adjustment of turbulent quantities to local conditions of the wall takes place at a slower rate than does the wall shear stress. The normalized spectra of the streamwise turbulent velocity remains largely unaffected by the surface configuration investigated here.

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## LIST OF SYMBOLS

Symbol	Definition	Dimension
А, А'	Constants	
В, В'	Constants	
Bw	Filter bandwidth	
С, с	Constants	
$^{C}\mathbf{f}$	Local skin friction coefficient	
D	Constant	
d	Height of rectangular-parallelepiped obstacle	L
d <sub>1</sub>	Diameter of Preston tube	L
E	Mean voltage	V
E(k)	One-dimensional wave-number spectrum funtion	$L^3/T^2$
E(n)	One-dimensional frequency spectrum function	Т
е	Voltage fluctuation	V
F	Function	
f	Funtion	
$\mathbf{f}_{\mathrm{p}}$	Function	
fr	Function	
G	Function	
g, g <sub>1</sub>	Functions	
Н	Shape factor	
∆h	Difference between dynamic and static pressure head	M/LT <sup>2</sup>
k	One-dimensional wave number	$L^{-1}$
<sup>k</sup> r	Roughness height	L
m	Exponent	
n	Frequency	. T <sup>-1</sup>

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### LIST OF SYMBOLS - Continued

Symbol	Definition	Dimension
Р	Mean static pressure	M/LT <sup>2</sup>
P <sub>∞</sub>	Static pressure outside boundary layer	M/LT <sup>2</sup>
$\Delta p$	Preston-tube reading	M/LT <sup>2</sup>
Red	Reynolds number based on d	
Rek	Reynolds number based on k	
s <sub>1</sub> , s <sub>2</sub>	Sensitivities of crossed wire at $\alpha = \pm 45^{\circ}$	VT/L
s <sub>U</sub>	Sensitivity of voltage with respect to velocity	VT/L
s <sub>α</sub>	Sensitivity of voltage with respect to angle	VT/L
U	Streamwise mean-flow velocity	L/T
U*	Shear velocity	L/T
U <sub>*s</sub>	Shear velocity corresponds to smooth-wall flow	L/T
U*r	Shear velocity corresponds to rough-wall flow	L/T
U <sub>∞</sub>	Free stream velocity	L/T
u,v,w	Velocity <b>f</b> luctuations in x,y,z direction	L/T
V,W	Secondary-flow components in y,z direction	L/T
x,y,z	Streamwise, spanwise, vertical coordinate	L
<sup>z</sup> o	Roughness parameter	L
α	Angle of attack	
δ	Boundary layer thickness	L
δ*	Displacement thickness	L
η	Non-dimensional spanwise coordinate	
ξ	Non-dimensional vertical coordinate	
θ	Momentum thickness	L
κ	Von Kármán constant	

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# LIST OF SYMBOLS - Continued

Symbol	Definition	Dimension
λ	Strip width	L
ν	Kinematic viscosity	$L^2/T$
П	Coles' pressure gradient parameter	
π	3.14159	
σ	Parameter	
τ	Non-dimensional streamwise coordinate	
τ <sub>w</sub>	Wall shear stress	ML/T <sup>2</sup>
φ	Non-dimensional mean velocity	
ω	Wake function of Coles	

examples of configurations which can produce cross flows. Near a roughness discontinuity the modification of mean velocity and turbulence profiles can be quite appreciable. At great lateral distances from the discontinuity, however, the chracteristics of the flow can be expected to be approximately the same as that over uniformly rough surfaces. A review of literature reveals that very little attention has been devoted to the effect of lateral roughness changes on the development of the atmospheric boundary layer. Neither theoretical treatments nor experimental data are available at the present stage.

The complex nature of atmospheric flow has stimulated efforts to study its characteristics in the laboratory under controlled conditions. In the lower layer of the atmosphere, the characteristics of the flow are found to be nearly analogous to that in the turbulent boundary layer along a flat plate. The feasibility of laboratory simulation of atmospheric flow for a wide range of situations has been established (Cermak, et al., 1966, Hidy, 1966, and Nemoto, 1968). The simulation of atmospheric flow over non-uniform surface conditions poses no new problems provided appropriate scaling relationships are available.

### 1.1 Purpose of Present Investigation

The objective of the present experimental investigation is to provide qualitative and quantitative information necessary for an analysis of the response of the turbulent boundary layer to lateral changes in surface roughness. In particular, information is sought regarding the boundary layer configuration and structure of the flow

over longitudinal strips of alternating surface roughness. The main purposes of this investigation are: (a) to study the effect of lateral roughness changes on mean velocity distribution, wall shear stress, turbulent intensities, turbulent shear stresses and energy spectra of the streamwise turbulent velocity by means of extensive pitot-tube and hot-wire surveys, and (b) to compare the experimental data with available data from two-dimensional flows over uniform surfaces. The experimental results will permit a better understanding of the flow characteristics, will provide a physical basis for developing an analytical model, and will have direct application to the study of atmospheric flows with complex boundary conditions which lead to three-dimensional mean motions.

### 1.2 Scope of Present Investigation

The present investigation deals with a turbulent boundary layer with zero pressure gradient over a flat plate consisting of alternate, longitudinal strips of smooth and rough surfaces, as shown in Fig. 1.1. The strips which are of equal width are placed in the streamwise direction starting from the leading edge of the plate. The rough strips are uniformly covered with sand grains. The experimental investigation is limited to one type of surface configuration with the free-stream velocity as the only variable.

A detailed survey of literature concerning turbulent boundary layers over both smooth and rough walls and research in the area of boundary layer flows over streamwise roughness changes is given in Chapter II.

The theoretical description of the flow is presented in Chapter III. The present study treats the non-uniform three dimensional surface conditions as a small perturbation about a twodimensional state. The effect of spanwise roughness changes on the mean velocity distribution is discussed. The distribution of the turbulent shear stress  $-p\overline{vw}$  which is responsible for the generation of a cross-flow in the boundary layer has been estimated.

The experimental equipment and procedure is described in detail in Chapter IV. Mean velocity profiles through the boundary layer at selected stations were made with a pitot tube and crosschecked by a single hot-wire. Measurement of the wall shear stress along a smooth strip of the flat plate in the boundary layer was obtained using a Preston tube. The intensities of turbulence and the turbulent shear stresses were measured by means of a single hot-wire and a crossed hot-wire. The spectra of the streamwise turbulent velocity were analyzed from the turbulence data.

The experimental results, given in Chapter V, are concentrated on the spanwise variation of mean velocity and turbulence profiles at distances far downstream from the leading edge of the flat plate where the local boundary layer thickness is comparable to the strip width, but is very large compared to roughness height. A summary of these results and discussions is given in Chapter VI.

### 1.3 Effect of Strip Width on Boundary Layer Development

The entire flow field produced by a surface configuration of the present type is very complicated. Near the leading edge of the plate the roughness elements contained in the rough strips will cause

large disturbances in the oncoming flow while the boundary layer forms downstream from these disturbances. Such disturbances associated with strong mixing and momentum diffusion will cause the flow to become turbulent at a shorter downstream distance from the leading edge than occurs naturally. The disturbances fade away gradually as the flow progresses downstream. Nevertheless, continuous production of small localized disturbances persists in a small region along side edges of the rough strips. Additionally, the spanwise inequality of the turbulent normal stresses in the boundary layer, due to the presence of roughness changes, will induce a weak secondary flow in the form of longitudinal vortices (Fig. 1.2). This secondary flow will cause a redistribution of momentum in planes perpendicular to the direction of the mean flow, which results in spanwise distortions of mean velocity and turbulence profiles.

The influence of strip width on the boundary layer characteristics can be determined by introducing the relative width  $\lambda/\delta$ , defined as the ratio of the strip width  $\lambda$  to the local boundary layer thickness  $\delta$ . Near the leading edge of the plate where  $\frac{\lambda}{\delta} >> 1$ , the flow field is unaffected by the presence of roughness changes except in the proximity of the wall near the roughness discontinuities. For  $\frac{\lambda}{\delta} \sim 1$ , that is, when the local boundary layer thickness is comparable to the strip width, the flow field will be influenced by the presence of roughness changes. The strip width then becomes a significant length scale and must be introduced as an additional parameter to characteristics of the flow. For  $\frac{\lambda}{\delta} << 1$ , that is, when the local boundary layer that is, when the local boundary layer that is, when the local boundary be introduced as an additional parameter to characteristics of the flow. For  $\frac{\lambda}{\delta} << 1$ , that is, when the local boundary layer that the strip width, the flow is again almost everywhere independent of the width.

#### 1.4 Effect of Surface Roughness on Boundary Layer Development

In turbulent flow, the proximity of the rigid boundary or wall has a direct bearing on the turbulence. For a smooth wall this effect occurs through the action of viscous stresses, for a rough wall through the action of more complicated forces resulting from the flow around the individual roughness elements. The action of the roughness can be interpreted as being equivalent to a reduction of the viscous sublayer. The roughness effect on boundary layer characteristics depends on the ratio of the roughness dimension to the thickness of the viscous sublayer. If the roughness elements are much smaller than the thickness of the viscous sublayer, the surface is considered as hydraulically smooth. For the hydraulically smooth surface, no effect of the roughness is felt. On the other hand, if the roughness elements are so large as to protrude through the viscous sublayer completely, the surface is considered to be completely rough. For the completely rough surface, the turbulent motion produced by the flow around the roughness elements dominates the wall shear stress and the viscous action has comparatively little influence. Between these two extremes there exists a transitional state where only a fraction of the roughness elements disturbs the viscous sublayer. Consequently, the wall shear stress in this case will be affected by both the viscosity and the roughness dimension.

For a surface configuration of the present type, the roughness height representing the scale length of the roughness on each of the rough strips remains constant, while the boundary layer thickness increases downstream. As far as the flow along the centerline of a rough strip in concerned, this circumstance causes the flow

characteristics near the leading edge of the flat plate to behave differently from that at very large distances downstream. The completely rough regime occurs over the forward portion, followed by the transition regime and eventually, the rough strip may become hydraulically smooth if the plate is sufficiently long.

#### CHAPTER II

#### LITERATURE REVIEW

A survey of the literature reveals that very little progress has been made concerning the structure of the turbulent boundary layer produced by the present surface configuration. Previous investigations have been primarily concerned with the effect of spanwise non-uniformity of surface roughness on turbulence suppression phenomena and on secondary flow phenomena. Concerning turbulence suppression, Liu, Kline and Johnston (1966) investigated the effect of extremely narrow roughness strips having various roughness spacings on a smooth surface using flow visualization techniques. Their results show that neither sufficient suppression of turbulent production nor significant reduction in wall shear stress have been found. The effect of a spanwise change from smooth-wall conditions to rough-wall conditions on secondary flow phenomena were studied by Hinze (1967). He concluded that the lateral change in surface roughness will cause a similar local secondary flow as the side edge of a finite flat plate (Elder, 1960).

In working with the turbulent boundary layer consisting of flow parallel to strips of smooth and rough surfaces, it is advantageous to do a parallel study of two-dimensional turbulent boundary layers over both smooth and rough walls. In this chapter, previous investigations of turbulent boundary layers over both smooth and rough walls and research in the area of boundary layer flows over streamwise roughness changes are reviewed.

#### 2.1 Turbulent Flow Over Smooth Walls

Close examination of a two-dimensional turbulent boundary layer over smooth walls reveals a characteristic which allows division of the boundary layer into (i) an inner layer, and (ii) an outer layer. The inner layer, whose thickness is roughly 10 to 20 percent of the boundary layer thickness derives its supply of turbulent energy almost exclusively from the interaction of the turbulent shear stress and the mean velocity gradient within itself. Within the layer the flow is dependent only on local wall conditions and the shear stress is nearly constant. This layer is also referred to as the constantstress layer. The mean velocity distribution across this layer is determined by the wall shear stress  $\tau_w$ , the density  $\rho$ , the kinematic viscosity  $\nu$ , and the vertical distance z from the wall. Application of dimensional analysis to these variables on which the mean velocity U depends gives the functional relationship

$$\frac{U}{U_{\star}} = f\left(\frac{zU_{\star}}{v}\right)$$
(2-1)

where  $U_* = \sqrt{\tau_w/\rho}$  is the shear velocity.

Equation (2-1) is known as Prandtl's law of the wall. This law implies that mean velocity profiles are independent of position along the plate when the mean velocity is measured in terms of its scale  $U_*$ , and the distance from the wall is measured in terms of its scale  $\vee/U_*$ .

In the case of a smooth wall there exists a sublayer adjacent to the wall where the flow is predominantly viscous and the turbulent shear stress is negligible. This is called the viscous sublayer,

whose thickness is of the order of 0.1 to 1.0 percent of the boundary layer thickness. The mean velocity distribution in this sublayer follows the linear relationship

$$\frac{U}{U_{\star}} = \frac{zU_{\star}}{v} \quad . \tag{2-2}$$

The outer layer, which contains 80 to 90 percent of the boundary layer, receives its greatest portion of turbulent energy by diffusion from the constant-stress layer. Within this layer the flow is completely turbulent and the effect of viscosity is negligible. It is practically independent of local wall conditions but depends instead on external conditions and on the upstream history of the flow. The velocity defect  $U_{\infty}$  - U is thus a function of the vertical distance z, the boundary layer thickness  $\delta$ , and the shear velocity  $U_{\star}$ . Dimensional analysis gives

$$\frac{U_{\infty} - U}{U_{\star}} = g(\frac{z}{\delta})$$
(2-3)

This similarity relation is known as Kärmän's velocity-defect law. This law implies that the asymptotic approach of the mean velocity profile to the free-stream velocity  $U_{\infty}$  is uniquely determined by the distance from the wall relative to the boundary layer thickness.

Since the division of the boundary-layer into inner and outer layers is arbitrary, there must be a region in which both Eqs. (2-1) and (2-3) are valid. This overlapping of the two similarity laws indicates that some relationship must exist between two distinct set of parameters used. To satisfy this conditions, the functions f and g must be logarithmic. Therefore, for the overlap region the mean velocity distribution can be expressed as

$$\frac{U}{U_{\star}} = \frac{1}{\kappa} \, \ln \frac{zU_{\star}}{v} + A \tag{2-4}$$

and

$$\frac{U-U_{\infty}}{U_{\star}} = \frac{1}{\kappa} \ln \frac{z}{\delta} + B$$
(2-5)

where  $\kappa$ , A and B are experimentally determined constants. In particular,  $\kappa$  is known as von Kármán's constant. The value of  $\kappa$  is usually taken as 0.41 for flat plate flows. A summary of the experimental determinations of  $\kappa$  has been made by Slotta (1963).

Telles and Dukler (1968) demonstrated that the law of the wall (Eq. 2-4) is the zeroth-order small perturbation type of solution of the boundary layer equations, but that the velocity-defect law (Eq. 2-5) does not represent a valid similarity type of solution of the boundary layer equations. The law of the wall has been generalized empirically to flows with arbitrary pressure gradients by Ludwieg and Tillmann(1950), and the velocity-defect law to a certain class of equilibrium flows by Clauser (1954).

Usually a single function is sought to describe the mean velocity distribution for the entire boundary layer. One of the earliest empirical descriptions is the power law profiles. The general form of the power law (Schlichting, 1968) is given by

$$\frac{U}{U_{\infty}} = \left(\frac{z}{\delta}\right)^{\frac{1}{m}}$$
(2-6)

where m is an empirical exponent.

Coles (1956) deduced a purely empirical correction function to account for the large observed deviations between the measured mean velocity profiles and the logarithmic form of the "law of the wall" in the outer region of the boundary layer,

$$\frac{U}{U_{\star}} = \frac{1}{\kappa} \ln \frac{zU_{\star}}{\nu} + A + \frac{\Pi}{\kappa} \omega \left(\frac{z}{\delta}\right)$$
(2-7)

where  $\Pi$  is a parameter which depends on pressure gradient and turbulence level near the edge of the boundary layer, and  $\omega(\frac{z}{\delta})$  is the wake function known as "the law of the wake". Coles originally gave the wake function in numerical form. Hinze (1959) suggested that the wake function can be accurately described by a closed form function

$$\omega\left(\frac{Z}{\delta}\right) = 1 - \cos\left(\frac{\pi Z}{\delta}\right) \tag{2-8}$$

Telles and Dukler (1968) demonstrated that the first-order small perturbation type solution of the boundary layer equations will generate a correction function similar to Coles' law of the wake.

Minor inconsistencies in Coles' formulation of the wake law were corrected by Bull (1968) who extend Eq. (2-8) by means of a simple power law. The resulting correction function was rather cumbersome and involved two parameters which had to be independently evaluated from experimental data. A similar approach was employed by Allan (1970) who deduced a generalized correction function for pipe flows, flat plate flows and diffusing flows.

There now exists a considerable amount of experimental work on smooth-wall turbulent boundary layers. The experimental results have

been accumulated in the past by Baines (1951), Ross (1953), Hama (1954), Clauser (1956), Townsend (1956), Hinze (1959) and Rotta (1962). Methods of computation of the turbulent boundary layers have been collected by Kline, et al. (1968).

#### 2.2 Turbulent Flow over Rough Walls

For a two-dimensional turbulent boundary layer over a rough wall, whose surface is uniformly covered by geometrically similar roughness of size small compared with the boundary layer thickness, there exists a constant-stress layer adjacent to the wall, where the motion is entirely determined by the wall shear stress  $\tau_w$ , density  $\rho$ , kinematic viscosity v and roughness height  $k_r$  (Rotta, 1962). The application of the similarity relation to the mean velocity distribution leads to

$$\frac{U}{U_{\star}} = F \left(\frac{zU_{\star}}{v}, \frac{k_{\rm r}U_{\star}}{v}\right)$$
(2-9)

Experimental studies by Moore (1950) indicated that the velocitydefect law as expressed in Eq. (2-5) is universal for either smooth or rough walls. Then Eq. (2-9) can thus be expressed for a rough wall by the relation

$$\frac{U}{U_{\star}} = \frac{1}{\kappa} \ln \frac{zU_{\star}}{v} + C \left(\frac{krU_{\star}}{v}\right)$$
(2-10)

where  $C(\frac{\mathbf{r} \cdot \mathbf{v}}{v})$  is an empirical constant depending upon the roughness geometry.

By virtue of Moore's conclusion, Clauser (1956) suggested that the mean velocity distribution in the constant-stress layer can be expressed generally as

$$\frac{U}{U_{\star}} = \frac{1}{\kappa} \ln \frac{zU_{\star}}{v} + A - \frac{\Delta U}{U_{\star}}$$
(2-11)

where  $\Delta U/U_*$  represents the downward shift of mean velocity profiles in the vicinity of the roughness elements. This shift is a function of the roughness Reynolds number  $\operatorname{Re}_k = k_r U_*/v$ . For completely rough flows Clauser (1956) showed that  $\Delta U/U_*$  takes the form

$$\frac{\Delta U}{U_{\star}} = \frac{1}{\kappa} \ln \frac{k_{\rm r} U_{\star}}{v} + D$$
(2-12)

where D is an empirical constant. However, Bettermann (see Dvorak, 1969) suggested that  $\Delta U/U_*$  must also depend upon the roughness spacing.

The "law of the wall" (Eq. 2-10) may be rewritten as

$$\frac{U}{U_{\star}} = \frac{1}{\kappa} \ln \frac{z}{z_{o}}$$
(2-13)

where  $z_0$  is a roughness parameter which is independent of the flow conditions and must be given as part of the boundary conditions. Summaries of values of  $z_0$  for various surface roughness are presented by Priestley (1959). In general, the value of  $z_0$  is an order of magnitude less than the average height of surface roughness elements.

The "power law" (Eq. 2-6) has been used to describe the mean velocity distribution for the turbulent boundary layer over rough walls by many investigators (e.g., Corrsin and Kistler, 1954; Chowdhury, 1966; and Liu, Kline and Johnston, 1966). The "law of the wake" has been modified to include the effect of surface roughness by Allan (1970). For rough-wall flows, the effective origin of z is undefined due to the presence of roughness elements. Intuitively, the effective origin can be located anywhere between the trough and crest of the roughness elements. Scottron (1967) argued that the location of the effective origin may be presumed to be associated with the generation of turbulence near the wall and depend on roughness geometry and local roughness Reynolds number  $\text{Re}_k$ . Based on the assumption of the existence of a logarithmic velocity profile near the wall, Perry and Joubert (1963) proposed a method for determining the effective origin of the wall. Their scheme consisted essentially of adjustment of the z origin until a best-fitting straight line could be drawn from a  $U/U_{\infty}$  versus z plot on a semi-logarithmic paper.

Experimental investigations on rough-wall turbulent boundary layers are those of Tillmann (1945), Baines (1950), Moore (1950), Hama (1954), Corrsin and Kistler (1954), Chanda (1958), Liu, Kline and Johnston (1966), Chowdhury (1966) and Perry, Schofield and Joubert (1969) in zero pressure gradient, and those of Perry and Joubert (1963) and Scottron (1967) in adverse pressure gradient. The experimental results of all these investigators indicate the existence of the logarithmic mean velocity profile in the constantstress layer.

### 2.3 Boundary Layer Flow over Streamwise Roughness Changes

Significant progress towards an understanding of turbulent boundary layer flow above an abrupt change in surface roughness under neutral stratification has been made in recent years. Al existing studies, which are mainly concerned with the lower layer of the

atmosphere, treat the problem in two dimensions, with the line of roughness discontinuity normal to the mean flow direction. The roughness heights on either side of the discontinuity are assumed to correspond to homogeneous terrain of infinite extent. Upstream of the roughness discontinuity the flow is assumed to be in equilibrium with the underlying surface roughness and to have a constant shear stress distribution independent of height and a logarithmic mean velocity profile. Downstream of the discontinuity the flow will gradually adjust itself at all levels to the downstream roughness. The effect of the downstream roughness is assumed to be confined to a so-called "internal boundary layer," which increases with downstream distance from the roughness discontinuity.

Several theories have been proposed for predicting the growth of the internal boundary layer by prescribing mean velocity and shear stress profiles inside the layer. Elliott (1958) assumed a constant shear stress distribution and a logarithmic velocity profile. This led to a discontinuity of shear stress at the edge of the internal boundary layer. Panofsky and Townsend(1964) removed this physically unrealistic discontinuity at the interface by assuming a linear variation of the shear stress with height and a log-linear velocity profile. Townsend (1965a, 1965b, 1966) introduced a self-preserving development of the flow modification induced by the downstream roughness change. The deviations of mean velocity from the upstresm distribution were expressed as the result of a flow acceleration and a vertical displacement of the streamline. Plate and Hidy (1967) added advection to Elliott's technique and incorporated Townsend's theory in a modified form, such that the result could be extended

to include a pressure gradient and a non-uniform shear stress at the surface. Blom and Wartena (1969) corrected a minor discrepancy between Townsend's resulting profile and his inner boundary condition, and further extended Townsend's theory to more than one discontinuity in surface roughness.

A different theoretical approach to the internal boundary layer problem was employed by Miyake (1961). He did not explicitly develop a theory for the velocity profile, but derived an expression for the height of the internal boundary layer. Miyake's theory is based on the simple assumption that the slope of the internal boundary layer is proportional to the ratio of the standard deviation of vertical velocity to the upstream velocity.

In conclusion, all theories agree in that an internal boundary layer develops after the roughness change, and the height of which varies approximately with the downstream distance from the roughness discontinuity, raised to the 4/5 power. The theories disagree with each other, however, in the shear stress distribution within the internal boundary layer, further, they differ in the downstream distance required to set up a new equilibrium up to a given height.

For flow from a smooth to a rough surface, the rough surface behaves in an "active" manner. The growth of the internal boundary layer may be dominated by the "wake" from the leading roughness elements. Because of the vigorous generation of turbulence close to the rough surface, adjustment of the wall shear stress to the downstream roughness is very rapid. For flow from a rough to a smooth surface, the smooth surface behaves in a "passive" manner. The growth of the internal boundary layer is inhabited in its growth by

diffusion of information from the "active" outer flow back towards the surface. This leads to a slower adjustment of the wall shear stress.

Wind-tunnel studies of the boundary layer flow above an abrupt change in surface roughness were performed by Jacobs (1940), Clauser (1956), Logan and Jones (1963), Plate and Hidy (1967), Yeh (1970), and Antonia and Luxton (1971a,1971b) following a smooth-to-rough change, and by Jacobs (1940) and Taylor (1962) following a rough-to-smooth change. Jacobs, Taylor, Plate and Hidy, and Antonia and Luxton conducted the experiments in channel flows, Logan and Jones in pipe flows, and Clauser and Yeh in flat plate flows. Atmospheric observations were made by Blackadar et al. (1967), Bradley (1968), and Blackadar and Panofsky (1970). The results of all these investigators indicated that the adjustment of the wall shear stress to downstream roughness occurs at a short distance downstream of the roughness discontinuity. A decelerated flow moving from smooth-to-rough surface and an accelerated flow moving from rough-to-smooth surface were observed.

Numerical studies of the internal boundary layer problem were treated by Wagner (1966), Onishi (1966), Nickerson (1968), Estoque and Bhumralkar (1968), Peterson (1969), Taylor (1969, 1970) and Huang (1971). These studies differ from each other mainly in the use of the governing equations and of the boundary conditions. Wagner's model is time-dependent, unlike the other models in which stationary is assumed. Nickerson considers a step change in wall shear stress, Taylor a step change in surface roughness parameter, and Peterson the shear stress proportional to the turbulent energy. Except Peterson's model, all other models depend on the assumed behavior of the mixing length theory or the momentum exchange coefficient.

# CHAPTER III THEORETICAL DESCRIPTION OF THE FLOW

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Since turbulent boundary layers involve the complex combination of boundary layer phenomena, non-linear mechanics and turbulence, a direct theoretical solution of the problem throughout the entire region occupied by the boundary-layer flow has not yet been attained. In order to obtain insight into the turbulent boundary layers, one is evitably led to strong dependence on experimental and empirical investigations.

In this chapter, a theoretical description of the flow is presented, and the turbulent shear stress  $-\rho \overline{vw}$ , which gives rise to a weak cross-flow within the boundary layer, has been estimated.

#### 3.1 Basic Equations

The present analysis considers a steady, incompressible turbulent boundary layer over alternate, longitudinal strips of smooth and rough surfaces with zero streamwise pressure gradient. In order to simplify the analysis and yet still retain the essential physics of the problem, the following assumptions are made:

- the roughness height is so small compared with the local boundary layer thickness that the mean flow and pressure field outside the boundary layer are not affected by the roughness changes,
- the disturbances generated along the side edges of each of the rough strips can be neglected,

- the rates of variation of quantities in the streamwise direction are small with respect to changes of the same quantities in the spanwise and vertical directions,
- the turbulent fluctuations are small with respect to the mean flow, and
- 5) the magnitude of the components of the secondary flow is small compared with that of the mean flow.

The Reynolds equation of motion then reduces to

$$U \frac{\partial U}{\partial x} + V \frac{\partial U}{\partial y} + W \frac{\partial U}{\partial z} + \frac{\partial \overline{uv}}{\partial y} + \frac{\partial \overline{uw}}{\partial z} = v \frac{\partial^2 U}{\partial z^2}$$
(3-1)

$$\frac{\partial v^2}{\partial y} + \frac{\partial \overline{vw}}{\partial z} = -\frac{1}{\rho} \frac{\partial P}{\partial y}$$
(3-2)

$$\frac{\partial \overline{vw}}{\partial y} + \frac{\partial w^2}{\partial z} = -\frac{1}{\rho} \frac{\partial P}{\partial z}$$
(3-3)

where x, y, z are the streamwise, spanwise and vertical coordinates of the cartesian coordinate system, as shown in Fig. 1.2. U is the streamwise mean-flow velocity, and V, W are defined as the secondaryflow components in y- and z-directions, respectively. u, v and w are the corresponding turbulent velocity fluctuations in x, y, z directions. P is mean static pressure. The terms  $\rho \overline{v^2}$ ,  $\rho \overline{w^2}$ ,  $-\rho \overline{uv}$ ,  $-\rho \overline{uw}$  and  $-\rho \overline{vw}$ are known as the turbulent normal and shear stresses, respectively, also known as Reynolds stresses. These stresses are the representation of the rate of momentum transport because of the turbulent velocity fluctuations.

The continuity equation for the mean motion is

$$\frac{\partial U}{\partial x} + \frac{\partial V}{\partial y} + \frac{\partial W}{\partial z} = 0$$
 (3-4)

Eqs. (3-1) through (3-4) with the appropriate boundary conditions do not form a complete set of equations from which a solution could be obtained. The closure of the system requires that all statistical quantities of turbulence be specified. So far, no satisfactory solution to this problem has been obtained.

#### 3.2 Mean Velocity Distribution

Cumpsty (1970) in his review of three-dimensional boundary layers made the following statement:

"When the three-dimensional effects, including cross-flow, are small, the use of two-dimensional descriptions for the flow is nature and needs no justification."

The turbulent boundary layer produced by surface configuration of the present type falls into this category.

The present analysis of the three-dimensional turbulent boundary layer presumes the existence of three different regions, the "smooth" region along the centerline of each smooth strip, the "rough" region along the centerline of each rough strip and the "intermediate" region lying between the smooth and the rough regions. The spanwise variation of the boundary layer thickness is assumed to be negligibly small.

In the smooth region the flow is assumed to behave like a twodimensional boundary layer over a smooth surface. The mean velocity distribution in this region may be described by the law of the wall

$$\frac{U}{U_{\star s}} = \frac{1}{\kappa} \ln \frac{z U_{\star s}}{v} + A$$
(3-5)

and the velocity-defect law

$$\frac{U - U_{\infty}}{U_{\star S}} = \frac{1}{\kappa} \ln \frac{z}{\delta} + B$$
(3-6)

where the subscript "s" refers to the smooth wall conditions. Eqs. (3-5) and (3-6) may be combined to give

$$\frac{U_{\infty}}{U_{*S}} = \frac{1}{\kappa} \ln \frac{\delta U_{*S}}{\nu} + (A - B)$$
(3-7)

In the rough region the flow is assumed to behave like a twodimensional boundary layer over a rough surface of uniform but identical roughness. The law of the wall for the rough-wall flows may be written, after Clauser (1956), as

$$\frac{U}{U_{\star r}} = \frac{1}{\kappa} \ln \frac{zU_{\star s}}{v} + A - \frac{\Delta U}{U_{\star r}}$$
(3-8)

where the subscript "r" refers to the rough wall conditions, and the term  $\Delta U/U_{*r}$  represents the vertical shift of the logarithmic profile caused by roughness. If the roughness elements are small compared with the boundary layer thickness, the universal velocity-defect law may be expressed as

$$\frac{U - U_{\infty}}{U_{\star r}} = \frac{1}{\kappa} \ln \frac{z}{\delta} + B$$
(3-9)

From the combination of Eqs. (3-8) and (3-9), it follows that

$$\frac{U_{\infty}}{U_{\star r}} = \frac{1}{\kappa} \ln \frac{\delta U_{\star s}}{\nu} + (A - B) - \frac{\Delta U}{U_{\star r}}$$
(3-10)

Thus, from Eqs. (3-7) and (3-10), there emerges a very significant relation between the smooth and rough regions, viz.,

$$\frac{U_{\infty}}{U_{\star s}} - \frac{U_{\infty}}{U_{\star r}} = \frac{\Delta U}{U_{\star r}}$$
(3-11)

In the intermediate region the roughness effect tends to increase from zero in the smooth region to a maximum value in the rough region. If the strip width  $\lambda$  is comparable to the boundary layer thickness, the relationship existing among the variables describing the roughness effect may be written, in general, as

$$\frac{\Delta U}{U_{\star r}} = \mathbf{f}_{\mathbf{r}} \left( \frac{\mathbf{k}_{\mathbf{r}} U_{\star \mathbf{r}}}{v}, \frac{\mathbf{y}}{\lambda} \right)$$
(3-12)

A knowledge of Eq. (3-12) thus permits the evalution of the complete correlation of geometry to roughness effect.

#### 3.3 Estimation of the Turbulent Shear Stress -pvw

Secondary flows in turbulent flow have been separated into two categories by Prandtl (1952). The secondary flow of the first kind is produced by mean flow skewing, while the secondary flow of the second kind is caused by the non-uniformity of anisotropic wall turbulence.

For the present study the secondary flow of Prandtl's second kind, produced by the inequality of turbulent normal stresses, predominates. Along the centerline of a smooth or rough strip,  $\overline{vw}$  vanishes by symmetry, so that integrating Eq. (3-3) with respect to z gives

$$\overline{w^2} = \frac{1}{\rho} (P_{\infty} - P)$$
 (3-13)

where  $P_{\infty}$  is the mean static pressure outside the boundary layer. Further, the integration of Eq. (3-2) with respect to y gives

$$\int_{0}^{\lambda} \frac{\partial \overline{vw}}{\partial z} dy = (\overline{v^{2}} + \frac{P}{\rho})_{y=0} - (\overline{v^{2}} + \frac{P}{\rho})_{y=\lambda}$$
(3-14)

If the flow field outside the boundary layer is assumed to be unaffected by the presence of the roughness changes, it follows from Eq. (3-13) that

$$\int_{0}^{\lambda} \frac{\partial \overline{vw}}{\partial z} dy = (\overline{v^{2}} - \overline{w^{2}})_{y=0} - (\overline{v^{2}} - \overline{w^{2}})_{y=\lambda}$$
(3-15)

For regions not too close to the wall, the distribution of  $(v^2 - w^2)$ in a normally two-dimensional boundary layer is strongly dependent on the vertical distance from the wall (Klebanoff, 1954). Non-dimensionalizing the turbulent normal stresses with  $U_*$ , the local shear velocity, it should be possible to write

$$\frac{\overline{v^2 - w^2}}{U_*^2} = G\left(\frac{z}{\delta}\right)$$
(3-16)

where G is a universal function. Perkins (1970) suggested that the behavior of the universal function can be adquately demonstrated by the approximation

G 
$$(\frac{z}{\delta}) = 1 - \frac{z}{\delta}$$
;  $\frac{z}{\delta} > 0.1$  (3-17)

with the aid of Eqs. (3-16) and (3-17), Eq. (3-15) becomes

$$\int_{0}^{\lambda} \frac{\partial \overline{vw}}{\partial z} dy = - \left( U_{\star r}^{2} - U_{\star s}^{2} \right) \left( 1 - \frac{z}{\delta} \right) \qquad ; \qquad \frac{z}{\delta} > 0.1 \qquad (3-18)$$

When the strip width is comparable to the local boundary layer thickness, it is reasonable to assume that  $\partial \overline{vw}/\partial z$  has its non-zero values distributed spanwisely between the centerlines of both smooth and rough strips. The distribution of  $\partial \overline{vw}/\partial z$  consistent with the symmetry of the flow about the centerlines, following Townsend (1956), is

$$\frac{\partial \overline{vw}}{\partial z} = \begin{cases} \left[\frac{\partial \overline{vw}}{\partial z}\right]_{y=\frac{\lambda}{2}} \frac{y}{\lambda} & ; & 0 \leq \frac{y}{\lambda} \leq \frac{1}{2} \\ \\ \left[\frac{\partial \overline{vw}}{\partial z}\right]_{y=\frac{\lambda}{2}} (1 - \frac{y}{\lambda}) & ; & \frac{1}{2} \leq \frac{y}{\lambda} \leq 1 \end{cases}$$
(3-19)

If this is introduced in Eq. (3-18) and by means of integration, we obtain

$$\left[ \frac{\partial \overline{vw}}{\partial z} \right]_{y=\frac{\lambda}{2}} = -\frac{4}{\lambda} \left( U_{*r}^2 - U_{*s}^2 \right) \left( 1 - \frac{z}{\delta} \right) ; \quad \frac{z}{\delta} > 0.1$$
 (3-20)

Then Eq. (3-19) upon integration gives

$$\overline{\mathbf{vw}} = \begin{cases} \frac{2\delta}{\lambda} \left( \mathbf{U}_{\mathbf{*r}}^2 - \mathbf{U}_{\mathbf{*s}}^2 \right) \frac{\mathbf{y}}{\lambda} \left( 1 - \frac{\mathbf{z}}{\delta} \right)^2 & ; \quad 0 \leq \frac{\mathbf{y}}{\lambda} \leq \frac{1}{2} \\ \\ \frac{2\delta}{\lambda} \left( \mathbf{U}_{\mathbf{*r}}^2 - \mathbf{U}_{\mathbf{*s}}^2 \right) \left( 1 - \frac{\mathbf{y}}{\lambda} \right) \left( 1 - \frac{\mathbf{z}}{\delta} \right)^2 & ; \quad \frac{1}{2} \leq \frac{\mathbf{y}}{\lambda} \leq 1 \end{cases}$$
(3-21)

This equation shows that the spanwise distribution of the turbulent shear stress  $-\rho \overline{vw}$  cannot be zero in this sort of flow, and the existence of this stress will tend to produce a cross-flow in the boundary layer.
### CHAPTER IV

### EXPERIMENTAL EQUIPMENT AND PROCEDURE

Experiments were performed in a low-speed wind tunnel located in the Fluid Dynamics and Diffusion Laboratory at Colorado State University. In this chapter, the experimental equipment, arrangement and procedure for this study are described.

### 4.1 Experimental Equipment

The basic experimental equipment consists of i) a wind tunnel, ii) a flat plate, iii) a carriage, iv) a manometer, v) two hot-wire anemometer systems, and vi) a wave analyzer. A brief description for each of the instruments is given below.

### 4.1.1 Wind Tunnel

The wind tunnel is a closed circulating type with a test section 30 ft. long, 6 ft. wide, and 6 ft. high. A schematic diagram of the wind tunnel is shown in Fig. 4.1. Air flow in the tunnel is driven by a constant pitch controlled fan. The tunnel can generate wind speeds ranging from 5 ft/sec to about 70 ft/sec. The turbulence level in the free stream is kept within 2 percent by providing damping screen and a 4:1 ratio entrance contraction. There is no arrangement for temperature control in the tunnel. A thermometer showed that the temperature varied from 70°F at the beginning of the test run to 90°F after a few hours of operation. It was thus necessary to warm up the tunnel for at least two hours to achieve a more or less steady temperature before starting the experiments. The temperature was recorded at the beginning and end of each experiment.

### 4.1.2 Flat Plate

The flat plate used for the experimental investigation was made of smoothed and polished aluminum plate. It is 12 ft. long, 4 1/2 ft. wide and 3/8 inch thick. The plate had a sharp leading edge which resulted in separation of the flow along that edge. The small region of separated flow was followed by reattachment in a transitional region; the flow became increasingly turbulent as it progressed further downstream. The plate was mounted on a steel stand at a height of 2 ft. above the wind tunnel floor in order to avoid the secondary motions resulting from the excessive loss of momentum in the corners of the tunnel, and was leveled accurately with a leveling instrument.

Four rough strips (sand paper, Grit 4), each of 12 ft. long, 1/2 ft. wide and 0.1 inch averaged roughness height, were attached on the plate longitudinally starting from the leading edge of the plate. The spacing between the two adjacent rough strips was kept at 1 foot. A total of fifty-four pressure tape with 1/32 inch in diameter holes, placed on either side of the centerline of the plate, was used to measure the wall static pressure along the plate. A plan view of the plate with pressure taps is shown in Fig. 4.2 and crosssectional elevation of the plate is shown in Fig. 4.3.

# 4.1.3 Carriage

Probes were mounted on a wind tunnel carriage located 3 ft. above the wind tunnel floor of the test section. The carriage moved along the wind tunnel on rails which were fastened to the vertical walls of the tunnel. Two small D.C. motors with driving mechanism were used to produce independent vertical and lateral movements. The

Vibration Analyzer and 1521 Graphic Level Recorder. The former consists essentially of a set of passive filters in the range of 2.5 to 25,000 HZ. The filters are of the octave type having 1/3-octave (23%) and 1/10-octave (7%) noise bandwidths. The latter can record automatically the spectrum of a region under analysis.

### 4.2 Experimental Arrangements

In order to obtain a condition of zero pressure gradient flow along the flat plate, a specially shaped false ceiling was mounted onto the wind tunnel ceiling, as shown in Fig. 4.4. The shape of the false ceiling is shown in Fig. 4.5. A dimensionless plot of static pressure distribution at the free stream along the centerline of the flat plate is shown in Fig. 4.6. The maximum magnitude of streamwise pressure gradients are about  $2.42 \times 10^{-6}$  psi/in. for  $U_{\infty}$ = 20 ft/sec and  $4.83 \times 10^{-6}$  psi/in. for  $U_{\infty}$ = 30 ft/sec, if the first two-foot section of the plate was excluded. These values are small enough to be considered as zero streamwise pressure gradient. A dimensionless plot of wall pressure distribution on the plate is shown in Fig. 4.7. The result indicates that the wall pressure distribution is of the same order of the corresponding free-stream pressure distribution. It also shows that the spanwise pressure gradient is very small.

For testing plane symmetric flow condition along the centerline of the flat plate, a comparison of similar mean velocity profiles extending to 1 ft on either side of the centerline was made for a streamwise distance of 5 ft. from the leading edge of the plate, as shown in Fig. 4.8. The results show no signs of three-dimensional

effects resulting from the whirling motion generated from the corner of the leading edge of the flat plate. It also indicates that a nearly plane symmetric flow about the centerline of the plate has been established.

#### 4.3 Experimental Procedure

The aspect of the experimental procedure consists of measurements of static pressure, mean velocity, wall shear stress, turbulence and frequency spectrum. The procedure of these measurements and the data reduction techniques are presented below.

### 4.3.1 Measurement of Static Pressure

The static pressure at the free stream along the flat plate was measured by using two pitot tubes with both static heads connected to an electronic pressure meter (MKS Baratron Type 77). Both tubes were first set at the same height, side by side in the free stream, near the leading edge and well above the surface of the plate. At this position, the pressure difference between these tubes is zero. One of the tubes was then used as a reference static pressure and the other tube was moved downstream along the plate. The streamwise pressure gradient, if any, was thus obtained.

The wall static pressure distribution on the surface of the plate were measured through pressure taps embedded in the plate. These pressure taps were connected successively through plastic tubing to the same pressure meter, where they were measured against a reference pressure of the tap located at x=6 ft. and y=1 ft. of the smooth surface of the plate. The accuracy of the static pressure measurement on the rough surface is uncertain due to large local static pressure variations.

### 4.3.2 Measurement of Mean Velocity

The mean velocity was measured with both a pitot tube and a single hot-wire. The hot-wire was also employed to cross-check the values obtained by pitot tube measurements.

#### (A) Pitot Tube for Mean Velocity

The pitot tube measurement of mean velocity was made with a 1/16 inch diameter standard pitot-static tube. The dynamic and static pressures were obtained from the dynamic and static taps of the tube, respectively. The two taps of the tube were connected by flexible tubings to the terminal of an electronic pressure meter (MKS Type 77). This instrument was calibrated against a standard micromanometer (Meriam Model 34 FB 2 TM). No detectable difference was observed between the two within the range of the present measurements.

The D.C. output of the pressure meter, which indicates the pressure difference in millimeters of mercury between the dynamic pressure and the static pressure, was connected to the Y-axis of a X-Y recorder (Moseley Type 135). The X-axis of the recorder can be either set in time to obtain the time integration for a point-by-point measurement, or connected to the voltage output of the potentiometer of the carriage to obtain a continuous vertical profile of the pressure difference between the dynamic and the static heads. At each station, the continuous profile was made and later checked by means of pointby-point measurements. After a smoothing process, the continuous profile was later digitized using an optical-traced digitizer, and computed to obtain mean velocities. From a knowledge of the pressure difference between dynamic and static pressures, the barometric

pressure, and the temperature, the mean velocity was calculated according to

$$U = 2.36 \int \Delta h/\rho \tag{4-1}$$

where U is the mean velocity in ft/sec,  $\rho$  is the density of the air in slugs/ft<sup>3</sup>, and  $\Delta h$  is the pressure difference between dynamic and static pressures in millimeters of mercury.

Possible sources of errors in pitot tube measurement of mean velocity have been discussed by Goldstein (1938), MacMillan (1956), Davies (1958), Bradshaw (1964) and Sandborn (1966). These errors are caused by the effects of the proximity of the wall, the local shear stress and the turbulence. Except very close to the wall, the possibility of errors due to shear and proximity of the wall are negligible. The turbulence effects can not be overlooked in the inner portion of the turbulent boundary layer where the turbulent intensities are very high. The importance of this effect decreases with distance from the wall. Correction for the turbulence effects of the pitot tube measurements following Goldstein, was found to vary from -1.7 to +0.1% of the local mean velocity. Hinze (1959) has criticized Goldstein's method which assumes that spanwise velocity fluctuations affect the pitot tube reading the same way as streamwise fluctuation, and has suggested that the effect of the former is appreciable less and may be even of opposite nature. For the present investigation no such correction was applied to the pitot tube measurement of mean velocity.

### (B) Hot-Wire for Mean Velocity

The hot-wire measurement of mean velocity was made with a  $2x10^{-4}$  inch diameter platinum coated tungsten wire, which was held parallel to the flat plate and perpendicular to the direction of the flow. The hot-wire was operated by a constant-temperature anemometer system (DISA 55D01). The voltage output of the anemometer was recorded on a X-Y recorder (Moseley Model 135) for a suitable time interval; later, an average was obtained from the records. The mean velocity was obtained from the D.C. voltage output through a calibration curve of the hot-wire. The procedure of the hot-wire calibration will be discussed in Section 4.3.4.

Fig. 4.9 illustrates a typical comparison of mean velocity profiles measured by the hot-wire and the pitot tube. The hot-wire results were generally in good agreement with pitot tube values, except in a region very close to the wall. This discrepency is primarily due to turbulence effects.

# 4.3.3 Measurement of Wall Shear Stress

Measurement of local wall shear stress along a smooth strip of the flat plate in the boundary layer was made by Preston's method using a round surface pitot tube attached to the surface (the so-called Preston tube). By assuming that close to the surface there exists a region in which the flow is substantially determined by the wall shear stress and the physical properties of the fluid, Preston (1954) obtained a universal non-dimensional relation between the wall shear stress and the difference between the pressure in the Preston tube and the static pressure on the surface

$$\frac{\tau_w d_1^2}{\rho v^2} = \mathbf{f}_p \left(\frac{\Delta p d_1^2}{\rho v^2}\right)$$
(4-2)

where  $\tau_w$  is the wall shear stress,  $d_1$  is the diameter of the Preston tube,  $\Delta p$  is the Preston tube reading (i.e. the difference between pitot tube and static pressures), and the function  $f_p$  is determined from the calibration of the Preston tube.

Experiments by Head and Rechenberg (1962), Ferriss (1965) and Patel (1965) confirmed the assumption of universal wall similarity common to flat-plate boundary layers and fully-developed pipe flows, and provided the convincing evidence of the soundness of Preston's method of measuring wall shear stress. However, Patel found that the Preston's method overestimates the wall shear stress in both strong favourable and adverse pressure gradients. The inaccuracy of the method is caused by large deviations from the wall similarity due to the existence of the pressure gradients.

For the present investigation a homemade pitot tube of diameter 1/16 inch was built and used to measure local wall shear stress in the turbulent boundary layer of nearly zero streamwise pressure gradient. Care was taken to ensure that the mouth of the tube was free from burrs and deformities and that the tube was lying along the surface of the flat plate as closely as possible. Additional pressure taps of 1/16 inch diameter hole at proper locations were constructed for the wall shear stress measurement. The tube was first placed on the surface, parallel to the centerline of the plate, on one side of a static pressure tap, and was 1/4 inch apart between the tube and the pressure tap. The pressure difference between the two was recorded by means of a pressure meter (MKS Baratron Type 77). The same arrangement was repeated again on the other side of the pressure tap. The average value of two readings was used to determine the wall shear stress at the station of the pressure tap. With the calibration formula given, the wall shear stress could be readily calculated. The calibration formula, provided by the staff of the Aerodynamics Division, N. P. L. (1958), for the flat plate

$$\log_{10} \frac{\tau_{w} d_{1}^{2}}{4\rho v^{2}} = \overline{2.647} + 0.875 \log_{10} \frac{\Delta p d_{1}^{2}}{4\rho v^{2}}$$
(4-3)

was used.

Measurement of wall shear stress using a Preston tube indicates good agreement with other measurements in the turbulent boundary layer along a smooth surface. Along a rough surface, however, the method casts considerable doubt on the accuracy of the measurement. The individual roughness produce local mean velocity fluctuations which in turn produce local static pressure variations near the wall. The surface pitot tube and wall pressure taps can not be used to give meaningful Preston tube readings.

### 4.3.4 Measurement of Turbulence

Two types of hot-wire probes were used to measure the turbulence quantities. One was a single wire placed normal to the flow, to measure  $\overline{u^2}$ . The same wire was also used on mean velocity measurement. Another was a crossed wire to measure  $\overline{v^2}$ ,  $\overline{w^2}$ ,  $\overline{uv}$  and  $\overline{uw}$ . The crossed wire was made of two hot-wires mounted on a single probe in

the shape of X as in conventional hot-wire anemometry. The measurements of these turbulence quantities depend on the plane of axis in which the wires were operated. All the wires were made of platinum coated tungsten wire of a  $2x10^{-4}$  inch diameter. A schematic diagram of the single wire measurement instrumentation is shown in Fig. 4.10. For crossed wire measurements similar arrangements were made with two hot-wire anemometer systems.

(A) Calibration of Hot-Wires

Prior to the hot-wire calibration, each hot-wire was subjected to the "curing" precess. During the curing process the hot-wire was operated in no flow condition for at least 24 hours. This operation would ensure a stable condition of the hot-wire characteristics before any calibration of measurement was made. The hot-wire was calibrated against a 1/16 inch diameter standard pitotstatic tube in the free stream outside the boundary layer. The calibration curve was obtained by measuring D.C. voltage output of the wire for a range of mean velocities.

During the experiments the hot-wire was calibrated every two hours. No detectable change on the calibration curve was found for a properly "cooked" wire. A typical calibration curve of the hot-wire is shown in Fig. 4.11.

(B) Calculations of Turbulent Intensities and Shear Stresses

Baldwin, Sandborn and Lawrence (1960) found that the heat loss from a hot-wire is a function of local velocity, temperature and fluid properties, and angle of attack. For the present investigation, temperature and fluid properties were considered to be constant. Following the derivations given by Sandborn (1967), the response

equation of the wire held in the x-z plane at an angle  $\alpha$  to the direction of the flow is given by

$$\mathbf{e} = \frac{\partial \mathbf{E}}{\partial \mathbf{U}} \mathbf{u} + \frac{1}{\mathbf{U}} \frac{\partial \mathbf{E}}{\partial \alpha} \mathbf{w}$$
(4-4)

where e represents the fluctuation output of the wire, and  $\frac{\partial E}{\partial U}$  and  $\frac{\partial E}{\partial \alpha}$  are respectively the sensitivities of the wire to local velocity and angle of attack.

Letting 
$$\frac{\partial E}{\partial U} = S_U$$
 and  $\frac{1}{U} \frac{\partial E}{\partial \alpha} = S_\alpha$ , one obtains  
 $e = S_U u + S_\alpha w$  (4-5)

Squaring and averaging this equation results in

$$\overline{e^2} = S_U^2 \overline{u^2} + 2 S_U S_\alpha \overline{uw} + S_\alpha^2 \overline{w^2}$$
(4-6)

Arya (1968) found that the relation between  ${\rm S}_{\rm U}$  and  ${\rm S}_{\alpha}$  should have the form

$$S_{\alpha} = c \cot \alpha S_{\mu}$$
(4-7)

and

$$c = \frac{1 - \sigma^2}{1 + \sigma^2 \cot \alpha}$$
(4-8)

where  $\sigma$  is a parameter, proposed by Hinze (1959), to indicate the effective heat transfer from the velocity component parallel to the wire. The variation of  $\sigma$  with wire sizes and velocities has been studies by Webster (1962) and Champagne et al. (1967). It is found that  $\sigma$  depends primarily on the wire length-to-diameter ratio. Wires of different materials have less effect in the values of  $\sigma$ . For the present study  $\sigma = 0.2$  is used.

Substituting Eq. (4-7) into (4-6), we obtain

$$\overline{e^2} = S_U^2 \left( \overline{u^2} + 2c \ \overline{uw} \ \cot\alpha + c^2 \overline{w^2} \ \cot^2 \alpha \right)$$
(4-9)

In the following discussion different wire configurations are considered for the turbulence quantities that are measured.

(i) <u>Calculating</u>  $\overline{u^2}$  from the Single Wire

The streamwise turbulent flux  $u^2$  was determined from the r.m.s. output to the local mean velocity. When the wire is held in a plane normal to the flow, i.e.  $\alpha = 90^{\circ}$ . it follows from Eq. (4-9) that

$$\overline{u^2} = \frac{\overline{e^2}}{S_U^2}$$
(4-10)

The velocity sensitivity  $S_U$  can be obtained either by measuring the slope of the hot-wire calibration curve, or by differentiating an empirical curve based on King's law. The latter approach has been found to be more convenient and to give better results. If one plots the squared voltage output of the wire  $E^2$  against the square root of corresponding velocity  $\sqrt{U}$ , a best fit curve can be obtained by a least-square curve fitting process,

$$E^2 = A' + B' \sqrt{U}$$
 (4-11)

where A' and B' are constants.

The velocity sensitivity  $S_{11}$  is then given by

$$S_{U} = \frac{\partial E}{\partial U} = \frac{1}{4E} \frac{B'}{\sqrt{U}}$$
(4-12)

Substituting Eq. (4-12) into Eq. (4-10), one obtains

$$\overline{u^{2}} = \frac{16U (A' + B'\sqrt{U})}{B'^{2}} \overline{e^{2}}$$
(4-13)

(ii) Calculating  $\overline{v^2}$ ,  $\overline{w^2}$ ,  $\overline{uv}$  and  $\overline{uw}$  from the Crossed Wire

For measuring turbulent quantities  $\overline{w^2}$  and  $\overline{uw}$ , the crossed wire was held in the x-z plane, as shown in Fig. 4.12. For convenience the angles between the direction of mean velocity and the crossed wire are kept at  $\alpha = \pm 45^{\circ}$ .

Eq. (4-9) follows, when  $\alpha = 45^{\circ}$ ,

$$\overline{e_1^2} = S_1^2 (\overline{u^2} + 2c \,\overline{uw} + c^2 \,\overline{w^2})$$
(4-14)

and when  $\alpha = -45^{\circ}$ ,

$$\overline{e_2^2} = S_2^2 (\overline{u^2} - 2c \,\overline{uw} + c^2 \,\overline{w^2})$$
 (4-15)

where  $\overline{e_1^2}$  and  $\overline{e_2^2}$  represent respectively the r.m.s. response of the crossed wire at  $\alpha = \pm 45^\circ$ , and  $S_1$  and  $S_2$  represent respectively the sensitivity of the crossed wire to local velocity at  $\alpha = \pm 45^\circ$ .

After rearranging Eqs. (4-14) and (4-15), and solving simultaneously for  $\overline{w^2}$  and  $\overline{uw}$ , one obtains

$$\overline{w^{2}} = \frac{1}{2c^{2}} \left\{ \left( \frac{\overline{e_{1}^{2}}}{S_{1}^{2}} + \frac{\overline{e_{2}^{2}}}{S_{2}^{2}} \right) - 2 \overline{u^{2}} \right\}$$
(4-16)

$$\overline{uw} = \frac{1}{4c} \left( \frac{\overline{e_1^2}}{s_1^2} - \frac{\overline{e_2^2}}{s_2^2} \right)$$
(4-17)

In the same fashion,  $v^2$  and  $\overline{uv}$  can be measured by operating the crossed wire in the x-y plane, and calculating equations similar to Eqs. (4-16) and (4-17).

#### (C) Errors in Hot-Wire Measurements

Hot-wire measurements of turbulence are subject to various type of errors. Random errors are produced in the process of calibration, reading instruments, alignment of probes, etc. Each of the random errors is independent of the others. The error of measured result can be estimated through the statistical method of repeating the measurements numerous times. The method for calculating the errors in single-sample experiments are discussed by Kline and McClintock (1953). Systematic errors are introduced due to finite length of the hot-wire, proximity of solid boundary, gradients of mean velocity and turbulent intensities, etc. Details of the systematic errors are discussed by Tieleman (1967). Most of these errors become significant only when measurements are made very close to the wall. Errors due to gradients of mean velocity and turbulent intensities are important only in measurements with crossed wires. The crossed wire measurements are furthermore affected by the correlation between the turbulent components over a small separation distance between two wires. In the absence of any suitable technique available for correcting the errors, no correction was applied in the present turbulence measurements.

### 4.3.5 Measurement of Frequency Spectra

For the purpose of measuring spectrum and various second order correlations of turbulent components, the fluctuating signals of hot-wire anemometers for single wire as well as crossed wire, at selected points in the boundary layer, were recorded on FM magnetic tape. Because of the low fluctuating outputs of the anemometer and a maximum of 1 volt r.m.s. recording limit to the tape recorder (Ampex FM 1300), the signals were required to pass through the low-level preamplifiers (Tektronix Type 122) and a proper attenuating stage before recording.

For the present study only the frequency spectra of the streamwise turbulent velocity were analyzed. The mean-square output between frequencies n and n + dn was automatically recorded by feeding the recorded hot-wire signal into a wave analyzer (General Radio Type 1911) having a set of passive filters from 2.5 to 25,000 Hz and two fixed bandwidths. Only the 1/3-octave (23%) bandwidth was used in the spectrum analyses.

The frequency spectrum of  $\overline{u^2}$  is

$$\frac{1}{u^2} \int_0^\infty E(n) \, dn = 1 \tag{4-18}$$

where E(n) is the fraction of energy between frequencies n and n + dn, and is known as the frequency-density function. In terms of the meansquare output signal of the wave analyzer, the fraction of energy at each frequency is

$$E(n) = \frac{1}{B_{W}} e^{2}(n, B_{W})$$
(4-19)

where  $e^2(n, B_W)$  is the mean-square output signal of the wave analyzer at any selected frequency n, and  $B_W$  is the filter bandwidth.

The spectral measurements are subject to errors due to finite hot-wire length, filter bandwidth, and noise. Due to finite length of the hot-wire, two or more eddies of small size may strike the wire simultaneously with the result that the measured power would be larger than the true power. This will cause the inaccuracy of the spectrum data at the high-frequency end. For a constant-percent bandwidth of 1/3-octave (23%), the side-band effect is negligible if the response time at each frequency (rise and fall time) is short enough.

# CHAPTER V

# ANALYSIS OF EXPERIMENTAL DATA

Extensive measurements of mean value and turbulence quantities for the turbulent boundary layer over a flat plate consisting of alternate, longitudinal strips of smooth and rough surfaces were obtained. These measurements included mean velocity distribution, wall shear stress, turbulent intensities, turbulent shear stresses and energy spectra of the streamwise turbulent velocity. The experimental investigation was limited to one type of surface configuration, as shown in Fig. 4.2. Two free-stream velocities of 20 and 30 ft/sec were chosen to give different flow conditions. The streamwise pressure gradient in the free stream was adjusted so as to be nearly zero. The spanwise periodic nature of the surface boundary conditions suggested that the entire flow field in the boundary layer is reproducible when the measurements are made in a region bounded by planes of symmetry of a smooth strip and an adjacent rough strip. In the present study, measurements were thus concentrated in a narrow region from the centerline of the smooth strip (y = 0 in.) to that of the rough strip (y = 6 in.). An analysis of the experimental data is presented in this chapter.

### 5.1 Boundary Layer Development

Various methods have been used to specify the height of a boundary layer. The boundary layer  $\delta$  is defined as the vertical distance from the wall where the mean velocity differs by one percent

from the free-stream velocity. The displacement thickness  $\delta^*$  and the momentum thickness  $\theta$  are defined by

$$\delta^* = \int_0^\infty \left(1 - \frac{U}{U_\infty}\right) dz$$
(5-1)

$$\theta = \int_{0}^{U} \frac{U}{U_{\infty}} \left(1 - \frac{U}{U_{\infty}}\right) dz$$
 (5-2)

The ratio  $H = \delta^* / \theta$ , known as the shape factor, is usually chosen as a parameter characterizing the mean velocity profile in the boundary layer.

The values of  $\delta^*$  and  $\theta$  can be computed with reasonable accuracy by numerical integration of the mean velocity profiles. On the other hand the uncertainty in  $\delta$  is large because the mean velocity in the boundary layer approaches the free-stream velocity asymptotically. A summary of boundary layer parameters that describe the experimental flow conditions at various locations is given in Tables 5.1a-b.

Figs. 5.1a-b show the boundary layer growth over the flat plate along the centerlines of the smooth strip (y = 0 in.) and the rough strip (y = 6 in.). Figs. 5.2 through 5.4 are the plots of  $\delta^*$ ,  $\theta$  and H as a function of x, respectively. The various thickness  $\delta$ ,  $\delta^*$  and  $\theta$  increase with increasing streamwise distance, but decrease with increasing free-stream velocity. The shape factor H appears to be nearly constant under the same conditions of the wall. The values of  $\delta^*$ ,  $\theta$  and H for the smooth surface are, in general, smaller than those for the rough surface. However,  $\delta$  is roughly the same for both smooth and rough surfaces. Since the boundary layer thickness cannot be accurately measured, a spanwise averaged value of  $\delta$  is used for the analysis of experimental data.

#### 5.2 Determination of Shear Velocity

The shear velocity  $U_*$ , which is an important velocity scale used in the analysis of mean velocity data, is defined by

$$U_{\star} = \int \frac{\tau_{w}}{\rho}$$
(5-3)

It is also related to the local skin-friction coefficient  $\ensuremath{\,^C_f}$  by the relation

$$\sqrt{\frac{C_{f}}{2}} = \frac{U_{\star}}{U_{\infty}}$$
(5-4)

The physical basis for the determination of the shear velocity is provided by the postulation of the logarithmic mean velocity profile in the region near the wall. The shear velocity is thus obtained from the slope of the measured logarithmic profiles in accordance with Eq. (2-4) or Eq (2-10). The results, as shown in Fig. 5.5, indicate that the shear velocity reacts rapidly to a lateral change in surface roughness.

For the smooth strip, the Preston-tube method (Section 4.3.3) has been used to measure the wall shear stress. The results are listed in Tables 5.2a-b. The values of the resulting shear velocity are in good agreement with those calculated from the well-known Ludwieg and Tillmann's (1950) empirical formula for two-dimensional turbulent boundary layer over smooth walls,

$$U_{\star} = U_{\infty} \{ 0.123 \times 10^{-0.678 \text{H}} \times \left( \frac{U_{\infty \theta}}{v} \right)^{-0.268} \}^{1/2}$$
(5-5)

but are a bit higher than those obtained using the logarithmic law of mean velocity method, as shown is Fig. 5.6. This is probably due to the small departure of measured mean velocity profiles from a twodimensional state.

In principle the shear velocity may be obtained by extraplating the turbulent shear stress  $-\rho \overline{uw}$ . For the smooth surface the values obtained by the extrapolation of the turbulent shear stress and by the logarithmic law of mean velocity method are in fair agreement. However, for the case of the rough surface such extrapolation of the turbulent shear stress gives considerably lower shear velocities in comparison with those determined from the mean velocity profiles. This is because the pressure distribution on the roughness elements can produce a form drag which may contribute significantly to the total wall shear stress. Therefore, the turbulent shear stress obtained from the crossed hot-wire anemometers does not represent the true wall shear stress on any rough surface.

### 5.3 Mean Velocity Field

The mean velocity distribution across the boundary layer has been measured at x-stations 1, 2, 3, 4, 6, 8 and 10 ft. with the corresponding y-stations 0, 1, 2, 2.5, 3, 3.5, 4, 5 and 6 inches using two free-stream velocities of 20 and 30 ft/sec. The mean velocity data are tabulated in Tables 5.3a-b.

5.3.1 Vertical Distribution of Mean Velocity

Since the secondary flow generated by the present surface configuration is very small, the two-dimensional descriptions for similarity velocity laws can approximately be applied (Cumpsty, 1970).

Three types of non-dimensional plots have been used to study and compare the mean velocity distribution normal to the wall. They are:

$$\frac{U}{U_{\infty}} = \left(\frac{z}{\delta}\right)^{1/m}$$
(5-6)

$$\frac{J}{J_{\star}} = f\left(\frac{zU_{\star}}{v}\right)$$
(5-7)

and

$$\frac{U_{\infty} - U}{U_{\star}} = g\left(\frac{z}{\delta}\right)$$
(5-8)

Eqs. (5-6) through (5-8) are known as the power law, the law of the wall and the velocity-defect law, respectively.

#### (A) Power Law Profiles

To test the validity of the power law (Eq. 5-6), mean velocity data were plotted as  $U/U_{\infty}$  versus  $z/\delta$ . The resulting velocity profiles along the centerline of the smooth strip (y = 0 in.) and the rough strip (y = 6 in.), as presented in Figs. 5.7a-b, are not similar but form a family of curves for different Reynolds number  $\operatorname{Re}_{x}=U_{\infty}x/v$  and degree of roughness of the wall. This result agrees with the findings of Hama (1954) and Clauser (1956). The parameter m in the power law, determined by the method of least squares, is plotted versus the streamwise coordinate x in Fig. 5.8. The value of m for the rough surface is always smaller than that for the smooth surface under the same flow conditions.

Spanwise comparisons of power law profiles at x = 6 and 10 ft. for  $U_{\infty} = 30$  ft/sec are presented in Figs. 5.9a-b. The resulting velocity profiles are similar in the outer 40 percent of the boundary layer but form a family of curves below this portion. This implies that the spanwise disturbing effect of wall roughness is limited to the inner 60 percent of the boundary layer. The spanwise variation of m in the power law profile takes place gradually from the centerline of the smooth strip to that of the rough strip.

### (B) "Law of the Wall" Profiles

From the semi-logarithmic plots of  $U/U_{\infty}$  versus log z as shown in Figs. 5.10a-b, two significant facts are noted. First, all mean velocity profiles coincide for large value of z and diverge near the wall. This indicates that the effect of the spanwise changes of wall roughness depends not only on the distance from the wall but also on the roughness difference between the rough strip and the smooth strip. Secondly, the mean velocity profiles in the inner portion of the boundary layer are logarithmic in character. It follows from this fact that there still exists a wall region where the law of the wall applies, resulting in a logarithmic velocity distribution.

To illustrate the validity of the wall law given in Eq. (5-7), the mean velocity data are plotted in the form  $U/U_*$  versus Log  $zU_*/v$ as shown in Figs. 5.11a-b, where  $U_*$  is the local shear velocity. Note that the experimental data have been forced to fit the twodimensional description for the wall law by the prodedure used to determine the shear velocity. Hence, a well-defined logarithmic region exists near the wall. In the wall region all velocity profiles appear to fall onto straight lines with the same logarithmic slope. The range of validity of the wall law extends from  $zU_*/v = 20$  to

 $zU_*/v = 150$  in the smooth part of the wall, and from  $zU_*/v = 100$  to  $zU_*/v = 500$  in the rough part of the wall. The thickness of the wall law region becomes smaller as it approaches the roughness discontinuity. At the plane of symmetry of the smooth strip (y = 0 in.), the law of the wall can be expressed as

$$\frac{U}{U_{\star}} = \frac{1}{\kappa} \ln \frac{zU_{\star}}{v} + A$$
(5-9)

where the values of the empirical constants  $\kappa$  and A are 0.41 and 6.6 respectively. For two-dimensional smooth-wall flows, there is appreciable scatter in the values of  $\kappa$  and A because different values of these constants have been proposed by different investigators. Opinions about the value of Kārmán's constant  $\kappa$  are less divergent. The value of  $\kappa$  is usually taken as 0.41 for flat-plate flows. The choice of the constant A, however, differs from a value of 4.9 (Clauser, 1956) to that of near 7 (Townsend, 1956). The present suggested value of A which is obtained from the best fit of the experimental data is still within this wide range of choice. It is therefore suggested that the mean velocity distributions along the centerline of the smooth strip is closely analogous to those in two-dimensional smooth-wall flows.

Away from the smooth-wall flow region, the law of the wall can be written, in general, in the following form:

$$\frac{U}{U_{\star}} = \frac{1}{\kappa} n \frac{zU_{\star}}{v} + A - \frac{\Delta U}{U_{\star}}$$
(5-10)

where  $\Delta U/U_{\star}$  represents the downward shift of a profile from the smooth

wall profile in the wall-law plot. The absolute value of  $\Delta U/U_{\star}$  depends on the shape and distribution of the roughness elements and is a function of  $k_r U_{\star}/v$ , where  $k_r$  is the roughness height. A relation exists between  $\Delta U/U_{\star}$  and  $k_r U_{\star}/v$  as given in Fig. 12. The roughness effect  $\Delta U/U_{\star}$  measured from Figs. 5.11a-b is plotted versus the spanwise coordinate y in Fig. 5.13. The result indicates that  $\Delta U/U_{\star}$  can be well predicted by the relation

$$\frac{\Delta U}{U_{\star}} = U_{\infty} \left( \frac{1}{U_{\star S}} - \frac{1}{U_{\star}} \right)$$
(5-11)

where the subscript "s" refers to the wall condition at the plane of symmetry of the smooth strip. It also reveals that the spanwise adjustment from a smooth-wall flow to a rough-wall flow takes place almost immediately in the vicinity of the roughness discontinuity.

### (C) Velocity-Defect Profiles

The mean velocity data, in conjuction with local shear velocity, can now be applied to test the validity of the velocitydefect law given in Eq. (5-8). Since  $\delta$  is not well defined by the experimental data, Rotta (1962) suggested that the velocity-defect law may be expressed as

$$\frac{U_{\infty} - U}{U_{\star}} = g_1 \left( \frac{zU_{\star}}{\delta^{\star} U_{\infty}} \right)$$
(5-12)

by introducing a non-dimensional wall distance  $zU_*/\delta^*U_{\infty}$  instead of  $z/\delta$ . The measured velocity data are therefore presented in the form  $(U_{\infty} - U)/U_*$  versus  $z/\delta$  in Figs. 5.14a-b, and versus  $zU_*/\delta^*U_{\infty}$  in Figs. 5.15a-b. The resulting velocity profiles show considerable

but systematic variations depending on the distribution of local wall shear stress. This implies that the universality of the velocitydefect law for the outer region of the boundary layer breaks down for the flow of this kind.

At the plane of symmetry of the smooth strip (y = 0 in.), the straight-line portion of the semi-logarithmic plots of mean velocity data can be described by the relation

$$\frac{U_{\infty} - U}{U_{+}} = -\frac{1}{\kappa} \ln \frac{z}{\delta} + 2.5$$
 (5-13)

or

$$\frac{U_{\infty} - U}{U_{\star}} = -\frac{1}{\kappa} \ln \frac{zU_{\star}}{\delta^{\star}U_{\infty}} - 0.4$$
(5-14)

The velocity profiles show deviations from the straight-line portion at  $z/\delta \approx 0.15$  or  $zU_*/\delta^*U_\infty \approx 0.04$ . Apparently this value of  $z/\delta$  or  $zU_*/\delta^*U_\infty$  marks roughly the boundary between the wall region and the outer region of the boundary layer. In the outer region of the boundary layer, the velocity profiles show fair agreement with the empirical relations for two-dimensional smooth-wall flows (Hama, 1954),

$$\frac{U_{\infty} - U}{U_{\star}} = 9.6 \left(1 - \frac{z}{\delta}\right)^2$$
(5-15)

or

$$\frac{U_{m} - U}{U_{\star}} = 9.6 \left(1 - 3.33 \frac{zU_{\star}}{\delta^{\star}U_{\infty}}\right)^{2}$$
(5-16)

The velocity correlations given in Eqs. (5-13) through (5-16) indicate that the outer portion of the boundary layer along the centerline of the smooth strip is also analogous to those in two-dimensional smooth-wall flows.

Away from the smooth-wall flow region, the velocity profiles form a family of curves for different wall shear stress. Several different versions of the velocity-defect plot have been examined. When the mean velocity data are replotted as  $2(U_{\infty}-U)/(U_{*s}+U_{*})$  versus  $z/\delta$  in Figs. 5.16a-b, and versus  $zU_{*s}/\delta_{s}^{*}U_{\infty}$  in Figs. 5.17a-b, all velocity profiles fall within a resonably narrow band in an apparently universal fashion. This suggested that all the velocity profiles in the outer portion of the boundary layer may be approximated by a modified velocity-defect expression of the form

$$\frac{U_{\infty} - U}{U_{\star}} = \frac{1}{2} \left( \frac{U_{\star s} + U_{\star}}{U_{\star}} \right) g\left( \frac{z}{\delta} \right)$$
(5-17)

or

$$\frac{U_{\infty} - U}{U_{\star}} = \frac{1}{2} \left( \frac{U_{\star s} + U_{\star}}{U_{\star}} \right) g_{1} \left( \frac{zU_{\star s}}{\delta_{s}^{\star} U_{\infty}} \right)$$
(5-18)

The functions g and  $g_1$  correspond to the velocity-defect distribution in a two-dimensional smooth-wall flow. The effect of spanwise variation of wall roughness enters through the variable  $(U_{*s}+U_{*})/2U_{*}$ .

### 5.3.2 Spanwise Distribution of Mean Velocity

Figs. 5.18a-b show the spanwise distribution of mean velocity at constant heights in vertical plane x = 6 and 10 ft. for  $U_{\infty} = 30$ ft/sec. The mean velocity distribution far above the wall is not affected by the presence of spanwise changes of wall roughness. At low elevations, however, the mean velocity varies in the spanwise direction. At very low elevations, only local surface conditions dominate the profile. The present study can thus be treated as a small perturbation about a two-dimensional smooth-wall flow, provided that the secondary flow in the boundary layer is small.

The effect of the transition from smooth-wall conditions to rough-wall conditions on the curves of constant mean velocity (or isovels) is shown in Figs. 5.19a-b. The spanwise distortion of the mean velocity distribution near the roughness discontinuity may be explained by the existence of a secondary flow of Prandtl's second kind (Section 3.3). This secondary flow, as suggested by Hinze (1967), is directed towards the wall in the plane of symmetry of a rough strip, then along the wall towards the smooth part of the wall, and away from the wall in the plane of symmetry of a smooth strip. In the present study, the measurement of secondary-flow components by means of a rotating crossed wire indicates that a weak cross-flow exists and is directed from a rough strip to a smooth strip near the wall. The magnitude of this cross-flow is everywhere less than 3 percent of the free-stream velocity. The strength of the cross-flow decreases with increasing wall distance. Due to various errors introduced in the hot-wire measurements (see Section 4.3.4), the accuracy of the

measured secondary flow is poor. The secondary flow data were thus excluded.

### 5.4 Turbulence Field

The measurements of the turbulent intensities  $\sqrt{u^2}$ ,  $\sqrt{v^2}$  and  $\sqrt{w^2}$ , and the turbulent shear stresses  $-\rho u \overline{v}$  and  $-\rho u \overline{w}$  were made in the boundary layer at x-stations 6 and 10 ft. with the corresponding y-stations 0, 1, 2, 2.5, 3, 3.5, 4, 5 and 6 inches using two freestream velocities of 20 and 30 ft/sec. The free stream had a turbulence level within 2 percent. The experimental data are given in Tables 5.4a-b, as well as in Figs. 5.20 through 5.22. No attempt has been made to measure the turbulent shear stress  $-\rho \overline{v} \overline{w}$  in the boundary layer. However, an estimate of the  $-\rho \overline{v} \overline{w}$  distribution based on a theoretical analysis is obtained (Section 3.3).

Figs. 5.20a-b show the intensities of the streamwise, spanwise and vertical turbulent velocities  $\sqrt{u^2}$ ,  $\sqrt{v^2}$  and  $\sqrt{w^2}$  relative to the free-stream velocity in the boundary layer along the centerline of the smooth strip (y = 0 in.) and of the rough strip (y = 6 in.), respectively. In all cases the values of the turbulent intensities for the smooth strip are smaller than the corresponding values for the rough strip. The turbulent intensities approach the free-stream values at a greater distance from the wall than does the mean velocity. The intensity of the streamwise turbulent velocity reaches a maximum value, while the other two turbulent intensities remain nearly constant, in a region close to the wall. The present high-turbulence data have been compared with the smooth-wall data of Chowdhury (1966) and the curve shows nearly the same trends and locations of maximum values of  $\sqrt{u^2}/U_{ex}$ . However, the measured turbulent intensities in the present case are found to be higher than Chowdhury's experimental data in the identical wind tunnel. This increase in the turbulent intensities may be due to continuous production of disturbances which diffuses outwards from the leading-edge roughness elements and along side edges of each of the rough strips as the flow progresses downstream. It may be noticed that the three turbulent intensities differ appreciably from one another over the inner portion of the boundary layer. The intensity of the streamwise turbulent velocity has the highest value; that of the vertical turbulent velocity has the smallest value. The difference between  $\overline{v^2}$  and  $\overline{w^2}$  is nearly constant in the region close to the wall, but depends strongly on the vertical distance over much of the whole boundary layer.

The distributions of  $-\overline{uv}$  and  $-\overline{uw}$  relative to the free-stream velocity in the boundary layer along the centerline of the smooth strip (y = 0 in.) and of the rough strip (y = 6 in.) are shown in Figs. 5.21a-b. The values of  $-\overline{uv}$  and  $-\overline{uw}$  increase gradually from near the wall to their maximum values at very short distances, then decrease. The value of  $-\overline{uv}$ , although very small, indicates that the turbulence field at the centerline of each strip is not "strictly" two-dimensional.

Spanwise comparisons of turbulence profiles  $\overline{u^2}$ ,  $\overline{v^2}$ ,  $\overline{w^2}$ ,  $-\overline{uv}$ , - $\overline{uw}$  and - $\overline{vw}$  are presented in Figs. 5.22a-f, where the - $\overline{vw}$  distribution is estimated theoretically according to Eq. (3-21). The results indicate that the spanwise distributions of these turbulence quantities depend not only on the underlying surface characteristics, but also on the roughness difference between the rough strip and the smooth strip.

The modifications of  $\overline{u^2}$ ,  $\overline{v^2}$ ,  $\overline{w^2}$  and  $-\overline{uw}$  takes place gradually from the centerline of the smooth strip to that of the rough strip. The values of  $-\overline{uv}$  and  $-\overline{vw}$  attain their maximum values at the roughness discontinuity, then decrease gradually towards the centerlines of both smooth and rough strips.

### 5.5 Energy Spectra of Streamwise Turbulent Velocity

It is possible to gain some further insight into the boundary layer flow by examining the energy spectrum of the streamwise turbulent velocity. The energy enters the spectrum through large eddies and is then transferred down the spectrum to smaller eddies where it is finally dissipated.

One-dimensional energy spectra of the streamwise turbulent velocity evaluated at x-stations 6 and 10 ft. with the corresponding y-stations 0, 1, 2, 2.5, 3, 3.5, 4, 5 and 6 inches for the free-stream velocity of 30 ft/sec are given in Table 5.5. The data have been normalized with respect to  $\overline{u^2}$  such that

$$\frac{1}{u^2} \int_0^\infty E(k) \, dk = 1 \tag{5-19}$$

where  $k = 2\pi n/U$  is the one-dimensional wave number, and E(k) is the one-dimensional wave-number spectrum function. Over most of the range, the accuracy is on the order of  $\pm 10$  %. The accuracy is somewhat less at the two extremes because of large-amplitude fluctuations at low frequencies and because of low signal-to-noise ratio at high frequencies. The error due to finite length of hot-wire increases as the scale of the turbulence decreases and becomes significant at the higher frequencies. However, because of the lesser accuracy of measurement in this range, no wire-length corrections were made.

Normalized spectra of the streamwise turbulent velocity at x = 6 ft, y = 0 and 6 inches for U = 30 ft/sec are shown in Figs. 5.23a-b. The normalized spectra show a strong dependence on  $z/\delta$ . Far away from the wall where the transfer among eddies predominates, there exists an extensive range where the spectra vary nearly according to  $k^{-5/3}$ . As the wall is approached where the production of turbulence becomes significant, the spectra show the tendency towards  $k^{-1}$  in the wavenumber range where  $k^{-5/3}$  normally exists when there is no mean-velocity gradient. This indicates a strong interaction between the mean-velocity gradient and the turbulent motion near the wall. The shape of the spectra shows a gradual transition from the -5/3 slope in the outer portion of the boundary layer to a -1 slope near the wall in accordance with the theoretical prediction by Tchen (1953). Nearly all of the spectra indicate the existence of a range varying as  $k^{-7}$  at the high wave-number end in agreement with Hisenberg's (1948) theory. The results also reveal that the contribution to the turbulent energy in the low wave-number range, that is, made by the large eddies, decreases as the wall approaches, but that the contribution in the high wavenumber range is increased.

Spanwise comparisons of the normalized spectra at x = 6 ft. for approximately equal values of  $z/\delta$  are given in Figs. 5.24a-c. Except at the low wave-number end, all spectra collapse on a single curve. At the low wave-number end, the turbulent energy contained in large eddies decreases gradually while going from a smooth to a rough strip.

# CHAPTER VI

# CONCLUSIONS

The present study has been carried out with the specific goal of increasing our knowledge of the flow characteristics of the turbulent boundary layer over a flat plate consisting of alternate, longitudinal strips of smooth and rough surfaces, at zero pressure gradient. The experimental investigation was limited to one type of surface configuration with the free-stream velocity as the only variable. Special attention has been given to mean velocity and turbulence distributions under the combined influence of wall roughness and strip width at distances far downstream from the leading edge of the plate, where the local boundary layer thickness is comparable to the strip width but is very large compared to the roughness height.

Results on measurements have been discussed in the preceding chapter. Several conclusions can be drawn from these results and discussions.

(1) The presence of the streamwise roughness strips in an otherwise two-dimensional smooth-wall flow will indeed perturb the flow. If the roughness height is very small compared to the local boundary layer thickness, the entire flow field can be treated as a small perturbation about a two-dimensional state. At the planes of symmetry of a smooth strip and a rough strip, the flow conditions are nearly analogous to those in two-dimensional boundary layers over a smooth wall and a rough wall, respectively.

(2) The adjustment of the wall shear stress to local conditions of the wall takes place very rapidly while undergoing an abrupt, spanwise change in surface roughness.

(3) The mean velocity distribution can be represented adquately by a power law, except near the wall. However, the power law profiles do not coincide but form a family of curves for different Reynolds number and degree of wall roughness.

(4) There still exists a wall region where the two-dimensional description for the law of the wall applies, resulting in a logarithmic velocity distribution. The thickness of this wall region becomes smaller as it approaches the roughness discontinuity. Away from the plane of symmetry of a smooth strip, the wall law shows a downward shift of mean velocity profiles from the smooth-wall profile, depending on the distribution of the local wall shear stress. The spanwise adjustment from a smooth-wall flow to a rough-wall flow occurs almost immediately in the vicinity of the roughness discontinuity.

(5) The universality of the two-dimensional velocity-defect law for the outer portion of the boundary breaks down for the flow of this kind. The velocity-defect profiles for the smooth part of the wall are higher than those for the rough part of the wall.

(6) A weak cross-flow which provides strong mixing and momentum transport in the spanwise direction exists near the wall. This crossflow is directed from a rough strip to a smooth strip.

(7) The intensities of the streamwise, spanwise and vertical turbulent velocities differ appreciable from one another over the inner portion of the boundary. The intensity of the streamwise turbulent velocity has the highest value and that of the vertical turbulent velocity has the smallest value. The values of the turbulent

intensities for the smooth strip are smaller than the corresponding values for the rough strip. The spanwise adjustment of the turbulent intensities takes place at a slower rate than does the wall shear stress.

(8) The turbulent shear stresses  $-\rho \overline{uv}$  and  $-\rho \overline{vw}$  attain their maximum values at the roughness discontinuity, then decrease gradually towards the centerlines of both smooth and rough strips. However, the turbulent shear stress  $-\rho \overline{uw}$  increase gradually from the centerline of the smooth strip to that of the rough strip.

(9) The normalized spectra of the streamwise turbulent velocity show strong dependence on  $z/\delta$ . The shape of the spectra varies from  $k^{-5/3}$  law in the outer portion of the boundary layer to  $k^{-1}$  law near the wall. The spectra exhibit the tendency towards  $k^{-7}$  law at the high wave-number end. The spanwise distribution of the spectra remains largely unaffected by the surface configuration of the kind investigated here.

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## APPENDIX

## BOUNDARY LAYER FLOW OVER A THREE-DIMENSIONAL OBSTACLE

A mathematical model for laminar flow over a semi-infinite, rectangular-parallelepiped obstacle located along the centerline of a flat plate, as shown in Fig. A-1, is formulated. The three-dimensional effects, including cross-flow, is assumed to be small such that the flow field is in a perturbed state about a two-dimensional boundary layer. The essential ideas underlying the solution are based on Oseen's linearization of the boundary layer equation. Far away from the obstacle the flow will approximate the two-dimensional flow past a semi-infinite flat plate. Stewartson and Howarth (1960) solved this two-dimensional flow by building up successively from the potential flow associated with the free-stream velocity and a shear layer introduced to restore the no-slip condition on the plate. Near the obstacle a three-dimensional diffusive flow is determined by introducing a negative volume source to give zero mean velocity inside the obstacle. The solution thus obtained, neglecting the effects of the edges, shall give a first-order approximation of the three-dimensional boundary layer problem.

In the absence of longitudinal pressure gradient, Oseen's linearization of the boundary layer equation with small cross-flow takes the form

 $U_{\infty} \frac{\partial U}{\partial \mathbf{x}} = v \nabla^{2} U$ (A-1)
where  $\nabla^{2} \equiv \frac{\partial^{2}}{\partial y^{2}} + \frac{\partial^{2}}{\partial z^{2}}$  is a Laplace operator.

The coordinates are now non-dimensionalized by the following transformation:

$$\phi = U/U_{\infty}$$
  

$$\tau = x/d$$
(A-2)  

$$\eta = y \int \overline{\text{Re}_{d}} / d$$

$$\xi = z \int \overline{\text{Re}_{d}} / d$$

where d is a characteristic length, and  $\text{Re}_{d} = U_{\infty}d/\nu$  is the Reynolds number based on length d. Upon substitution, Eq. (A-1) becomes

$$\frac{\partial^2 \phi}{\partial \tau^2} = \frac{\partial^2 \phi}{\partial \eta^2} + \frac{\partial^2 \phi}{\partial \xi^2}$$
(A-3)

For the case of an semi-infinite, rectangular-parallelepiped obstacle of height d and width 2ed placed on a flat plate longitudinally starting from the leading edge of the plate, the boundary conditions are:

$$\tau < 0, \quad \phi = 1 \qquad \text{for all values of } \eta \text{ and } \xi$$
  
$$\tau > 0, \quad \phi = \begin{cases} 0 \qquad |\eta| \le e \overline{\text{Re}_d} \text{ and } 0 \le \xi \le \sqrt{\text{Re}_d} \qquad (A-4) \\ 0 \text{ otherwise} \end{cases}$$

The linearity of Eq. (A-3) with appropriate boundary conditions allows the solution to be separated into the potential flow associated with the free-stream velocity, a two-dimensional shear flow to restore the no-slip condition, and a three-dimensional perturbed, diffusive flow:

$$\phi = 1 + \phi_1 + \phi_2 \tag{A-5}$$

The two-dimensional shear flow is described by the eqaution

$$\frac{\partial \phi_1}{\partial \tau} = \frac{\partial^2 \phi_1}{\partial \xi^2}$$
(A-6)

with the boundary conditions:

$$\tau < 0, \qquad \phi_1 = 0 \qquad \text{for all values of } \xi$$
  
$$\tau > 0, \qquad \phi_1 = \begin{cases} -1 & \text{for } \xi = 0 \\ 0 & \text{as } \xi \neq \infty \end{cases}$$
(A-7)

This leads to the solution

$$\phi_1 = -\operatorname{erfc}\left(\frac{\xi}{2\sqrt{\tau}}\right) \tag{A-8}$$

where

$$\operatorname{erfc}(z) = \frac{2}{\sqrt{\pi}} \int_{z}^{\infty} e^{-t^{2}} dt = 1 - \operatorname{erf}(z)$$

Three-dimensional perturbed flow is described by the equation

$$\frac{\partial \phi_2}{\partial \tau} = \frac{\partial^2 \phi_2}{\partial \eta^2} + \frac{\partial^2 \phi_2}{\partial \xi^2}$$
(A-9)

with the boundary conditions:

$$\tau < 0, \quad \phi_2 = 0 \quad \text{for all values of } \eta \text{ and } \xi$$
  
$$\tau > 0, \quad \phi_2 = \begin{cases} -\operatorname{erf}(\frac{\xi}{2\sqrt{\tau}}) & |\eta| \le e\sqrt{\operatorname{Re}_d} \text{ and } 0 \le \xi \le \sqrt{\operatorname{Re}_d} \\ 0 & \text{otherwise} \end{cases}$$
(A-10)

Eq. (A-9) with the boundary conditions describes the diffusion from a semi-infinite, rectangular-parallelepiped volume source with variable strength. Following the technique given by Carslaw and Jaeger (1959), the solution of Eq. (A-9) is, in general, given by

$$\phi_2 = \frac{1}{4\pi} \iiint_{-\infty}^{\infty} \frac{\phi_2(\mathbf{t}, \alpha, \beta)}{\tau - \mathbf{t}} \left\{ e^{-\frac{(\eta - \alpha)^2 + (\xi - \beta)^2}{4(\tau - \mathbf{t})^2}} - e^{-\frac{(\eta - \alpha)^2 + (\xi + \beta)^2}{4(\tau - \mathbf{t})}} \right\} dt d\alpha d\beta$$

Substituting the boundary conditions yields

$$\phi_{2} = -\frac{1}{4\pi} \int_{-\sqrt{Re_{d}}}^{\sqrt{Re_{d}}} \int_{-e/Re_{d}}^{e/Re_{d}} \int_{0}^{\infty} \frac{erf(\frac{\beta}{2/t})}{\tau - t} \left\{ e^{-\frac{(\eta - \alpha)^{2} + (\xi - \beta)^{2}}{4(\tau - t)}} - e^{-\frac{(\eta - \alpha)^{2} + (\xi + \beta)^{2}}{4(\tau - t)}} \right\} dt d\alpha d\beta$$

or

$$\phi_{2} = \frac{1}{4\sqrt{\pi}} \int_{0}^{\sqrt{Re}d} \int_{0}^{\infty} \frac{\operatorname{erf}(\frac{\beta}{2\sqrt{t}})}{\tau - t} \left\{ \operatorname{erf}(\frac{\eta - e\sqrt{Re}d}{2\sqrt{\tau - t}}) - \operatorname{erf}(\frac{\eta + e\sqrt{Re}d}{2\sqrt{\tau - t}}) \right\}$$
$$\left\{ e^{-\frac{(\xi - \beta)^{2}}{4(\tau - t)}} - e^{-\frac{(\xi + \beta)^{2}}{4(\tau - t)}} \right\} dt d\alpha d\beta$$
(A-11)

The final solution for the boundary layer flow over the threedimensional obstacle then reads

$$\frac{U}{U_{\infty}} = \operatorname{erf}(\frac{\xi}{2\sqrt{\tau}}) + \frac{1}{4\sqrt{\pi}} \int_{0}^{\overline{Re}} \int_{0}^{\infty} \frac{\operatorname{erf}(\frac{\beta}{2\sqrt{\tau}})}{\tau - t} \left\{ \operatorname{erf}(\frac{n - e\sqrt{Re}}{2\sqrt{\tau - t}}) - \operatorname{erf}(\frac{n + e/\overline{Re}}{2\sqrt{\tau - t}}) \right\}$$

$$\left\{ e^{-\frac{(\xi - \beta)^{2}}{4(\tau - t)}} - e^{-\frac{(\xi + \beta)^{2}}{4(\tau - t)}} \right\} dt d\alpha d\beta \qquad (A-12)$$

x	Y	\$	5*	0	н	X	Y	8	5*	0	Ħ
(ft)	(1n.)	(in.)	(in.)	(in.)		(ft)	(in.)	(in.)	(in.)	(in.)	
1	0	1.80	268	201	1.33	6	0	4.34	.544	.418	1.30
•	1	1.90	.284	.213	1.33		1	4.35	. 551	.422	1.31
	2	1.75	.260	.193	1.31		2	4.42	.592	.448	1.32
	2.5	1.86	.300	.223	1.35		2 5	6 30	615	454	1.35
	3	1.80	.320	.220	1.45		3	4.59	676	477	1 42
	3.5	1.86	-307	.216	1.42		3.5	4 48	.602	.440	1.37
	4	1.82	.282	.206	1.37		4	4.59	.607	.446	1.36
	5	1.76	.306	.213	1.44		5	4.41	-605	.441	1.37
	6	1.81	.316	.216	1.46		6	4.41	.632	.455	1.39
2	0	2 56	340	268	1 20		0	5 01	622	107	1 20
-	1	2.50	334	260	1.29	°	,	5.05	.033	-407	1 31
	2	2.45	333	258	1.20		2	5.20	.000	518	1 31
	2 5	2 51	356	272	1.29		2 5	5.10	.079	.5.5	1.31
	3	2.42	302	273	1.01		2.5	5.30	.003	.505	1.36
	3.5	2.57	372	274	1 36		3 5	5.10	620	-515	1.34
	4	2.46	338	252	1.34		4	5 22	730	528	1 38
	5	2.61	387	281	1 38	4	5	5 17	713	521	1 37
	6	2.51	.360	.265	1.36		6	5.03	.737	.525	1.40
3	0	3.27	.413	.322	1.28	10	0	5.50	.774	. 573	1.35
	1	3.25	.407	.319	1.28		1	5.70	.805	. 595	1.35
	2	3.33	.409	.319	1.28		2	5.76	.782	- 578	1.35
	2.5	3.38	.440	.335	1.31		2.5	5.57	.765	.571	1.34
	3	3.00	- 451	.324	1.39		3	5.91	.840	.604	1.39
	3.5	3.24	.462	.335	1.38		3.5	5.77	.773	.571	1.35
	4	3.26	.444	.326	1.36	<i>u</i>	4	5.57	.797	.574	1.39
	5	3.22	.485	.341	1.42		5	5.43	.779	.566	1.33
	6	3.20	.465	.333	1.40	3	6	5.24	.825	.580	1.42
4	0	3.80	.451	.357	1.26						
	1	3.72	.452	.357	1.27						
	2	3.76	.457	.359	1.27						
	2.5	3.72	.486	.373	1.30						
	з	3.93	. 533	.330	1.40						
	3.5	3.85	.504	.375	1.34						
	4	3.56	.496	.365	1.36						
	5	3.59	.527	.380	1.39						
	6	3.62	.528	.381	1.39						

Table 5.1a Summary of boundary layer parameters,  $U_{\rm \infty}\text{=}$  20 ft/sec

X (ft)	Y (in.)	\$ (in.)	5* (in.)	θ (in.)	н	X (ft)	¥ (in.)	8 (in.)	5* (in.)	0 (in.)	н
	0	1.71	.240	185	1.29	6	0	5.99	486	378	1.29
	1	1.58	.222	173	1.28	1	1	4.12	. 507	.392	1.20
	2	1.71	248	.193	1.28	1	2	4.08	.534	407	1.31
	2 5	1 70	246	190	1.20		2.5	4.07	546	411	1 33
	3	1 78	207	210	1.41	1	3	4.13	506		1.60
	3 5	1 76	260	105	1 37		3 5	4.10	556	413	1 37
	4	1 73	205	102	1.30		4	4.20	594	.415	1.30
	5	1.80	263	180	1.39	1	5	4.22	618	.420	1 42
	5	1.00	207	207	1.03		6	4.00	613	.404	1.42
	0	1.02	.297	.207	1.45	i.	0	4.00	.015	• 427	1.44
2	0	2.46	.308	.241	1.28	8	0	4.75	. 592	.460	1.29
	1	2.44	.309	.242	1.28		1	4.94	.632	.486	1.30
	2	2.24	.295	.231	1.28		2	4.88	.652	.494	1.32
	2.5	2.33	.329	.252	1.31		2.5	5.02	.646	.493	1.31
	з	2.30	.368	.257	1.43		з	4.99	.672	.492	1.37
	3.5	2.44	.347	.254	1.37		3.5	5.07	.675	.496	1.36
	4	2.37	.330	.242	1.36		4	5.01	.686	.494	1.39
	5	2.38	.364	.257	1.42	1	5	4.97	.702	.498	1.41
	6	2.38	.359	.256	1.40		6	5.18	.701	.503	1.39
3	0	3.15	.371	.292	1.27	10	0	5.19	.740	.546	1.36
	1	2.99	.353	.282	1.25	1. A A A A A A A A A A A A A A A A A A A	1	5.42	.772	.565	1.37
	2	2.97	.371	.291	1.27		2	5.55	.767	.563	1.36
	2.5	3.07	.330	.295	1.29		2.5	5.32	.784	.565	1.39
	з	2.95	.441	.315	1.40		3	5.16	.801	.556	1.44
	3.5	2.98	.413	.302	1.38		3.5	5.20	.766	.541	1.42
	4	2.95	.449	.316	1.42		4	5.46	.748	.539	1.39
	5	2.88	.450	.312	1.44		5	5.30	.772	.541	1.43
	6	2.96	.478	.328	1.46		б	5.08	.761	.529	1.44
4	0	3.83	.415	.334	1.24						
	1	3.63	.418	.336	1.24						
	2	3.62	.416	.331	1.26						
	2.5	3.46	.434	.336	1.29						
	з	3.28	.484	.339	1.43						
	3.5	3.53	.474	.346	1.37						
	4	3.44	.477	.349	1.37						
	5	3.59	.508	.362	1.40						
	6	3.53	.502	.359	1.40						

Table 5.1b Summary of boundary layer parameters,  $\rm U_{\infty}$  = 30 ft/sec

et men municipaniae	X	=3.0 ft $=2.88 \times 10^5$			X	=4.0 ft =3.85x10 <sup>5</sup>			X R	=5.0 ft $x = 4.81 \times 10^5$	
Y (in.)	U <sub>*</sub> (fps)	(lbs/ft <sup>2</sup> )	$c_{f}$	¥ (in.)	U <sub>*</sub> (fps)	(lbs/ft <sup>2</sup> )	$c_{f}$	¥ (in.)	U <sub>*</sub> (fps)	(lbs/ft <sup>2</sup> )	¢f
0.5 1.0 2.0 2.5 2.75	.890 .861 .910 .853 .769	$1.48 \times 10^{-3}$ $1.38 \times 10^{-3}$ $1.54 \times 10^{-3}$ $1.36 \times 10^{-3}$ $1.10 \times 10^{-3}$	$3.97 \times 10^{-3}$ 3.72 \text{10}^{-3} 4.15 \text{10}^{-3} 3.65 \text{10}^{-3} 2.97 \text{10}^{-3}	0.5 1.0 2.0 2.5 2.75	.868 .827 .867 .816 .755	$1.40 \times 10^{-3}$ $1.23 \times 10^{-3}$ $1.40 \times 10^{-3}$ $1.24 \times 10^{-3}$ $1.06 \times 10^{-3}$	$3.78 \times 10^{-3}$ $3.43 \times 10^{-3}$ $3.77 \times 10^{-3}$ $3.34 \times 10^{-3}$ $2.86 \times 10^{-3}$	0.5 1.0 2.0 2.5 2.75	.817 .305 .844 .781 .734	$1.24 \times 10^{-3}$ $1.21 \times 10^{-3}$ $1.33 \times 10^{-3}$ $1.14 \times 10^{-3}$ $1.01 \times 10^{-3}$	3.35x10 <sup>-3</sup> 3.25x10 <sup>-3</sup> 3.57x10 <sup>-3</sup> 3.06x10 <sup>-3</sup> 2.70x10 <sup>-3</sup>
STREET, Street											
	X R	=6.0 ft $x = 5.77 \times 10^{5}$			X R	=8.0 ft =7.69x10 <sup>5</sup>			X	=10.0 ft =9.61x10 <sup>5</sup>	
¥ (in.)	X R U*	=6.0 ft $x = 5.77 \times 10^{5}$ $T_{w}$ (lbs/ft <sup>2</sup> )	¢f	¥ (in.)	X R U <sub>*</sub> (fps)	=8.0 ft x=7.69x10 <sup>5</sup> <sup>T</sup> w (lbs/ft <sup>2</sup> )	¢f	¥ (in.)	X R U* (fps)	=10.0 ft =9.61x10 <sup>5</sup> $T_{w}^{T}$ (lbs/ft <sup>2</sup> )	¢f
Y (in.) 0.5 1.0 2.0	X R U* (fps) .782 .816 .816	=6.0 ft $x = 5.77 \times 10^{5}$ $T_{W}$ (lbs/ft <sup>2</sup> ) 1.14×10 <sup>-3</sup> 1.12×10 <sup>-3</sup> 1.24×10 <sup>-3</sup>	$c_{f}$ 3.07x10 <sup>-3</sup> 3.03x10 <sup>-3</sup> 3.34x10 <sup>-3</sup>	¥ (in.) 0.5 1.0 2.0	X R U* (fps) .725 .735 .766	=8.0 ft $x^{T}$ ,69x10 <sup>5</sup> $T_{W}$ (lbs/ft <sup>2</sup> ) 9.81x10 <sup>-4</sup> 1.01x10 <sup>-3</sup> 1.09x10 <sup>-3</sup>	$c_{f}$ 2.64x10 <sup>-3</sup> 2.71x10 <sup>-3</sup> 2.94x10 <sup>-3</sup>	¥ (in.) 0.5 1.0 2.0	X R (fps) .679 .695 .731	=10.0 ft =9.61x10 <sup>5</sup> $T_w$ (lbs/ft <sup>2</sup> ) 8.60x10 <sup>-4</sup> 9.00x10 <sup>-4</sup> 9.96x10 <sup>-4</sup>	$c_{f}$ 2.31x10 <sup>-3</sup> 2.42x10 <sup>-3</sup> 2.68x10 <sup>-3</sup>

Table 5.2a Preston tube measurement of wall shear stress,  $U_{\infty}$  = 20 ft/sec

Contraction of the	Statement and a statement	the second s	NAMES OF TAXABLE PARTY.	Contraction of the local division	outputs of the lot of	STATES OF THE SALAR STATES	The survey of the second s			And a state of the second	and the state of the second
	X	= 3.0  ft = 4.36x10 <sup>5</sup>			X	x = 4.0  ft $x = 5.31 \times 10^5$			X	= 5.0  ft = 7.26x10 <sup>5</sup>	
¥ (in.)	U <sub>*</sub> ) (fps)	(lbs/ft <sup>2</sup> )	$\mathtt{c}_{\mathtt{f}}$	¥ (in.)	U <sub>*</sub> (fps)	τ <sub>w</sub> (lbs/ft <sup>2</sup> )	$c_{f}$	¥ (in.)	U <sub>*</sub> (fps)	τ <sub>w</sub> (lbs/ft <sup>2</sup> )	$c_{f}$
0.5	1.36	3.43x10 <sup>-3</sup>	4.00x10 <sup>-3</sup>	0.5	1.32	3.21x10-3	3.74x10 <sup>-3</sup>	0.5	1.28	3.05x10 <sup>-3</sup>	3.55x10 <sup>-3</sup>
1.0	1.36	3.39x10 <sup>-3</sup>	3.95x10 <sup>-3</sup>	1.0	1.31	3.17x10 <sup>-3</sup>	3.69x10 <sup>-3</sup>	1.0	1.28	3.01x10 <sup>-3</sup>	3.50x10 <sup>-3</sup>
2.0	1.35	3.36x10 <sup>-3</sup>	3.91x10 <sup>-3</sup>	2.0	1.32	$3.19 \times 10^{-3}$	3.72x10 <sup>-3</sup>	2.0	1.25	2.90x10 <sup>-3</sup>	3.38x10 <sup>-3</sup>
2.5	1.30	3.11x10 <sup>-3</sup>	3.62x10 <sup>-3</sup>	2.5	1.26	2.93x10 <sup>-3</sup>	$3.41 \times 10^{-3}$	2.5	1.23	2.81x10 <sup>-3</sup>	$3.27 \times 10^{-3}$
2.75	1.19	$2.61 \times 10^{-3}$	3.04x10 <sup>-3</sup>	2.75	1.17	2.51x10 <sup>-3</sup>	$2.93 \times 10^{-3}$	2.75	1.13	2.36x10 <sup>-3</sup>	2.75x10 <sup>-3</sup>
decen antenna	x	= 6.0 ft			X	= 8.0 ft			x	=10.0 ft	
6	X R	= 6.0  ft = 8.71x10 <sup>5</sup>			X R	= 8.0 ft =1.12210 <sup>6</sup>			X R	=10.0 ft =1.45x10 <sup>6</sup>	
<u></u> У	х <sup>R</sup> х <sub>v</sub>	= 6.0  ft $= 8.71 \times 10^5$	°,	¥	X R V	= 8.0 ft $x = 1.12210^{6}$	° <sub>f</sub>	Y	х <sub>R</sub> х v.	=10.0 ft =1.45x10 <sup>6</sup> $\tau_{W}$	c <sub>f</sub>
Y (in.)	X R V (fps)	= 6.0 ft = $8.71 \times 10^{5}$ (lbs/ft <sup>2</sup> )	¢ f	¥ (in.)	X R U_* (fps)	= 8.0 ft = $1.12 \times 10^{6}$ (lbs/ft <sup>2</sup> )	¢f	¥ (in.)	X R <sub>x</sub> U <sub>*</sub> (fps)	=10.0 ft =1.45x10 <sup>6</sup> $\tau_{W}^{2}$ (lbs/ft <sup>2</sup> )	¢f
¥ (in.) 0.5	X R U (fps) 1.24	$= 6.0 \text{ ft}$ $= 8.71 \times 10^{5}$ $(1 \text{ bs/ft}^{2})$ $2.85 \times 10^{-3}$	¢ f 3.32x10 <sup>3</sup>	¥ (in.) 0.5	X R U* (fps) 1.18	= 8.0 ft $x^{=1.12\times10^{6}}$ $(1bs/ft^{2})$ 2.55x10 <sup>3</sup>	¢ <sub>f</sub> 2.98x10 <sup>-3</sup>	¥ (in.) 0.5	X R <sub>x</sub> U <sub>*</sub> (fps) 1.10	=10.0 ft =1.45x10 <sup>6</sup> $\tau_{W}$ (lbs/ft <sup>2</sup> ) 2.25x10 <sup>-3</sup>	¢ <sub>f</sub> 2.62x10 <sup>-3</sup>
¥ (in.) 0.5 1.0	X R U (fps) 1.24 1.24	$= 6.0 \text{ ft}$ $= 8.71 \times 10^{5}$ $(1 \text{ bs/ft}^{2})$ $2.85 \times 10^{-3}$ $2.83 \times 10^{-3}$	c <sub>f</sub> 3.32x10 <sup>-3</sup> 3.30x10 <sup>-3</sup>	¥ (in.) 0.5 1.0	X R U <sub>*</sub> (fps) 1.18 1.17	= 8.0 ft = 1.12x10 <sup>6</sup> $T_w$ (lbs/ft <sup>2</sup> ) 2.55x10 <sup>3</sup> 2.52x10 <sup>3</sup>	c <sub>f</sub> 2.98x10 <sup>-3</sup> 2.94x10 <sup>-3</sup>	¥ (in.) 0.5 1.0	X R <sub>x</sub> U <sub>*</sub> (fps) 1.10 1.11	=10.0 ft =1.45x10 <sup>6</sup> $\tau_{W}$ (1bs/ft <sup>2</sup> ) 2.25x10 <sup>-3</sup> 2.29x10 <sup>-3</sup>	¢ <sub>f</sub> 2.62x10 <sup>-3</sup> 2.67x10 <sup>-3</sup>
¥ (in.) 0.5 1.0 2.0	X R U (fps) 1.24 1.24 1.22	= 6.0 ft = $8.71 \times 10^{5}$ (1bs/ft <sup>2</sup> ) 2.35x10 <sup>-3</sup> 2.83x10 <sup>-3</sup> 2.73x10 <sup>-3</sup>	c <sub>f</sub> 3.32x10 <sup>-3</sup> 3.30x10 <sup>-3</sup> 3.18x10 <sup>-3</sup>	¥ (in.) 0.5 1.0 2.0	X R U* (fps) 1.18 1.17 1.15	= 8.0 ft $x = 1.12 \times 10^{6}$ $T_{W}$ (lbs/ft <sup>2</sup> ) 2.55x10 <sup>-3</sup> 2.52x10 <sup>-3</sup> 2.42x10 <sup>-3</sup>	c <sub>f</sub> 2.98x10 <sup>-3</sup> 2.94x10 <sup>-3</sup> 2.82x10 <sup>-3</sup>	¥ (in.) 0.5 1.0 2.0	X R <sub>x</sub> (fps) 1.10 1.11 1.08	=10.0 ft =1.45x10 <sup>6</sup> $T_{W}$ (lbs/ft <sup>2</sup> ) 2.25x10 <sup>-3</sup> 2.29x10 <sup>-3</sup> 2.16x10 <sup>-3</sup>	$c_{f}$ 2.62x10 <sup>-3</sup> 2.67x10 <sup>-3</sup> 2.52x10 <sup>-3</sup>
¥ (in.) 0.5 1.0 2.0 2.5	X R U (fps) 1.24 1.24 1.22 1.19	$= 6.0 \text{ ft}$ $= 8.71 \times 10^{5}$ $(1 \text{ bs/ft}^{2})$ $2.85 \times 10^{-3}$ $2.83 \times 10^{-3}$ $2.73 \times 10^{-3}$ $2.60 \times 10^{-3}$	c <sub>f</sub> 3.32x10 <sup>-3</sup> 3.30x10 <sup>-3</sup> 3.18x10 <sup>-3</sup> 3.02x10 <sup>-3</sup>	¥ (in.) 0.5 1.0 2.0 2.5	X R U* (fps) 1.18 1.17 1.15 1.14	= 8.0 ft $x = 1.12 \times 10^{6}$ $\tau_{W}$ (1bs/ft <sup>2</sup> ) 2.55x10 <sup>-3</sup> 2.52x10 <sup>-3</sup> 2.39x10 <sup>-3</sup>	¢ <sub>f</sub> 2.98x10 <sup>-3</sup> 2.94x10 <sup>-3</sup> 2.82x10 <sup>-3</sup> 2.79x10 <sup>-3</sup>	¥ (in.) 0.5 1.0 2.0 2.5	X R <sub>x</sub> (fps) 1.10 1.11 1.08 1.10	=10.0 ft =1.45x10 <sup>6</sup> $T_W$ (1bs/ft <sup>2</sup> ) 2.25x10 <sup>-3</sup> 2.29x10 <sup>-3</sup> 2.16x10 <sup>-3</sup> 2.21x10 <sup>-3</sup>	<sup>C</sup> f 2.62x10 <sup>-3</sup> 2.67x10 <sup>-3</sup> 2.52x10 <sup>-3</sup> 2.58x10 <sup>-3</sup>

Table 5.2b Preston tube measurement of wall shear stress,  $U_{\infty}$  = 30 ft/sec

X= 1.0 FT	X= 1.0 FT	X= 1.0 FT	x= 1.0 FT	X= 1.0 FT	X= 1.0 FT	X= 1.0 FT	X= 1.0 FT	x= 1.0 FT
Y= 0.0 IN	Y= 1.0 IM	Y= 2.0 IN	r= 2.5 IN	Y= 5.0 IN	Y= 3.5 IN	Y= 4.0 IN	Y= 5.0 IN	Y= 6.0 IN
Z U	Z U	Z U	Z U	Z U	Z U	Z U	2 U	Z U
(IN) (FPS)	(IN) (FPS)	(IN) (FPS:	(IN) (FPS)	(IND (FPS)	UNI IFPSI	(IN) (FPS)	(IN) (FPS)	(112 (FPS)
.040 10.046 .056 10.466 .072 11.074 .091 11.844 .106 12.376 .130 12.94 .161 13.441 .201 13.831 .248 14.184 .305 14.510 .482 15.376 .482 15.376 .612 16.05 .684 16.45 .746 16.80 .808 17.236 .863 17.58 .919 18.02 .989 18.41 1.69 19.17 1.169 19.10 1.269 19.377 1.372 19.58 .1497 19.797 1.372 20.302 2.000 20.365 2.220 20.404	.037 9.38: .050 10.206 .066 10.819 .084 11.425 .107 12.015 .126 12.775 .156 15.317 .208 13.317 .208 13.317 .208 13.514 .596 14.198 .393 14.508 .393 14.508 .527 15.588 .527 15.588 .527 15.588 .527 15.645 .604 16.102 .675 16.453 .761 16.8355 .178 18.355 .178 18.355 .178 18.355 .178 18.355 .149 19.726 .095 20.315 .728 20.767 .2050 20.417 .2051 20.447 .217 20.447	.035 9.688 .645 10.515 .057 11.931 .095 12.567 .121 15.150 .150 12.567 .205 14.079 .265 14.460 .596 15.275 .452 15.589 .524 15.967 .595 16.554 .659 16.672 .755 17.059 .809 17.422 .869 17.421 .966 18.046 1.026 18.046 1.026 18.046 1.026 18.051 1.217 19.097 1.307 19.571 1.217 19.097 1.307 19.571 1.217 19.097 1.307 19.571 1.217 19.097 1.307 19.522 .410 19.628 1.514 19.862 1.641 20.112 1.765 20.322 2.012 20.407 2.216 20.414	.038 9.533 .052 10.325 .071 11.009 .098 11.725 .127 12.350 .155 12.813 .196 13.528 .258 13.758 .268 14.195 .548 14.480 .414 14.769 .414 14.769 .414 14.769 .414 14.769 .414 14.769 .414 14.769 .415 15.828 .632 15.828 .636 16.186 .632 15.828 .636 16.186 .632 15.828 .636 16.186 .632 15.828 .636 16.186 .16.55 19.105 1.553 19.105 1.553 19.379 1.569 19.687 1.963 20.267 1.963 20.267 1.963 20.259 2.084 20.359	.036 7.411 .063 8.176 .069 9.041 .114 9.772 .145 10.485 .171 11.216 .201 11.910 .231 12.474 .266 12.973 .520 15.429 .585 15.900 .445 14.535 .501 14.788 .706 16.554 .772 16.752 .851 17.127 .919 17.605 1.655 19.142 1.565 19.142 1.561 19.944 1.561 19.944 1.561 19.944 1.561 19.944 1.561 19.944 1.561 19.944 1.561 19.944 1.561 20.056 1.811 20.220 1.955 20.522 2.077 20.595	.127 9.614 .167 9.743 .199 10.168 .226 10.725 .264 11.433 .302 12.523 .548 12.914 .405 15.589 .463 14.116 .516 14.602 .575 15.059 .651 15.599 .725 16.037 .792 16.451 .865 16.901 .935 17.528 .002 17.690 1.090 18.110 1.165 18.994 1.251 18.697 1.547 19.766 1.647 19.766 1.647 19.786 1.647 19.786 1.647 19.786 1.647 19.786 1.647 19.786 1.647 19.788 2.134 20.397 2.238 20.414 2.258 20.414	$\begin{array}{c} .129 & 9.962 \\ .161 & :0.253 \\ .199 & 10.814 \\ .214 & 11.449 \\ .255 & 11.983 \\ .265 & 12.620 \\ .296 & 13.150 \\ .341 & 13.610 \\ .390 & 14.004 \\ .443 & 14.785 \\ .551 & 15.218 \\ .600 & 15.531 \\ .551 & 15.262 \\ .715 & 16.254 \\ .805 & 16.746 \\ .872 & 17.965 \\ .945 & 17.581 \\ 1.028 & 17.995 \\ 1.018 & 18.929 \\ 1.78 & 18.656 \\ 1.247 & 18.925 \\ 1.516 & 19.127 \\ 1.403 & 19.579 \\ 1.522 & 19.687 \\ 1$	.128 9.355 .167 9.772 .197 10.253 .224 10.921 .259 11.620 .296 12.346 .334 12.847 .500 14.092 .500 14.095 .500 14.095 .500 14.096 .503 15.075 .701 15.593 .753 16.419 .815 16.445 .889 16.915 .699 16.915 .259 18.871 1.334 19.280 1.399 19.538 1.491 19.799 1.597 19.997 1.698 20.141 1.815 20.279 1.931 20.364 2.655 20.402 2.235 20.402	.129 8.087 .161 8.533 .195 9.179 .227 10.035 .249 10.810 .280 11.566 .324 12.296 .374 12.296 .374 12.296 .541 14.266 .541 14.266 .541 14.266 .546 14.202 .665 15.372 .802 16.457 .869 16.901 .944 17.311 1.009 17.644 1.079 18.046 1.157 18.466 1.228 18.762 1.298 18.999 1.584 19.315 .478 19.560 1.589 19.807 1.711 20.041 1.813 20.562 1.933 20.522 2.084 20.397 2.236 20.397
x= 2.0	FT X= 2.0 FT	x= 2.0 FT	x= 2.0 FT	x= 2.0 FT	x= 2.0 FT	x= 2.0 FT x	= 2.0 FT x=	2.0 FT
Y= 0.0	IN Y= 1.0 IM	x= 2.0 IN	T= 2.5 IN	x= 3.0 IN	r= 3.5 IN	Y= 4.0 IN Y	= 5.0 IN Y=	6.0 IN
C IND	U Z	U Z U	Z U	Z U	Z U	Z U	Z U	U
	(FPS) (IN) (F	PSI (IN) (FPS)	(IN) (FPS)	(IN) (FPS)	(IN) (FPS)	(IN) (FPS)	(IN) (FPS) (	(NJ) (FPS)
.056 077 090 106 1 120 1 140 1 160 1 229 1 279 1 279 1 279 1 537 1 613 1 792 1 1.052 1 1.265 1 1.265 1 1.265 1 1.265 1 1.265 1 1.265 1 1.268 1 2.462 1 2.462 1 2.462 1 2.465 2 3.150 2	8,556, .657, 9, 9,426, .095, 10, 0,387, .108, 11, 1,202, .127, 12, 1,946, .153, 12, 2,657, .184, 15, 5,281, .208, 15, 3,759, .259, 14, 4,192, .328, 14, 4,1526, .409, 15, 5,582, .429, 14, 4,1526, .409, 15, 5,582, .429, 14, 1,526, .429, 15, 5,582, .575, 16, 6,444, .555, 16, 6,446, .553, 16, 6,446, .553, 16, 6,446, 1,067, 17, 6,942, 1,78, 17, 7,229, 1,355, 17, 1,546, 18, 8,180, 1,671, 18, 8,180, 1,671, 18, 1,625, 2,425, 14, 9,414, 2,216, 19, 9,414, 2,216, 19, 1,105, 5, 131, 20, 0,106, 5, 151,	421         .057         0.816           660         .076         9.406           523         .099         10.113           325         .095         10.901           0553         .109         10.113           325         .095         10.911           055         .109         12.202           462         .159         12.844           959         .233         13.821           699         .281         14.212           069         .281         14.212           069         .281         14.421           059         .283         15.962           054         .579         15.621           576         .492         15.296           054         .579         15.621           533         .670         15.992           0504         .944         16.905           0511         .009         17.168           948         1.117         17.497           229         1.226         17.765           043         1.638         18.766           043         1.638         19.766           0451         19.956 </td <td>.059 9.118 .085 9.487 .093 10.267 .104 10.870 .127 11.655 .149 12.266 .177 12.699 .218 12.266 .355 14.255 .402 14.657 .478 15.556 .644 15.743 .720 16.313 .884 16.615 .900 16.933 .884 16.615 .900 16.933 .884 16.615 .900 16.933 .884 16.615 .900 16.933 .044 17.278 1.210 17.597 1.520 17.895 1.445 18.255 1.566 18.522 .2.619 19.623 .2.196 19.670 1.211 19.453 .2.196 19.670 .2.345 19.856 .2.475 19.949 2.647 20.020 .2.345 19.856 .2.150 20.186 .2.150 20.186 .150 20.186 .150 20.186 .150 20.186 .150 20.186</td> <td>056         4,647           077         5,289           095         6,158           107         7,296           121         8,190           146         9,411           175         10,396           201         11,308           2255         12,009           276         12,589           529         13,111           583         13,662           447         14,182           516         14,705           5697         15,211           589         15,639           -765         16,039           -862         16,443           -967         16,218           1.964         17,527           1.303         17,852           1.400         18,119           1.522         18,427           1.647         18,716           1.7527         1,203           1.642         18,427           1.642         18,427           1.643         19,675           2.268         19,800           2.435         19,9250           2.268         19,800           2.435</td> <td><math display="block">\begin{array}{c} .113 &amp; 9.371 \\ .155 &amp; 9.646 \\ .156 &amp; 10.202 \\ .188 &amp; 10.905 \\ .225 &amp; 11.500 \\ .269 &amp; 12.167 \\ .318 &amp; 12.725 \\ .393 &amp; 13.334 \\ .444 &amp; 13.787 \\ .515 &amp; 14.262 \\ .594 &amp; 14.750 \\ .682 &amp; 15.256 \\ .775 &amp; 15.715 \\ .969 &amp; 16.145 \\ .945 &amp; 16.841 \\ .134 &amp; 17.195 \\ .245 &amp; 17.921 \\ .468 &amp; 18.263 \\ .945 &amp; 17.921 \\ .468 &amp; 18.583 \\ .755 &amp; 18.875 \\ .965 &amp; 19.148 \\ .161 &amp; 18.583 \\ .755 &amp; 18.875 \\ .965 &amp; 19.148 \\ .161 &amp; 18.583 \\ .1755 &amp; 18.875 \\ .965 &amp; 19.148 \\ .161 &amp; 19.410 \\ .124 &amp; 19.575 \\ .260 &amp; 19.697 \\ .2428 &amp; 19.834 \\ .2575 &amp; 19.915 \\ .2.925 &amp; 20.082 \\ .131 &amp; 20.112 \end{array}</math></td> <td><math display="block">\begin{array}{cccccccccccccccccccccccccccccccccccc</math></td> <td>134         9.139           162         9.152           .196         10.115           .225         10.640           .225         10.640           .201         11.210           .502         11.745           .503         12.108           .543         12.108           .543         12.108           .543         12.108           .543         12.108           .543         12.108           .543         12.108           .543         12.108           .543         12.108           .543         12.108           .452         13.612           .445         13.612           .559         14.026           .559         14.026           .559         14.026           .512         15.251           .014         16.578           .021         16.578           .021         16.578           .032         19.575           .503         19.506           .503         19.506           .503         19.506           .503         20.101           .132</td> <td>133 8.962 159 9.525 175 9.895 175 9.895 175 9.895 175 9.895 175 9.895 175 9.895 175 9.895 175 9.895 175 9.895 175 12.384 17.355 12.384 10.154 14.335 12.384 14.335 15.706 1</td>	.059 9.118 .085 9.487 .093 10.267 .104 10.870 .127 11.655 .149 12.266 .177 12.699 .218 12.266 .355 14.255 .402 14.657 .478 15.556 .644 15.743 .720 16.313 .884 16.615 .900 16.933 .884 16.615 .900 16.933 .884 16.615 .900 16.933 .884 16.615 .900 16.933 .044 17.278 1.210 17.597 1.520 17.895 1.445 18.255 1.566 18.522 .2.619 19.623 .2.196 19.670 1.211 19.453 .2.196 19.670 .2.345 19.856 .2.475 19.949 2.647 20.020 .2.345 19.856 .2.150 20.186 .2.150 20.186 .150 20.186 .150 20.186 .150 20.186 .150 20.186	056         4,647           077         5,289           095         6,158           107         7,296           121         8,190           146         9,411           175         10,396           201         11,308           2255         12,009           276         12,589           529         13,111           583         13,662           447         14,182           516         14,705           5697         15,211           589         15,639           -765         16,039           -862         16,443           -967         16,218           1.964         17,527           1.303         17,852           1.400         18,119           1.522         18,427           1.647         18,716           1.7527         1,203           1.642         18,427           1.642         18,427           1.643         19,675           2.268         19,800           2.435         19,9250           2.268         19,800           2.435	$\begin{array}{c} .113 & 9.371 \\ .155 & 9.646 \\ .156 & 10.202 \\ .188 & 10.905 \\ .225 & 11.500 \\ .269 & 12.167 \\ .318 & 12.725 \\ .393 & 13.334 \\ .444 & 13.787 \\ .515 & 14.262 \\ .594 & 14.750 \\ .682 & 15.256 \\ .775 & 15.715 \\ .969 & 16.145 \\ .945 & 16.841 \\ .134 & 17.195 \\ .245 & 17.921 \\ .468 & 18.263 \\ .945 & 17.921 \\ .468 & 18.583 \\ .755 & 18.875 \\ .965 & 19.148 \\ .161 & 18.583 \\ .755 & 18.875 \\ .965 & 19.148 \\ .161 & 18.583 \\ .1755 & 18.875 \\ .965 & 19.148 \\ .161 & 19.410 \\ .124 & 19.575 \\ .260 & 19.697 \\ .2428 & 19.834 \\ .2575 & 19.915 \\ .2.925 & 20.082 \\ .131 & 20.112 \\ .131 & 20.112 \\ .131 & 20.112 \\ .131 & 20.112 \\ .131 & 20.112 \end{array}$	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	134         9.139           162         9.152           .196         10.115           .225         10.640           .225         10.640           .201         11.210           .502         11.745           .503         12.108           .543         12.108           .543         12.108           .543         12.108           .543         12.108           .543         12.108           .543         12.108           .543         12.108           .543         12.108           .543         12.108           .452         13.612           .445         13.612           .559         14.026           .559         14.026           .559         14.026           .512         15.251           .014         16.578           .021         16.578           .021         16.578           .032         19.575           .503         19.506           .503         19.506           .503         19.506           .503         20.101           .132	133 8.962 159 9.525 175 9.895 175 9.895 175 9.895 175 9.895 175 9.895 175 9.895 175 9.895 175 9.895 175 9.895 175 12.384 17.355 12.384 10.154 14.335 12.384 14.335 15.706 1

Table 5.3a Mean velocity data,  $\rm U_{\infty}$  = 20 ft/sec

X= 3.0 FT	X= 3.0 FT	X= 3.0 FT	X= 3.0 FT	x= 3.0 FT	X= 3.0 FT	X= 5.0 FT	x= 5.0 FT	x= 3.0 FT
Y= 0.0 IN	Y= 1.0 IN	Y= 2.0 IN	Y= 2.5 IN	Y= 3.0 IN	Y= 3.5 IN	Y= 4.0 IN	Y= 5.0 IN	t= 6.0 IN -
Z U	Z U	Z U	Z U	Z U	Z U	Z U	Z U	Z U
(IN) (FPS)	(IN) (FPS)	(IN) (FPS)	(IN) (FPS)	(IND (FPS)	(IN) (FPS)	([N) (FPS)	(IN) (FPS)	(IN) (FPS)
.055 8.551 .060 8.740 .069 9.345 .083 10.066 .097 10.733 .120 11.466 .143 12.147 .168 12.810 .203 15.306 .248 15.765 .266 14.135 .556 14.533 .452 14.965 .504 15.501 .562 15.627 .680 15.971 .582 15.627 .880 15.971 .108 17.090 1.274 16.778 1.08 17.729 1.458 17.729 1.458 17.729 1.458 17.729 2.265 19.255 .275 19.452 2.715 19.452 2.735 19.452 2.735 19.452 2.356 19.571 2.529 19.568 .3409 19.889 3.629 19.958 3.671 20.008	.038 9.227 .068 9.065 .096 10.773 .111 11.660 .140 12.430 .178 13.026 .230 13.582 .303 14.138 .578 14.578 .469 15.031 .574 15.462 .701 15.899 .850 16.344 .010 16.818 1.490 17.162 1.490 17.857 1.651 18.148 1.809 18.416 1.993 18.673 2.162 18.906 .2348 19.125 2.532 19.311 2.705 19.478 2.875 19.619 5.074 19.745 5.299 19.841 5.469 19.915 5.675 19.980 5.909 20.018 4.187 20.025 4.187 20.025	.040 $9.929$ .053 $9.497$ .073 10.185 .098 10.918 .110 11.700 .141 12.469 .187 15.079 .248 15.654 .529 14.222 .419 14.754 .529 14.222 .419 14.754 .529 15.656 .739 16.048 .850 16.425 .974 16.770 1.102 17.104 1.251 17.426 1.509 17.751 1.556 19.055 1.743 18.567 1.745 18.567 1.745 18.567 1.745 18.567 1.745 18.567 1.745 18.567 1.959 19.695 5.224 19.852 5.568 19.400 5.040 19.755 5.224 19.852 5.667 20.017 4.196 20.077 4.196 20.077	.040 8.531 .056 9.087 .069 9.081 .069 10.281 .115 11.022 .146 11.745 .180 12.318 .215 12.799 .271 15.256 .343 15.714 .422 14.178 .517 14.660 .613 15.073 .759 15.600 .902 16.089 1.052 16.581 1.216 17.054 1.206 18.754 1.400 17.558 1.587 17.995 1.773 18.406 1.986 18.754 1.906 19.267 2.686 19.463 2.991 19.624 3.205 19.733 5.402 19.827 2.555 19.915 5.906 19.982 4.184 20.020 4.184 20.020	.042 6.060 .065 6.945 .085 7.854 .099 8.735 .129 9.601 .167 10.564 .205 11.366 .250 12.057 .306 12.646 .372 13.251 .447 13.787 .527 14.312 .629 14.914 .736 15.459 .874 16.027 1.016 16.512 1.79 16.909 1.334 17.591 1.490 17.817 1.671 18.206 1.969 18.585 2.259 19.106 2.465 19.532 2.869 19.718 5.075 19.901 3.535 19.901 3.555 19.901 3.655 19.901 3.655 19.901 3.655 19.901 3.655 19.907 4.190 19.970 4.190 19.970	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	.157 8.598 .169 9.962 .195 9.651 .226 10.442 .271 11.210 .315 11.870 .536 12.445 .534 13.564 .552 15.969 .659 14.595 .711 14.799 .805 15.270 .888 15.649 .977 16.012 1.103 16.489 1.557 17.328 1.557 17.328 1.557 18.227 1.652 18.573 1.557 18.227 1.652 18.573 1.557 18.227 1.852 18.573 2.224 18.658 2.212 19.118 2.379 19.306 2.572 19.729 3.201 19.855 3.490 19.948 2.766 19.610 2.972 19.729 3.200 20.003 3.922 20.041 4.200 20.070	.153 6.772 .156 7.511 .180 8.271 .210 9.061 .244 9.785 .297 10.557 .356 11.296 .396 11.910 .462 12.472 .538 15.119 .619 15.233 .973 15.724 1.090 16.207 .184 16.554 1.578 17.761 1.716 18.056 2.069 19.774 2.422 19.534 2.492 19.534 2.492 19.534 2.492 19.534 2.492 19.534 2.494 19.905 2.206 19.827 3.419 19.901 3.206 19.827 3.419 19.0015 2.0052 4.199 20.032 4.199 20.032	$\begin{array}{cccccccccccccccccccccccccccccccccccc$
X- 4.0 FT	X= 4.0 FT	x= 4.0 FT	x= 4.0 FT	x= 4.0 FT	x= 4.0 FT	x= 4.0 FT	X= 4.0 FT	X- 4.0 FT
Y- 0.0 IN	Y= 1.0 IN	y= 2.0 IN	r= 2.5 IN	r= 3.0 IN	Y= 3.5 IN	T= 4.0 IN	Y= 5.0 IN	Y- 6.0 [1]
Z U	Z U	Z U	Z U	Z U	Z U	Z U	Z U	Z U
(IN) (FPS	(IN) (FPS	I (IN) (FPS)	(INI (FPS)	(IN) (FPS)	(IN) (FPS)	([hi) (FPS)	(1N) (FPS)	(1HJ (FPS)
.029 8.62 .044 9.02 .055 9.49 .064 10.07 .078 10.66 .094 11.28 .107 11.87 .131 12.53 .157 13.09 .192 13.52 .235 13.90 .303 14.95 .368 14.76 .553 15.54 .659 15.88 .782 16.27 .892 16.56 1.018 16.85 1.168 17.17 1.328 17.46 1.483 17.72 1.667 18.02 1.958 18.43 2.160 18.69 2.415 19.97 2.689 19.24 2.938 19.44 5.165 19.63 5.411 19.77 2.662 19.94 5.165 19.63 5.954 20.09 4.565 20.14 4.825 20.15	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	.042 9.416 .057 9.055 .070 9.755 .086 10.340 .105 10.955 .152 11.686 .135 12.176 .227 12.776 .238 13.265 .375 13.775 .463 14.265 .556 14.727 .657 15.240 .795 15.644 .898 16.012 1.029 16.412 1.029 16.412 1.029 16.412 2.557 17.166 .477 17.166 .477 17.166 .477 17.166 .477 17.166 .477 17.851 1.845 18.202 2.079 19.547 1.552 19.661 2.569 19.812 2.551 19.951 4.204 20.021 4.515 20.081 4.815 20.081	.036         2.492           .051         4.056           .065         5.416           .095         6.826           .110         8.012           .130         9.991           .5         .150         9.905           .150         9.906           .150         9.901           .208         12.206           .415         13.491           .208         12.206           .415         13.491           .415         15.461           .415         15.461           .415         15.461           .416         .416           .710         15.111           .845         15.661           .1010         16.211           .527         17.967           .1527         17.967           .2567         19.272           .2701         18.682           .2567         19.927           .2567         19.927           .2567         19.927           .2567         19.927           .2567         19.927           .2567         19.926           .2567         19.927      .	124         9.759           154         9.325           199         10.665           231         11.319           3.403         12.563           403         12.563           5.58         1.973           403         12.563           5.58         1.975           6.49         14.300           6.72         15.646           6.105         16.759           7.58         1.975           6.49         14.300           6.77         15.646           6.1065         16.754           7.755         1.975           9.158         17.944           1.944         18.281           2.139         19.552           2.139         19.552           2.139         19.655           3.544         19.655           3.544         19.655           3.544         19.655           3.544         19.655           3.642         19.655           3.642         19.655           3.642         19.625           4.828         20.094           4.828         20.094           4	$\begin{array}{c} .131 & 7.175\\ .147 & 8.076\\ .164 & 8.654\\ .189 & 9.692\\ .216 & 10.503\\ .265 & 11.717\\ .319 & 11.782\\ .376 & 12.333\\ .425 & 12.892\\ .531 & 13.499\\ .617 & 14.005\\ .707 & 14.533\\ .821 & 15.096\\ .928 & 15.571\\ .041 & 16.016\\ .190 & 16.544\\ .355 & 17.66\\ .928 & 17.316\\ .1695 & 17.66\\ .928 & 17.316\\ .1695 & 17.66\\ .928 & 17.316\\ .1695 & 17.66\\ .271 & 19.022\\ .212 & 18.894\\ .541 & 19.168\\ .317 & 18.322\\ .212 & 18.894\\ .541 & 19.168\\ .193 & 27.20\\ .2312 & 18.894\\ .541 & 19.56\\ .425 & 19.816\\ .425 & 20.076\\ .825 & 20.076\\ .825 & 20.076\\ .825 & 20.076\\ .825 & 20.076\\ \end{array}$	$\begin{array}{c} .129 & 8.552\\ .166 & 9.152\\ .207 & 9.751\\ .256 & 10.394\\ .314 & 11.091\\ .380 & 11.796\\ .462 & 12.466\\ .544 & 15.077\\ .634 & 15.665\\ .757 & 14.264\\ .851 & 14.865\\ .961 & 15.364\\ .961 & 15.364\\ .961 & 15.364\\ .205 & 16.322\\ .205 & 16.322\\ .205 & 16.322\\ .205 & 16.322\\ .205 & 16.322\\ .205 & 16.322\\ .205 & 16.322\\ .205 & 16.322\\ .205 & 16.322\\ .205 & 16.322\\ .205 & 16.322\\ .205 & 16.322\\ .205 & 16.322\\ .205 & 16.322\\ .205 & 15.362\\ .205 & 16.322\\ .205 & 10.352\\ .2$	134 7.953 167 8.541 195 9.178 253 9.909 279 10.601 255 11.297 363 11.297 363 11.295 4.42 12.409 511 12.954 552 13.547 5687 14.195 1.961 12.954 3.92 15.601 1.116 16.064 1.241 16.484 1.383 16.912 2.155 18.542 2.405 18.849 2.215 18.542 2.405 18.849 5.241 19.155 2.834 19.555 5.241 19.155 2.834 19.555 5.241 19.701 1.5527 19.846 5.3241 19.759 3.452 20.352 3.4829 20.075 3.4829 20.075

Table 5.3a Mean velocity data,  $U_{\infty} = 20$  ft/sec -- Continued

Tal	ole !	5.3a x= 6.1	Me FT	an x= 6.0 Y= 2.0	yelo FT	city ** 6.0	r da ri	ta, x= 6.0 x= 5.0	U <sub>00</sub> =	20 ** 6.0 ** 5.5	ft/	sec x= 6.0 y= 4.0	( FT 111	Cont x= 6.0 r= 5.0	inue M	d *= 6.3 *= 6.0	F T 194
(ÎN)	U (FPS)		U (FPS)	Cho	U (FPS)	2 (110	U (FPS)	, Z (110	U (FPS)	The	U (FPS)	z (Ita	U (FPS)	z (III)	U (FPS)	נווט	U (FPS)
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x= 8.0 Y= 0.0	FT	x= 0.0 Y= 1.0	FT IN	x= 8.0 Y= 2.0	FT IN	x= 8.0 Y= 2.5	FT 111	x= 8.0 Y= 5.0	151 151	x= 8.0 Y= 5.5	) FT 5 [N	X= 8.0 Y= 4.0	) FT ) IN	x= 8.0 T= 5.0	) FT ) 111	x= 9.1 Y= 6.1	0 FT 0 [1]
Z (1N)	U (FPS)	Z (1NJ	(FPS)		(FPS)	Z (IN)	(FPS)	2 (1)10	U (FPS)	2 (110	U (FPS)	Z ( IND	(FPS)	Z (IN)	U (FPS)	Z (110	U (FPS)
.029 .039 .044 .056 .069 .092 .117 .150 .196 .530 .540 .540 .540 .540 .540 .540 .540 .54	8.921 9.166 9.426 9.426 9.426 10.445 10.445 11.265 11.665 12.103 12.461 12.465 13.166 13.166 13.166 13.166 13.166 13.562 14.44 14.918 15.552 15.562 16.592 16.592 17.596 16.592 17.592 17.592 18.581 18.581	022 024 037 048 067 122 126 260 260 260 260 260 260 260 260 271 846 982 1.282 1.282 1.282 1.282 1.282 1.282 1.282 1.282 1.282 1.282 1.282 1.285 2.172 2.597	7,998 8,726 9,876 9,891 10,441 11,005 11,554 12,142 12,608 14,076 14,545 14,545 14,545 14,545 14,545 15,658 14,076 14,545 15,658 15,614 16,614 16,615 17,204 17,209 18,355 18,078 18,566 18,579 18,579	.025 .030 .054 .072 .108 .150 .190 .251 .414 .501 .587 .689 .825 .587 .587 .587 .587 .1.558 .2157 .1.558 .2.265 .2.265 .2.2652 .2.662	8,491 9,120 9,644 10,104 10,517 11,010 11,495 11,495 12,281 12,694 13,145 13,563 13,945 14,549 14,807 15,255 16,416 16,595 16,416 17,121 17,417 17,455 17,939 18,175 18,175 18,392	.052 .049 .065 .100 .158 .107 .245 .590 .475 .574 .781 .245 .5674 .781 .205 1.206 1.276 1.206 1.574 1.642 1.804 1.974 2.141 1.642 2.141 2.954 2.814	9.026 9.0516 10.026 10.485 11.766 12.227 12.627 12.627 12.627 12.627 12.627 12.627 12.627 12.627 12.627 12.627 15.524 15.524 15.526 16.536 17.289 15.536 17.289 17.299 17.	029 040 069 084 105 200 505 609 128 505 505 505 505 505 505 505 505 505 50	5.722 6.633 8.596 9.213 9.737 10.336 10.795 11.307 12.522 13.556 14.056 14.559 15.510 15.510 15.510 15.539 17.2452 17.572 17.572 17.572 17.877 18.170	128 150 164 192 229 255 554 416 555 6554 6554 6554 6554 6554	8.491 8.915 9.510 9.776 10.762 11.235 11.235 11.235 12.111 12.578 13.451 13.860 14.290 14.624 14.991 14.9420 15.814 16.538 14.920 15.814 16.538 16.538 16.538 16.538 17.569 1	.128 .148 .164 .185 .211 .244 .277 .527 .412 .567 .653 .653 .756 .653 .756 .823 .904 1.025 1.139 1.256 1.590 1.590 1.692 1.857 2.025 2.198	7.557 7.653 8.562 8.785 9.279 9.786 10.854 11.307 12.748 15.615 15.982 14.800 15.534 15.997 16.6614 16.946 17.27572 17.572	.128 .1454 .154 .174 .233 .271 .361 .361 .417 .539 .690 .777 .882 1.022 1.128 1.278 1.422 1.579 1.727 1.899 2.040 2.451	7.727 8.115 8.396 8.830 9.269 9.791 10.255 10.255 11.210 11.624 12.056 12.488 12.952 13.437 15.866 14.320 14.772 15.182 15.622 15.622 16.016 16.564 16.564 17.010 17.327 17.609 17.923	.120 .133 .146 .164 .209 .240 .270 .516 .501 .501 .501 .516 .551 .654 .501 .742 .825 .931 1.035 1.294 1.427 1.566 1.717 1.896 2.054 2.254	6.554 6.974 7.415 7.950 8.557 8.921 9.834 10.350 10.771 11.361 11.361 11.351 11.351 12.442 12.961 13.545 15.495 15.495 15.482 16.594 16.594 16.969 17.522 17.635

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Table 5.3a Mean velocity data,  $U_m = 20$  ft/sec -- Continued

X= 1.0 FT Y= 1.0 IN X= 1.0 FT Y= 0.0 IN x= 1.0 FT Y= 2.0 IN X= 1.0 FT Y= 3.0 IN X= 1.0 FT Y= 4.0 IN X= 1.0 FT X= 1.0 FT X= 1.0 FT Y= 3.5 IN X= 1.0 FT Y= 2.5 1N T= 5.0 111 T= 6.0 111 (IN) (FPS) (IN) (FPS) (IN) (FPS) (IN) (FPS) (IN) (FPS) (IN) (FPS) (TH) (FPS) (IN) (FPS) (IN) (FPS) .031 16.429 .048 17.075 .065 17.981 .034 16.534 .050 17.356 .075 18.145 .095 19.045 .120 13.715 .152 14.427 .184 15.158 .034 12.525 .059 13.451 .079 14.215 .099 15.120 .031 16.365 .052 17.024 .128 14.916 .128 15.202 .128 13.259 .128 14.916 .163 15.476 .198 16.106 .223 16.747 .244 17.785 .055 17.236 .159 13.962 .162 15.806 .191 16.365 .060 18.281 .084 19.145 .107 19.798 .131 20.408 .157 20.917 .095 19.005 .084 18.944 .227 15.963 .208 16.189 .125 16.047 17.906 .232 17.131 .118 20.015 .103 19.619 .115 19.751 .248 .258 17.103 .127 20.431 .152 20.515 .272 18.741 .284 18.721 .261 18.081 149 20.738 .290 18.028 .186 21.387 .184 18,044 .214 18,782 .255 19,716 .297 20,594 .595 21,221 .402 21,777 .469 22,452 .527 23,057 .549 25,764 .662 24,493 .732 25,122 .805 25,679 .877 26,221 .977 26,221 .977 26,221 .040 27,874 1,21 27,849 1,244 28,474 1,329 28,904 .184 18.044 .329 18.818 .229 21.904 .281 22.435 .531 22.824 .591 23.288 .194 21.640 .199 21.445 .232 21.596 .559 20.295 .558 20.091 .351 19.924 .369 19.570 .234 22.182 .275 22.706 .315 23.144 .421 20.417 .471 21.090 .527 21.782 .285 22.121 .400 20.758 .242 21.845 .299 22.571 .551 22.849 .404 23.276 .467 25.716 .554 24.171 .354 22.630 .395 25.078 .436 23.501 .442 21.956 .449 21.456 .430 21.632 .495 22.630 .501 22.115 .490 22.503 .315 23.144 .365 23.627 .427 24.088 .492 24.587 .562 25.027 .623 25.519 .695 25.994 .776 26.540 .861 27.139 .941 27.635 1.024 28.088 1.126 28.701 1.230 29.177 1.331 29.545 .456 23.828 .585 22.494 .546 23.194 .602 23.812 .670 24.403 .746 25.091 .830 25.709 .912 26.279 1.011 26.930 1.112 27.579 .554 22.769 .616 23.505 .678 24.203 .757 24.819 .818 25.590 .896 26.308 .980 27.005 1.062 27.593 .554 25.284 .615 25.976 .695 24.825 .778 25.590 .858 26.334 .942 27.040 1.020 27.590 .646 23.239 .715 23.944 .776 24.552 .840 25.187 .917 25.898 .999 26.556 .519 24.325 .508 24.020 .534 24.171 .609 24.668 .684 25.145 .769 25.668 .848 26.166 .940 26.709 1.019 27.304 1.105 27.873 .568 24.497 .628 24.962 .692 25.402 .773 25.961 .857 26.508 .938 27.107 .661 25.406 .804 26.435 1.082 27.227 1.163 27.773 1.245 28.241 .969 27.426 1,122 28.139 1.047 27.814 1.164 28.204 1.266 28.754 1.369 29.193 1.491 29.581 1.215 28.180 1.221 28.584 1.326 28.986 1.427 29.321 1.021 27.652 1.216 28.561 1.184 28.546 1.228 28.691 1.420 29.298 1.323 28.644 1.426 29.105 1.529 29.454 1.650 29.715 1.746 29.928 1.529 29.642 1.634 29.874 1.752 30.059 1.435 29.464 1.555 29.842 1.677 50.091 1.622 29.915 1.622 29.915 1.725 30.115 1.827 30.255 1.599 29.288 1.611 29.851 1.331 29.545 1.392 29.255 1.329 28.904 1.494 29.591 1.617 29.896 1.740 50.125 1.429 29.275 1.553 29.623 1.677 29.683 1.799 30.075 1.946 30.224 2.130 30.355 1.455 29.922 1.978 30.240 1.624 29.909 1.558 30.119 1.838 30.198 1.740 50.125 1.863 50.318 1.989 50.356 2.133 50.382 1.827 50.255 1.927 30.353 2.031 50.426 2.152 50.475 2.152 50.475 2.152 50.475 1.748 30.141 1.892 30.271 2.041 30.341 2.169 30.382 1.640 30.236 1.744 30.300 1.852 30.341 1.971 30.416 1.800 50.240 1.957 30.309 1.873 30.113 2.098 30.378 2.079 30.394 2.180 30.451 2.180 30.451 2.180 30.451 2.180 30.451 2.088 30.391 2.166 30.429 2.189 50.391 2.189 50.391 2.189 50.391 2.189 50.391 2.104 30.300 2.198 30.344 2.198 30.344 2.169 30.382 2.097 30.466 2.135 50.582 2.166 50.429 2,150 30,553 X= 2.0 FT Y= 1.0 IN x= 2.0 FT Y= 3.0 IN X= 2.0 FT Y= 3.5 IN X= 2.0 FT Y= 6.0 IN X= 2.0 FT Y= 2.0 IN T= 2.5 IN Y= 4.0 IN Y= 0.0 IN Y= 5.0 IN Z U Z U (IN) (FPS) Z U (IN) (FPS) (IN) (FPS) (THE IFPS) (IN) IFPSI (IN) (FPS) (IN) (FPS) (IN) (FPS) .125 14.513 .155 15.135 .103 15.721 .125 13.787 .147 14.447 .168 15.148 .125 12.098 .151 13.137 .176 14.087 .202 15.028 .030 15.830 .050 15.974 .052 15.066 ,055 10.835 .125 12.567 .030 14.876 .055 16.414 .075 16.961 .089 17.627 .154 13.317 .179 14.074 .202 14.946 .057 15.360 .055 16.403 .054 15.582 .064 11.450 .105 15.721 208 16.379 245 17.145 282 17.911 325 18.777 .416 20.179 .464 20.792 .520 21.467 .096 16.854 .115 17.535 .135 18.284 .157 19.105 .184 19.745 .119 12.940 .141 13.952 .157 14.914 .092 16.356 .096 17.884 .189 15,879 .202 15.028 .231 15.836 .263 16.593 .294 17.278 .322 17.826 .365 18.537 .233 15.920 .271 16.922 .308 17.847 .113 18.751 .108 18.465 .216 16.707 .122 18.001 157 14.914 175 16,139 199 17,190 222 17,905 257 18,802 296 19,663 140 19.970 169 20.609 202 21.50 204 21.687 505 22.306 547 22.716 400 25.147 522 24.003 565 24.375 5666 24.837 748 25.269 643 25.687 5666 24.837 748 25.269 643 25.682 1.035 26.552 1.055 26.552 1.552 7.505 1.527 28.460 .275 18.048 .318 18.955 .557 19.495 .134 19.043 .176 20.558 .308 17.847 .342 18.496 .306 19.213 .427 19.884 .472 20.580 215 21.150 261 21.756 519 22.258 560 22.720 443 25.180 525 23.716 607 24.285 687 24.798 770 25.259 872 25.769 975 26.651 1.199 27.090 1.321 27.519 1.446 27.961 1.691 28.697 1.691 28.697 1.691 28.697 1.979 20.342 .184 19.745 .216 20.529 .258 20.901 .510 21.480 .415 22.412 .475 22.885 .554 25.283 .627 25.880 .709 24.392 .792 24.871 .147 19.836 .175 20.608 .204 21.226 .234 21.749 .279 22.310 .400 20.150 .406 19.178 .299 19.663 .340 20.576 .395 21.119 .445 21.740 .497 22.310 .557 22.891 .618 25.482 .687 25.999 .462 21.046 .523 21.827 .585 22.480 .649 23.127 .725 23.824 .792 24.470 .867 25.005 .958 25.479 .585 22.186 .645 22.899 .708 25.445 .771 25.979 .496 20.446 .525 21.320 .329 22.812 .377 23.197 .440 23.611 .592 21,600 .631 22.527 .631 22.527 .687 23.086 .747 23.627 .810 24.137 .876 24.613 .641 22.221 .842 24.520 .925 25.103 1.007 25.599 .695 22.845 .758 23.479 .842 24.236 .501 23.999 .565 24.902 .767 25.367 .869 25.872 .972 26.310 1.075 26.676 ,792 24.971 .871 25.055 .974 25.045 1.076 26.567 1.200 26.871 1.322 27.520 1.445 27.736 1.570 28.148 1.712 28.553 1.955 29.928 1.998 29.264 1.955 29.928 1.998 29.264 .687 25.999 .766 24.559 .853 25.122 .955 25.717 1.036 26.139 1.139 26.697 1.241 27.171 1.343 27.563 1.037 25.599 1.091 26.066 1.172 26.476 1.275 26.967 1.378 27.446 1.482 27.872 1.608 28.529 1.753 28.697 1.855 29.035 .842 24.236 .922 24.896 1.005 25.520 1.111 26.165 1.214 26.715 1.318 27.209 1.441 27.718 1.565 28.191 1.020 26.008 1.103 26.533 1.208 27.083 .932 25.027 .998 25.476 1.068 25.898 1.152 26.346 1.236 26.793 1.340 27.307 1.440 27.749 1.511 27.567 1,511 27,567 1,414 27,978 1,517 28,336 1,641 28,740 1,766 29,062 1,891 29,507 2,034 29,527 2,180 29,732 2,525 29,898 2,409 30,022 2,675 30,123 2,865 30,123 1.199 27.122 1.444 27.885 1.586 28.262 1.751 28.667 1.872 28.961 2.058 29.535 1.979 29.342 2.142 29.575 2.305 29.754 1.668 28.842 1.815 29.180 1.957 29.450 1.466 28.005 1.588 28.420 1.735 28.849 1.855 29.036 2.001 29.365 2.145 29.604 1.689 28.620 1.835 29.019 1.979 29.516 1.566 28.229 1.690 28.617 1.834 28.980 1.957 29.450 2.079 29.674 2.225 29.685 2.366 30.000 2.529 50.072 2.712 50.155 2.918 50.185 5.122 50.202 8.122 50.202 2.145 29.604 2.309 29.777 2.475 29.914 2.620 50.054 2.826 30.107 2.986 50.167 5.135 50.189 5.135 50.189 1,979 29,516 2,125 29,578 2,269 29,761 2,454 29,935 2,619 50,057 2,789 30,117 2,953 30,161 3,135 30,167 3,135 30,167 2.145 29.559 2.307 29.786 2.477 29.939 2.669 30.031 2.819 30.082 2.305 29.154 2.489 29.955 2.672 30.057 2.856 30.117 3.024 30.167 1.875 29.232 2.021 29.525 2.185 29.767 1.962 29.297 2.102 29.568 2.036 29.333 2.220 29.617 2.405 29.822 2.606 29.993 2.773 30.069 2.948 30.145 3.122 30.173 2.102 29.568 2.281 29.806 2.441 29.958 2.587 30.069 2.765 30.136 2.954 30.167 2.185 29.767 2.346 29.949 2.511 30.055 2.715 30.151 2.918 30.185 3.122 30.202 2.587 2.765 2.954 5.121 5.121 5.121 3.121 30.189 30.189 30.189 2.863 30.192 2.996 30.101 5.122 50.110 5.122 50.110 5.046 30.218

3,155 30,189

5, 127 50, 250

5.157

Table 5.3b Mean velocity data, U = 30 ft/sec

30.175 5 121 50,189 3.122 50.202

X- 3.0 FT	X= 5.0 FT	X= 3.0 FT	X= 3.0 FT	X= 3.0 FT	X= 3.0 FT	X= 3.0 FT	X= 3.0 FT	x= 3.0 FT
Y- 0.0 IN	Y= 1.0 IN	Y= 2.0 IN	Y= 2.5 IN	Y= 5.0 IN	Y= 3.5 IN	Y= 4.0 IN	Y= 5.0 11	Y= 6.0 111-
Z U	Z U	Z U	Z U	Z U	Z U	Z U	Z U	Z U
(IN) (FPS)	(IN) (FPS)	(INI (FPS)	(IN) (FPS)	(INJ (FPS)	(IN) (FPS)	(1N) (FPS)	(IN) (FPS)	(INJ (FPS)
.050 14.033 .054 14.935 .072 15.872 .094 16.922 .117 17.905 .145 18.905 .177 20.013 .226 20.935 .272 21.520 .538 22.208 .539 22.762 .471 23.511 .554 23.740 .608 24.173 .672 24.552 .756 24.996 .851 25.393 .962 25.821 1.006 26.288 1.213 26.719 1.366 27.150 1.526 27.555 1.678 27.903 1.678 27.903 1.678 27.903 1.678 27.903 1.678 29.9177 2.244 28.914 2.430 29.177 2.642 29.687 5.556 29.879 5.556 29.879 5.558 30.202 5.558 30.202 5.584 30.249	.030 16.027 .055 16.684 .073 17.524 .096 18.496 .120 19.346 .153 20.287 .187 21.051 .251 21.733 .339 22.569 .426 25.164 .534 23.804 .428 24.27 .775 25.051 .1025 26.065 1.274 26.967 1.400 27.345 1.525 27.712 1.682 28.107 1.837 28.417 1.992 28.720 2.150 28.990 2.340 29.268 2.525 29.539 2.715 28.417 1.992 28.720 2.150 28.990 2.340 29.268 2.525 29.539 2.715 20.783 2.952 29.974 3.155 30.295 3.572 30.218 5.561 30.265 5.780 30.509 5.780 30.509 5.780 30.509	.030 15.434 .055 16.414 .074 17.262 .096 19.143 .127 18.985 .165 19.675 .202 20.524 .251 21.029 .319 21.608 .400 22.229 .491 22.899 .491 22.899 .587 24.082 .772 24.582 .678 24.082 .772 24.582 .678 24.082 .772 24.585 1.110 26.245 1.110 26.245 1.244 26.785 1.244 26.785 1.242 28.053 1.832 28.053 1.832 28.053 1.832 28.053 1.832 28.053 1.832 29.665 2.026 28.642 2.026 28.645 2.059 30.0151 5.559 30.215 5.812 50.259 5.940 30.271 5.940 30.271	.051 14.721 .046 15.416 .066 16.216 .097 16.975 .109 17.761 140 19.695 .177 19.669 .226 20.415 .302 21.298 .392 21.927 .469 22.602 .554 25.221 .651 23.852 .765 24.505 .980 25.610 1.108 26.172 1.228 26.644 1.357 27.076 1.513 27.563 1.669 28.022 1.825 28.453 2.015 28.806 2.305 29.354 2.516 29.671 2.760 29.844 5.327 50.149 5.549 50.243 5.782 50.521 5.782 50.521	.032 11.265 .056 12.260 .080 13.281 .117 14.565 .149 15.594 .180 16.512 .250 17.611 .296 18.670 .556 19.556 .420 20.534 .482 21.046 .560 20.534 .482 21.046 .565 22.611 .732 25.449 .827 24.157 .044 25.529 .1.475 26.194 1.275 26.194 1.275 26.194 1.275 26.194 1.275 26.194 1.275 29.069 2.419 29.404 2.642 29.550 2.014 28.750 2.022 29.069 2.419 29.404 2.645 30.277 3.787 30.512 3.787 30.512 3.787 30.512	.125 12.799 .159 13.676 .189 14.650 .226 15.588 .267 16.512 .515 11.481 .359 18.546 .414 19.556 .474 20.296 .534 21.042 .601 21.740 .684 22.484 .767 23.160 .945 24.4548 1.059 25.273 1.195 25.975 1.575 26.694 1.510 27.293 1.656 29.742 1.955 29.828 2.252 29.199 2.448 29.514 2.655 29.423 3.100 30.082 3.578 30.293 3.758 30.293 3.758 30.293 3.758 30.293 3.956 30.506 3.956 30.506 3.956 30.506	.126 13,331 .164 13,828 .207 14,421 .248 15,229 .292 16,116 .543 17,023 .406 17,985 .472 18,918 .529 19,711 .598 20,557 .664 21,556 .752 22,074 .814 22,862 .900 23,595 .979 24,157 1,073 24,157 1,073 24,157 1,075 24,760 1,198 25,491 1,359 26,216 1,478 26,836 2,010 28,604 2,195 28,036 2,010 28,604 2,195 28,036 2,010 28,604 2,195 28,038 2,383 29,320 2,662 29,620 2,618 29,831 3,098 29,987 3,412 30,117 3,720 30,180 3,951 30,205 3,951 30,205 3,951 30,205	126 10.791 165 11.565 195 12.655 225 13.938 271 15.410 511 6.506 565 17.627 434 18.812 492 19.619 566 20.529 .644 21.396 1.957 22.293 1.037 24.633 1.140 25.333 1.254 25.979 1.599 27.256 1.659 26.665 1.559 27.259 1.659 28.205 2.059 28.677 2.229 29.690 2.445 29.427 2.652 29.690 2.455 29.290 2.455 29.290 2.455 29.290 2.455 29.290 2.455 29.30 3.507 30.211 3.817 50.249 3.817 50.24	.126 11.769 .179 12.662 .218 13.494 .255 14.381 .505 15.551 .575 16.581 .424 17.681 .504 18.711 .507 19.82 .664 20.699 .738 21.566 .836 22.577 .927 25.391 1.022 24.160 1.207 25.351 1.317 25.951 1.429 26.515 1.429 26.515 1.652 27.044 2.206 28.865 2.018 28.444 2.206 28.865 2.531 29.199 2.610 29.517 3.408 30.111 2.862 29.856 5.105 29.977 3.408 30.201 3.858 30.201 3.858 30.201 3.858 30.201
X- 4.0 FT	X= 4.0 FT	X= 4.0 FT	X= 4.0 FT	x= 4.0 FT	x= 4.0 FT	x= 4.0 FT	K= 4.0 FT X	= 4.0 FT
Y- 0.0 IN	Y= 1.0 IN	Y= 2.0 IN	T= 2.5 IN	r= 3.0 IN	Y= 3.5 IN	y= 4.0 IN	T= 5.0 IN T	= 6.0 IN
Z U	Z U	Z U	Z U	Z U	Z U	Z U	Z U	Z U
	(IN) (FPS)	(IN) (FPS)	(IN) (FPS)	(1NJ (FPS)	(IND (FPS)	(IN) (FPS)	(1N) (FPS)	(INU (FPS)
.031 14.79 .056 16.04 .072 17.06 .171 19.84 .171 19.86 .151 19.84 .173 22.52 .466 23.18 .573 22.52 .466 23.18 .573 22.52 .466 23.18 .547 25.67 .640 24.22 .731 24.67 .588 25.23 .983 25.72 .1.39 26.28 1.295 26.73 1.454 27.13 1.659 27.56 1.828 27.94 2.048 28 .276 30.4 3.565 30.11 4.537 30.22 5.566 30.4 5.276 30.4 5.276 30.4 5.276 30.4	.051         15.276           .048         16.622           .070         17.697           .097         18.701           7         .142         19.801           6         .073         17.697           7         .142         19.801           6         .275         21.484           2         .363         22.291           9         .571         25.624           9         .952         24.926           9         .952         25.893           9         1.515         27.259           9         1.556         26.721           9         1.552         28.535           10         2.264         28.656           2.90.62         3.729           17         4.579         3.529           19         1.515         2.755           10         2.264         28.656           2.704         29.062         3.773           17         4.529         30.591           16         3.428         30.591           17         4.648         30.421           16         5.275         3.441	.051 14.797 .040 16.219 .054 17.193 .060 18.094 .116 19.057 .164 19.916 .227 20.704 .503 21.455 .1366 22.156 .479 22.769 .576 25.403 .698 24.151 .825 24.777 .953 25.374 1.422 26.098 1.257 26.643 1.455 27.159 1.611 27.556 2.578 29.325 1.455 29.016 2.578 29.325 2.578 29.325 2.578 29.325 2.578 29.325 3.804 30.190 3.804 30.382 2.527 30.451 2.527 30.451 2.527 30.451 2.527 30.451	.050 14.016 .049 15.709 .069 16.724 .096 17.617 .139 19.505 .192 19.51 .254 19.855 .302 20.520 .563 21.135 .441 21.786 .525 22.382 .615 22.382 .615 22.382 .615 22.382 .615 22.501 .707 25.572 .958 24.954 1.081 25.422 .958 24.954 1.566 26.497 1.525 27.028 1.566 26.497 1.565 28.554 2.246 29.698 2.245 29.091 2.715 29.585 2.966 23.614 5.214 28.802 5.498 29.963 3.782 50.296 4.655 30.246 5.5286 30.246	.030 5.689 .065 7.915 .084 9.728 .107 12.126 .128 14.076 .151 15.606 .180 16.730 .218 17.724 .265 18.630 .513 19.545 .577 20.129 .456 20.954 .529 21.616 .625 22.577 .720 23.069 .815 23.664 .928 24.299 1.035 24.80 1.142 25.463 1.270 26.040 1.555 27.129 1.555 27.129 1.555 27.129 1.712 27.647 1.868 28.089 2.577 28.089 2.577 28.089 2.577 50.255 3.509 50.007 5.526 50.255 4.257 50.257	$\begin{array}{c} .128 \ 15.209 \\ .175 \ 15.969 \\ .266 \ 14.804 \\ .242 \ 15.925 \\ .278 \ 16.747 \\ .577 \ 17.590 \\ .527 \ 17.590 \\ .426 \ 19.176 \\ .486 \ 19.954 \\ .566 \ 20.859 \\ .426 \ 19.176 \\ .486 \ 19.954 \\ .566 \ 20.859 \\ .425 \ 22.577 \\ .835 \ 22.5688 \\ .176 \ 25.2688 \\ .176 \ 25.2688 \\ .176 \ 25.2688 \\ .176 \ 25.2688 \\ .176 \ 25.2688 \\ .1798 \ 27.759 \\ .287 \ 29.288 \\ .2795 \ 29.289 \\ .2795 \ 29.289 \\ .2795 \ 29.289 \\ .2795 \ 29.289 \\ .2795 \ 29.289 \\ .2795 \ 29.289 \\ .295 \ 29.280 \\ .295 \ 29.280 \\ .295 \ 29.534 \\ .295 \ 29.280 \\ .295 \ 29.280 \\ .295 \ 29.534 \\ .295 \ 29.280 \\ .295 \ 29.534 \\ .295 \ 29.280 \\ .295 \ 29.534 \\ .295 \ 29.280 \\ .295 \ 29.534 \\ .295 \ 29.280 \\ .295 \ 29.534 \\ .295 \ 29.280 \\ .295 \ 29.534 \\ .295 \ 29.534 \\ .295 \ 29.280 \\ .295 \ 20.524 \\ .295 \ 20.$	.131 13.894 .168 14.539 .209 15.562 .247 16.114 .294 16.944 .541 17.751 .509 20.072 .571 20.786 .656 21.950 .729 22.220 .806 22.867 .875 24.013 1.771 24.590 1.712 25.092 1.725 25.711 1.411 25.092 1.293 25.711 1.411 25.092 1.293 25.711 1.411 25.092 1.293 25.711 2.065 28.295 2.255 28.712 2.441 29.035 2.655 20.127 3.872 50.243 4.855 50.127 3.875 50.329 4.895 50.327 4.835 50.321	$\begin{array}{c} .151 \ 12.566\\ .167 \ 13.159\\ .200 \ 14.157\\ .242 \ 15.258\\ .291 \ 16.196\\ .546 \ 17.133\\ .410 \ 17.995\\ .494 \ 18.922\\ .552 \ 19.769\\ .629 \ 20.667\\ .715 \ 21.572\\ .805 \ 22.474\\ .906 \ 23.244\\ 1.121 \ 24.578\\ 1.612 \ 25.946\\ 1.212 \ 24.578\\ 1.242 \ 25.130\\ 1.349 \ 25.806\\ 1.471 \ 26.594\\ 1.212 \ 24.578\\ 1.258 \ 26.691\\ 1.775 \ 27.531\\ 1.959 \ 26.004\\ .2126 \ 28.497\\ 2.514 \ 29.925\\ 2.751 \ 29.623\\ 5.003 \ 29.956\\ 5.243 \ 50.265\\ 4.496 \ 50.450\\ 5.245 \ 50.265\\ 4.496 \ 50.450\\ 5.121 \ 50.477\\ 5.121 \ 50.477\\ 5.121 \ 50.477\\ 5.121 \ 50.477\\ 5.121 \ 50.477\\ \end{array}$	$\begin{array}{c} 129 & 12.111 \\ 171 & 15.007 \\ .200 & 14.150 \\ .240 & 15.145 \\ .275 & 16.055 \\ .513 & 16.820 \\ .513 & 16.820 \\ .513 & 16.820 \\ .513 & 16.820 \\ .513 & 16.820 \\ .513 & 16.820 \\ .513 & 16.820 \\ .513 & 10.820 \\ .513 & 10.820 \\ .513 & 10.820 \\ .513 & 10.820 \\ .513 & 10.820 \\ .513 & 20.913 \\ .702 & 21.541 \\ .708 & 22.501 \\ .702 & 21.541 \\ .708 & 22.501 \\ .702 & 21.541 \\ .708 & 22.501 \\ .702 & 21.541 \\ .708 & 22.501 \\ .702 & 21.541 \\ .702 & 21$

Table 5.3b Mean velocity data,  $U_{\infty}$  = 30 ft/sec -- Continued

Table 5	.3b Me	an yeloo	ity dat	a, $U_{\infty} =$	30 ft/	sec	Continu	ed
X= 6.0 FT	X= 6.0 FT	X= 6.0 FT	X= 6.0 FT	x= 6.0 FT	X= 6.0 FT	x= 6.0 FT	X= 6.0 FT	x= 6.0 FT
Y= 0.0 IN	T= 1.0 IN	Y= 2.0 IN	T= 2.5 IN	Y= 3.0 IN	Y= 3.5 IN	Y= 4.0 IN	Y= 5.0 IN	T= 6.0 IN
Z U	Z U	Z U	Z U	Z U	Z U	Z U	Z U	Z U
(IN) (FPS)	(IN) (FPS)	(INI (FPS)	(1N) (FPS)	(110) (FPS)	(INI (FPS)	(IN) (FPS)	(IN) (FPS)	(114) (FPS)
.028 14.035 .037 14.699 .043 15.523 .067 15.922 .064 16.477 .088 16.992 .110 17.577 .129 18.081 .163 18.621 .201 29.059 .401 20.972 .404 21.644 .599 22.266 .690 22.915 .396 22.266 .690 22.915 .396 25.518 .907 24.111 .051 24.683 1.895 25.690 1.520 25.601 .1491 26.358 1.695 26.663 .1491 26.358 1.695 26.663 .1491 27.591 2.566 28.259 2.566 28.259 2.566 29.922 3.704 29.628 3.070 29.110 3.529 29.322 3.704 29.628 3.070 29.110 3.529 29.322 3.704 29.628 3.070 29.110 3.529 29.322 3.704 29.628 3.070 29.110 3.529 29.322 3.704 29.628 3.029 5.518 4.055 29.591 5.826 30.229 5.826 30.229 5.826 30.229 5.826 30.229	.050 13.435 .045 14.028 .052 14.731 .062 15.341 .072 15.946 .068 16.477 .110 17.069 .152 17.619 .162 18.132 .189 18.641 .254 19.215 .275 19.650 .216 22.517 .275 19.650 .217 21.427 .584 22.004 .668 22.517 .1050 24.562 .1050 24.562 .1050 24.562 .135 24.842 .135 24.842 .135 24.845 .1050 24.562 .1557 26.285 1.941 27.253 .256 28.529 .5575 28.528 .576 28.528 .576 28.528 .576 29.595 .576 29.599 .557 26.295 .576 29.599 .557 26.295 .576 29.599 .557 26.295 .576 29.599 .557 26.295 .576 29.599 .575 29.784	.028 13.330 .034 14.075 .040 14.718 .050 15.317 .070 16.045 .105 17.063 .131 17.552 .162 19.132 .204 18.6579 .162 19.132 .204 18.6145 .258 19.113 .528 19.630 .595 21.032 .440 20.548 .529 21.220 .596 21.657 .664 22.076 .654 22.076 .654 22.075 .664 22.075 .664 22.076 .1554 25.757 .552 28.428 .1554 26.142 1.555 27.757 2.555 28.076 2.552 28.424 2.742 28.766 3.202 29.200 .1655 27.757 2.555 28.076 2.552 28.424 2.742 28.766 3.202 29.200 .1655 27.757 2.555 28.076 2.552 28.424 2.742 28.766 3.202 29.200 .1655 27.757 2.555 28.424 2.742 28.766 3.202 29.200 .165 29.457 3.799 29.652 4.106 29.852 4.517 29.653 4.965 50.097 .144 30.116 5.414 30.116	$\begin{array}{c} .055 \ 15, 196 \\ .046 \ 15, 995 \\ .062 \ 14, 725 \\ .062 \ 14, 725 \\ .062 \ 14, 725 \\ .062 \ 14, 725 \\ .062 \ 14, 725 \\ .062 \ 14, 725 \\ .061 \ 15, 916 \\ .175 \ 16, 485 \\ .196 \ 17, 597 \\ .238 \ 18, 101 \\ .269 \ 19, 917 \\ .238 \ 18, 101 \\ .269 \ 19, 917 \\ .259 \ 19, 916 \\ .198 \ 17, 597 \\ .259 \ 19, 916 \\ .199 \ 17, 597 \\ .259 \ 19, 916 \\ .199 \ 17, 597 \\ .295 \ 22, 916 \\ .199 \ 12, 915 \\ .145 \ 24, 455 \\ .145 \ 24, 455 \\ .159 \ 25, 956 \\ .159 \ 25, 956 \\ .159 \ 25, 956 \\ .159 \ 25, 956 \\ .159 \ 25, 956 \\ .159 \ 25, 956 \\ .1690 \ 25, 956 \\ .1690 \ 25, 956 \\ .1690 \ 25, 956 \\ .1690 \ 25, 956 \\ .1690 \ 25, 956 \\ .1690 \ 25, 956 \\ .1690 \ 25, 956 \\ .1690 \ 25, 956 \\ .177 \ 28, 939 \\ .1826 \ 26, 800 \\ .1997 \ 27, 198 \\ .205 \ 27, 568 \\ .2995 \ 23, 177 \ 28, 939 \\ .1826 \ 26, 800 \\ .1997 \ 27, 198 \\ .205 \ 27, 568 \\ .2995 \ 23, 177 \ 28, 939 \\ .1620 \ 29, 922 \\ .177 \ 28, 939 \\ .1620 \ 29, 922 \\ .156 \ 29, 922 \\ .1636 \ 29, 922 \\ .16$	.055 9.620 .055 10.700 .062 11.554 .078 12.527 .055 15.109 .120 13.784 .144 14.454 .168 15.082 .190 15.709 .227 16.529 .256 16.876 .511 17.556 .511 17.556 .570 16.240 .455 9.181 .542 20.044 .625 23.849 .725 21.690 .794 22.667 .877 22.834 .956 23.849 .726 23.849 .726 23.849 .726 23.849 .726 23.849 .726 23.849 .735 24.297 .1510 25.729 1.649 26.210 1.810 26.69 .1649 26.210 1.810 26.69 .1978 27.719 .257 28.390 .260 24.812 .355 29.218 .260 24.812 .355 29.218 .260 24.812 .584 20.635 .584 20.431 .584	, 128 11.639 , 135 12.155 , 148 12.737 , 164 15.435 , 189 12.737 , 164 15.435 , 219 15.145 , 255 15.849 , 353 16.55 , 158 19, 751 , 645 20.57 , 188 25 , 752 21.532 , 189 22.160 , 962 22.979 , 1.94 23.780 , 1.625 26.578 , 1.455 26.578 , 1.455 26.578 , 1.455 26.578 , 1.455 26.578 , 1.455 26.578 , 1.665 26.840 2.760 27.921 , 1.666 26.350 , 1.855 26.840 2.766 27.921 , 1.666 26.350 , 1.855 26.93 , 1.955 26.93 , 1.855 26.93 , 1.855 26.935 26.95	, 150 11.447 , 141 11.954 , 152 12.563 , 158 12.714 , 174 15.129 , 179 13.622 , 218 14.959 , 265 15.776 , 357 16.429 , 367 17.125 , 420 17.561 , 357 16.429 , 367 17.125 , 420 17.561 , 357 16.429 , 367 17.125 , 420 17.561 , 357 16.429 , 367 16.429 , 36	.128 11.597 .146 12.050 .166 12.496 .210 13.531 .245 14.159 .290 14.993 .551 15.933 .407 16.712 .498 17.680 .568 18.551 .651 19.438 .735 20.238 .825 21.031 1.034 22.617 1.142 23.342 1.266 24.151 1.539 25.496 1.694 26.096 1.694 26.096 1.662 26.626 0.298 27.092 2.200 27.500 2.411 27.931 2.623 28.329 2.880 28.724 3.684 29.568 4.019 29.806 4.402 29.970 5.590 30.160 5.590 30.160 5.590 30.160	$\begin{array}{c} 130 \ 10.056\\ 141 \ 11.597\\ 152 \ 11.009\\ 167 \ 12.571\\ 190 \ 12.931\\ 211 \ 15.545\\ 244 \ 14.247\\ 274 \ 14.903\\ 515 \ 15.464\\ 363 \ 16.219\\ 390 \ 16.750\\ 453 \ 17.541\\ 502 \ 18.464\\ 606 \ 19.065\\ 760 \ 20.201\\ 19.665\\ 760 \ 20.201\\ 19.665\\ 760 \ 20.201\\ 19.625\\ 2352 \ 16.92\\ 2352 \ 16.92\\ 2352 \ 16.92\\ 2352 \ 27.500\\ 2.102 \ 27.500\\ 2.552 \ 29.534\\ 3.572 \ 29.571\\ 3.202 \ 29.749\\ 4.257 \ 29.883\\ 5.572 \ 29.571\\ 3.806 \ 29.749\\ 4.257 \ 29.883\\ 5.572 \ 29.571\\ 3.806 \ 29.749\\ 4.257 \ 29.883\\ 5.574 \ 29.990\\ 5.594 \ 30.094\\ 5.594 \ 30.094\\ 5.594 \ 30.094\\ \end{array}$
X= 8.0 FT	X= 8.0 FT	x= 8.0 FT	X= 8.0 FT	x= 8.0 FT	x= 0.0 FT	X= 8.0 FT	X= 8.0 FT	x= 8.0 FT
Y= 0.0 IN	Y= 1.0 IN	Y= 2.0 IN	Y= 2.5 IN	T= 3.0 IN	Y= 3.5 [N	Y= 4.0 IN	Y= 5.0 IN	Y= 6.0 IN
Z U	Z U	Z U	Z U	Z U	Z U	Z U	Z U	Z U
LINI (FPS	INI IFPS	(INI (FPS)	(IN) (FPS)	(INJ (FPS)	(IN) (FPS)	LINI (FPS)	(IN) (FPS)	(INJ (FPS)
. 052 14,09 .039 14,63 .041 15,01 .045 15,56 .066 16,31 .064 16,96 .109 17,47 .141 18,09 .177 18,64 .216 19,11 .216 19,11 .217 19,69 .348 20,53 .426 20,94 .507 21,50 .603 22,05 .603 22,05 .605 22,57 .811 25,12 .952 25,76 .1,854 26,54 .1,854 26,54 .1,854 26,54 .1,854 26,54 .1,854 26,54 .1,854 26,54 .1,854 26,54 .5,295 28,87 .5,295 28,57 .5,295 28,57 .5,29	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} .050 \\ .050 \\ .045 \\ .055 \\ .14, 439 \\ .065 \\ .14, 439 \\ .065 \\ .14, 439 \\ .065 \\ .14, 439 \\ .065 \\ .120 \\ .16 \\ .120 \\ .16 \\ .120 \\ .16 \\ .120 \\ .16 \\ .120 \\ .16 \\ .16 \\ .16 \\ .17 \\ .16 \\ .16 \\ .17 \\ .16 \\ .16 \\ .17 \\ .16 \\ .1$	.054 10.460 .041 11.172 .050 11.922 .067 12.514 .094 13.775 .120 14.461 .150 15.155 .185 15.752 .232 16.547 .230 17.550 .457 19.237 .549 18.650 .457 19.237 .549 18.650 .457 19.237 .529 19.036 .634 23.487 .764 21.784 .954 22.878 .954 22.878 .555 25.251 1.659 25.251 1.659 25.251 1.659 25.261 2.502 27.975 .5350 28.628 .550 29.806 .420 29.635 .556 33.133 .576 33.233 .578 33.238	, 126 12, 414 , 151 15, 020 , 168 13, 529 , 197 14, 182 , 223 14, 767 , 249 15, 256 , 222 14, 767 , 249 15, 256 , 222 15, 256 , 222 15, 256 , 221 16, 437 , 355 17, 486 , 441 19, 090 , 452 18, 670 , 546 19, 193 , 601 19, 795 , 612 20, 329 , 752 20, 814 , 6, 792 21, 914 , 702 21, 924 , 704 20, 265 , 705 20, 927 , 754 29, 927 , 545 29, 927 , 545 29, 927 , 545 29, 927 , 546 29, 957 , 546 29, 95	, 127 11, 120 , 145 12, 208 , 161 12, 654 , 192 13, 129 , 209 13, 639 , 239 14, 374 , 274 14, 374 , 17, 225 , 17, 15, 656 , 372 16, 476 , 451 17, 225 , 497 18, 015 , 559 18, 746 , 450 19, 458 , 668 20, 558 , 627 19, 458 , 620 19, 458 , 770 20, 470 , 172 22, 457 , 172 22, 457 , 172 22, 457 , 172 22, 457 , 175 22, 564 , 175 22, 575 , 511 428, 227, 875 , 512 28, 613 , 512 29, 920 , 512 29, 200 , 512 29, 200	.124 11.597 .149 12.225 .169 12.225 .169 12.207 .234 14.030 .234 14.030 .541 14.030 .541 15.445 .566 16.100 .415 16.781 .463 17.980 .519 17.965 .569 18.603 .569 18.603 .569 18.603 .569 18.603 .569 18.603 .780 20.510 .864 21.153 .956 21.812 1.246 23.542 1.375 24.187 1.605 24.773 1.672 25.421 1.847 26.038 2.105 26.598 2.105 26.598 2.105 26.592 .105 26.592 .8584 29.805 3.552 4.90 3.075 28.880 3.552 4.90 3.075 28.880 3.552 4.90 3.552 4.90 3.075 28.880 3.552 3.548 2.955 3.154 5.551 30.280 5.748 30.378 6.214 50.451 6.643 50.482 6.643 50.482	. 125 11, 160 . 146 11, 946 . 176 12, 960 . 208 13, 755 . 275 14, 467 . 278 15, 256 . 379 15, 256 . 428 17, 402 . 450 18, 109 . 556 18, 751 . 614 19, 329 . 701 20, 062 . 790 20, 756 . 872 21, 350 . 954 21, 895 . 1, 252 22, 905 . 1, 253 22, 454 1, 259 22, 456 1, 125 22, 905 1, 263 25, 464 1, 596 25, 164 2, 504 27, 506 2, 525 27, 869 2, 164 29, 105 . 3, 96 29, 555 . 4, 539 2, 37 . 4, 650 29, 37 . 4, 650 29, 37 . 4, 650 29, 37 . 4, 650 29, 37 . 5, 360 29, 575 . 5, 350 35, 457 . 6, 951 30, 457

Table 5.3b Mean velocity data,  $U_{\infty}$  = 30 ft/sec -- Continued

x=10.0 Y= 0.0	FT IN	X=10.0 FT Y= 1.0 IN	x=10. Y= 2.	O FT O IN	x=10.0 T= 2.5	FT IN	x=10.0 T= 3.0	FT	x=10.0 T= 3.5	FT	x=10.0 Y= 4.0	FT	x=10.0 r= 5.0	FT	x=10.0 Y= 6.0	FT 14
Z	U (FPS)	Z U LINJ (FP	Z SI (IN)	U (FPS)	2 1110	U (FPS)	2 (110	U (FPS)	2 (110	U (FPS1	Z	U FPSI		(FPS)	2 t [NI	U (FPS)
.050	11.554	.030 11.8	.052	12.064	.031 1	1,796	.051	9.371	.130	9.021	.136	11.469	.128	11.057	. 126	9.421
.037	12.079	.040 12.4	12 .045	12.824	.054 1	2.255	.051	9.977	.146	9.737	.155	12.150	.150	11,600	.143	10.202
.049	12.571	.055 13.0	.060	15.495	.051 1	2.658	.069	10.531	. 162	10.656	.170	12.668	.170	12.056	. 164	11.040
.060	13.029	.001 15.8	.086	14.245	.075 1	3.094	.091	11,185	.165	11,551	.206	15,479	.190	12.480	.188	11,914
.074	15.598	.100 14.5	.114	14.895	.097 1	5,565	.117	11.842	.223	12.661	.258	14.552	.207	12.905	.212	12.609
.087	14.105	.127 14.9	, 160	15.625	.142 1	4.409	.145	12.586	.254	15,451	. 322	15.705	.237	13.521	.249	13,479
	14.772	.159 15.5	.204	16.271	.194 1	5.120	. 169	15.202	.297	14.329	. 587	16.594	.275	14.157	.295	14.569
.159	15.546	.200 16.2	.271	17.061	.259 1	5,745	.215	14,150	. 559	15.120	.466	17.612	. 509	14, 720	. 541	15.170
.1/6	16.285	.265 17.1	6 .328	17.755	.285 1	6.581	.262	14.9/5	. 592	15.942	.558	18.569	. 335	15.356	. 5%2	15.919
.222	17.121	.351 17.8	408	18.492	.329 1	6.916	. 502	15.661	.456	17.066	.632	19.269	.405	16.015	.448	16. 702
.2/5	17.820	.415 18.6	.500	19.328	1392 1	1.639	. 545	16.518	.517	17.910		19,985	.400	10.113	.324	10,000
100	10.100	.490 19.3	2 . 394	20.010	.4/8 1	8.528	. 391	10.9/1	.002	10,002	.817	20.131	. 540	10 472	. 270	10.067
. 390	10 964	.5/6 19.9	54 .103 75 000	21. 808	. 200 1	9.209	617	10 800	.007	10 607	1 025	22 071	-02/	10 245	770	10 964
503	20 506	701 21 2	52 015	21 020	750 2	0 754	BCC.	18 047	700	20 508	1 150	22 726	906	20 002	875	20 603
679	21 102	018 21 8	10.1	27 575	057 2	1 505	636	10 557	800	21 879	1 276	25 550	016	20.705	.060	21 200
BIR	21.795	1.040 22.4	28 1.16	25 150	054 2	1.050	721	20 215	1.020	22 182	1 405	25.875	1.019	21 480	1.086	22 045
.951	22.582	1,165 22.9	37 1.290	23.668	1.069 2	2.521	805	20.781	1,108	25.060	1.569	24.555	1.147	22.288	1,191	22.645
1.075	22,920	1,292 23.4	99 1.454	24.526	1,191 2	5.059	.905	21.585	1.366	25,795	1.745	25,246	1.506	25,208	1.297	23,224
1.219	23.503	1,461 24.1	27 1.626	24.911	1.519 2	5.628	1.018	22.084	1.584	24.659	1.912	25.843	1.485	24.017	1.415	23.764
1.558	24.041	1.627 24.7	24 1.818	25.504	1.460 2	4,174	1,142	22.758	1.794	25.321	2.085	26.507	1.654	24.686	1.548	24.554
1.530	24.616	1.841 25.4	44 2.005	26.014	1.612 2	4.708	1.267	25.350	2.009	25.934	2.294	26.865	1.872	25.470	1.656	24.774
1.698	25.156	2.044 26.1	01 2.211	26.562	1.778 2	5.276	1,309	25.926	2.216	26.462	2.496	27.555	2.081	26.130	1.798	25.510
1.909	25.777	2.261 26.6	99 2.42	27.065	1.949 2	5.777	1.550	24.597	2.470	27.052	2.762	27.916	2.335	26.873	1.944	25.846
2.114	26.361	2.467 27.1	85 2.6%	5 27,506	2.118 2	6.246	1.729	25.228	2.721	27.575	5.015	28.547	2.592	27.506	2.120	26.504
2.565	27.024	2.724 27.6	98 2.886	27.950	2.322 2	6.810	1.897	25.784	5.019	28.149	5.268	28.705	2.845	28.095	2.295	27.075
2.614	27.623	2.972 28.1	55 3.15	28.404	2,535 2	7.279	2,104	26.594	5.515	28.652	5.564	29.062	3.099	28.545	2.508	27.599
2.062	28.109	5.222 28.5	85 5.591	28.807	2.744 2	7.719	2.520	26.912	5.655	29.091	3,907	29.585	3.442	29.016	2.718	28.011
5.155	28.586	5.472 28.8	96 3.68	29,182	2.951 2	8.078	2.529	27.400	5,950	29.396	4.250	29.655	5, 781	29.580	2.972	28.450
5,486	29.072	5.724 29.1	76 4.01	29.521	5.204 2	8.475	2.741	27.818	4.290	29.690	4,584	29.881	4.163	29.665	5.226	28. 781
2.821	29.434	4.019 29.4	21 4.440	29.161	5.456 2	8. 191	2.935	28.189	4.626	29.909	4.929	50.059	4,545	29.849	5.482	29.121
4.191	29.132	4.555 29.6	4.85	29.950	5.708 2	9.101	5.199	28.580	5,005	50.089	5.309	50.158	4.929	50.056	5.779	29.457
6 116	29.990	4.694 29.8	00 0.2/3	\$0.102	1 215 2	9.544	5.447	28,889	5.428	30.245	5,687	50.268	5.268	50.149	4.121	29.716
5.533	BA REE	5.032 30.0	00 5.102	BO .209	4.617 2	9.355	8 040	20. 464	6 818	80.333	0.113	80.309	5.691	30.204	4.420	29.903
5.001	80.480	5.409 30.1	00 C. E.E.	EA 474	4 005 7	9.109	1 205	20, 722	6 780	TA 446	6.007	50.470 BAG	6.510	80.000	4.805	80.100
6 700	BA 6.12	6 120 83 8	77 7 0.33	50.410	5 266 5	1 121	4 620	20 061	7 157	KA 464	7 808	80.508	6 666	NA ARE	5.505	81 201
7.528	\$1 5.12	6 550 50 4	52 7 57	51 508	5 685 5	1 202	4.966	30.136	7 375	30 464	7 585	50.508	7. 390	\$1 464	6 029	\$0.200
7.528	\$2.542	6 957 31 4	86 7 53	51 518	6.099 3	0.299	5.417	50.262	7. 375	50.464	7.585	50.508	7. 390	50 464	6.455	30 161
7.529	30.542	7. 579 50 4	89 7.379	50.508	6.522 3	0.371	5.841	50.402	7. 373	30,464	7.585	30.508	7.390	50.464	6.966	30. 100
7.528	50.542	7.579 50.4	89 7.57	50.508	6.946 3	0.436	6.298	50,418	7.375	50.464	7.385	50.508	7.590	30.464	6.966	50.492
7.529	50.542	7. 579 50.4	89 7.575	30.508	7.560 5	0.456	6.765	50.452	7.375	30,464	7.385	30.508	7.590	30.464	6.965	30. 122
7.528	32.542	7. 579 30.4	6 7.57	30.508	7,560 5	1.456	7,145	50,489	7.575	30.464	7.585	\$0.508	7.390	30, 464	6.965	50. 197
7 820	84 6.13	" ETC EA /	7 1 1	81 519	7 560 8	A 486	7 575	\$1 490	7 878	80 464	7 505	80 500	2 800	80 101	6 0.00	84 1.37

Table 5.4a Turbulence data at x = 6 ft,  $U_{\infty}$  = 20 ft/sec

	Y=0.0 (IN)							Y=1.0	(11)					1=2.0	(174)		121
Z (IN)	(FPS)	IFPST	(FPS)	-W (FPS)	-UN	z cho	(FPS)*	(FPS)2	(FPS)2	-W	-UH	Z (110	(FPS)*	(FPS) <sup>a</sup>	(FPS)	-W (FPS)2	-JH (FPS) <sup>2</sup>
.060 .139 .206 .300 .565 .461 .608 .797 .964 1.107 1.326 2.207 3.172 4.452 5.445	5.72 4.78 4.07 5.75 5.58 3.35 5.158 2.91 2.66 2.52 2.33 1.81 1.51 1.03 .15	1.37 1.42 1.48 1.57 1.54 1.26 1.26 1.26 1.26 1.26 1.26 1.20 1.22 .06	.86 .87 .92 .95 .97 .99 .89 .74 .71 .57 .21 .15	16 17 22 23 25 22 20 19 16 05 02	.58 .62 .60 .57 .52 .48 .41 .25 .14 .05 .02	.036 .085 .156 .290 .548 .413 .522 .631 .759 .956 1.173 1.607 2.474 3.342 4.426 5.511	5.76 5.12 5.41 5.44 5.25 5.14 2.85 2.55 2.37 1.88 1.29 .90 .59 .15	1.96 1.99 2.03 1.92 1.86 1.86 1.70 1.52 1.43 1.05 .78 .31 .05	1.04 1.12 1.25 1.20 1.16 1.15 1.00 .91 .81 .54 .29 .15	26 28 50 51 55 51 55 28 25 19 15 08 01	523,56 58 54 54 54 54 54 54 54 54 54 54 54 54 54	.045 .098 .161 .290 .415 .522 .631 .739 .956 1.67 2.474 3.342 4.426 5.511	5.985 5.355 5.676 5.24 5.09 2.80 2.257 2.09 1.38 89 5.15	2.18 2.20 2.25 2.26 2.20 2.20 2.20 2.20 2.20 2.20 2.20	1.57 1.59 1.60 1.58 1.55 1.41 1.55 85 .77 .25 5.15	41 36 33 28 26 24 22 13 10 06 03 01	.65 .66 .68 .67 .65 .60 .54 .21 .16 .05 .01
		Y=2.5	5 (IN)					Y=3.(	) (III)					¥=3.5			
Z (110	(FPS)2	(FPS) <sup>a</sup>	(FPS)"	-W	-UN (FPST	(ÎN)	(FPS)	(FPS)2	IFPST	-W	-UN (FPS)2	Z (1N)	(FPS)2	(FPST	(FPS) <sup>2</sup>	-W (FPS) <sup>2</sup>	-UH (FPS)2
.035 .081 .147 .290 .348 .413 .522 .631 .739 .956 1.173 1.607 2.474 2.474 5.511	5.62 5.29 4.81 4.11 3.92 5.72 5.51 3.51 3.50 2.68 2.18 1.46 1.00 .49 .16	5.74 5.78 5.72 5.67 5.28 5.17 2.78 2.465 1.65 .95 .36 .19	2.65 2.68 2.71 2.75 2.70 2.62 2.29 2.19 1.59 .24 .15	46 41 30 26 25 21 18 14 10 05 01	1.05 1.06 1.07 1.09 1.05 1.05 1.00 .92 .74 47 .18 .06 .01	.005 .129 .192 .299 .340 .415 .522 .651 1.739 .956 1.173 1.607 2.474 3.542 4.426 5.511	5.44 5.90 5.97 5.97 5.97 5.97 5.97 5.97 5.97 5.97	5,45 5,54 5,62 5,38 2,84 2,84 2,48 1,66 .95 .58 .21 .02	1.62 1.64 1.65 1.65 1.45 1.20 .98 .80 .47 .28 .13	53 54 54 50 47 37 35 215 07 03 01	.95 .97 .94 .89 .82 .74 .64 .47 .26 .14 .06	.130 .177 .251 .571 .435 .544 .652 .761 .978 1.195 1.629 2.496 3.364 4.448 5.533	5.41 5.66 5.94 5.67 5.46 5.07 4.39 3.67 1.65 .46 1.05 .46 .16	3.16 3.12 2.95 2.83 2.44 2.19 1.63 1.54 .70 .27 .05	1.68 1.75 1.66 1.61 1.95 1.40 1.16 .92 .47 .27 .15	52 529 46 29 158 03 01	1.28 1.29 1.34 1.26 1.18 1.02 .88 .57 .22 .08
		Y=4.0	(1N)					1=5.0	(1)a					1-6.0	(1)0		
	(FPS)2	(FPS)2	(FPS)2	-W (FPS)2	-UH IFPST	Z (1N)	(FPSI <sup>2</sup>	IFPS?	(FPS)2	-W	-UH (FPS) <sup>2</sup>		UFPS)2	TV FPSr	FPS)2	-W	-UH
.140 .185 .251 .371 .435 .542 .761 .978 1.195 1.628 2.496 3.364 4.448 5.535	5.35 5.18 5.20 5.17 4.99 7 4.77 2.68 1.68 1.68 1.68 1.03 .15	5.29 5.26 5.02 2.95 2.55 2.21 1.81 1.08 .77 .25	1.69 1.71 1.74 1.70 1.65 1.35 1.15 1.00 .86 .64 .26	58 42 59 57 54 27 25 07 05 01	.97 .99 .96 .92 .88 .77 .68 .52 .28 .17 .06	.195 .180 .244 .571 .435 .544 .652 .761 .978 1.195 1.628 2.496 5.553	5.41 5.55 5.09 5.15 5.02 4.81 4.51 5.76 2.92 4.81 1.09 4.51 5.76 2.92 4.81 1.09 4.51 5.76 2.92 4.51 5.76 2.92 4.51 5.76 5.76 5.76 5.76 5.76 5.76 5.76 5.76	4.02 4.02 4.02 4.02 4.02 4.02 4.02 4.02	2.05 2.07 2.08 1.91 1.65 1.75 1.04 .79 .44 .12	26 29 25 25 21 18 09 05 01	1.03 1.04 1.05 1.03 .98 .89 .89 .85 .57 .26 .16 .01	.135 .100 .240 .571 .435 .544 .652 .761 .978 1.195 1.628 2.496 3.364 4.448 5.533	5.34 5.13 5.13 5.37 5.34 5.21 5.04 4.81 4.27 7.2.99 1.69 98 .14	5.67 5.57 5.48 5.43 5.11 2.73 2.34 1.91 1.34 .76 .25 .07	2.16 2.21 2.26 2.15 2.00 1.90 1.40 .94 8 .12	25 28 24 22 19 17 15 03 01	1.09 1.12 1.14 1.08 1.04 .90 .67 .29 .16 .05

Table 5.4a Turbulence data at x = 10 ft,  $U_{\infty}$  = 20 ft/sec

	Y-0.0 (IN)							Y=1.0	(1)					Y=2.0	(1N)		
	W (FPSf	W (FPS)2	IFPSP	-W (FPS)2	-UN	Z (1N)	(FPS)2	(FPS)2	(FPS)2	-W	-UH (FPS)2	2 (1N)	(FPS/ <sup>2</sup>	(FPS)2	HH (FPS)2	-W (FPS)2	-UH (FPS)2
.035 .084 .170 .296 .535 .701 .912 1.629 1.629 2.284 3.340 4.817 6.084	8.42 5.88 4.16 5.89 5.70 5.70 5.50 2.95 1.95 1.12 .62 .25	2.09 2.11 2.12 2.08 2.06 1.77 1.61 1.32 .16 .16	1.55 1.38 1.39 1.42 1.43 1.44 1.25 1.12 .87 .30 .16	18 18 15 15 14 12 07 01	.59 .61 .65 .62 .59 .34 .21 .08 .02	.035 .084 .170 .311 .385 .533 .704 .912 1.651 2.284 3.540 4.817 6.084	6.79 5.01 4.59 5.99 5.81 5.65 5.56 2.56 1.93 1.15 .55 .25	2.75 2.80 2.73 2.61 2.55 2.48 1.87 1.44 1.56 .64 .17	1.15 1.15 1.29 1.31 1.26 .89 .61 .58 .16	23 27 24 19 14 12 07 04 01	.61 .62 .64 .62 .61 .53 .26 .15 .10	.035 .084 .169 .524 .533 .704 .912 1.229 1.651 2.284 5.540 4.917 6.084	8.11 6.42 4.45 4.55 4.55 4.21 5.99 5.69 5.27 2.78 2.08 1.17 .65 .27	2.48 2.56 2.39 2.37 1.96 1.79 1.52 1.29 .47 .21	1.19 1.21 1.28 1.40 1.42 1.44 1.04 .92 .96 .56 .22	28 27 26 20 12 07 07 05	.54.58.59.527 .59.59.527 .2005
		¥=2.5	(IN)					¥=5.0	(1N)					¥=3.5	(IN)		
Z (IN)	w rPsr	(FPS)2	(FPS)2	-W (FPS)2	-UH	Z (1N)	(FPST	(FPSI <sup>2</sup>	(FPS)2	-W	-UH	Z (IN)		W (FPS) <sup>2</sup>	(FPS)2	-W (FPS)2	-UH (FPS)2
.051 .084 .170 .565 .401 .549 .717 .929 1.245 1.667 2.501 3.356 4.834 6.100	8.53 6.69 5.63 5.19 5.14 4.86 4.52 4.09 5.48 2.75 1.93 1.24 .59 .21	2.16 2.21 2.27 2.36 2.41 2.17 2.11 1.51 1.05 .52 .23	1.09 1.12 1.14 1.30 1.39 1.44 1.35 1.21 .93 .57 .21	57 56 25 19 08 04 01	.66 .68 .69 .72 .64 .54 .54 .51 .07 .02	.105 .168 .274 .590 .549 .717 .929 1.245 1.667 2.501 3.356 4.834 6.100	8.35 6.47 5.94 5.68 5.34 4.91 4.27 5.46 2.77 1.82 1.18 .51 .21	2.15 2.27 2.32 2.34 2.28 2.06 1.45 1.12 .45 1.12 .18	1.52 1.70 1.80 1.95 1.90 1.55 1.44 .91 .39	37 45 16 13 09 07 04 01	.84 .96 .90 .76 .58 .37 .24 .08 .02	.126 .170 .276 .443 .549 .717 .929 1.245 1.667 2.301 3.956 4.836 6.100	7.766.90 6.21 5.28 5.29 8.99 8.99 8.99 8.20 5.25 1.61 .21	3.10 3.39 3.37 2.90 1.97 1.43 .67 .22	1.22 1.31 1.43 1.46 1.33 1.21 1.10 .54 .37 .17	47 39 22 09 03 01	.90 .92 .94 .80 .52 .18 .09 .03
		1-4.	0 (IN)					۲=5.	0 (11)					¥=6.	0 (IND		
	W (FPS/	(FPS)	(FPS)	-W (FPS)	-UH	Z (Ita)	(FPS)	v (FPS)	(FPS)	-W	-UR (FPS)2		UU IFPSI	W IFPS	(FPS)2	-W (FPS)	-Uk (FPSI <sup>2</sup>
.126 .170 .276 .445 .549 .717 .929 1.245 1.667 2.301 3.356 4.667 4.667	7.36 6.90 6.60 5.67 5.31 4.75 5.31 2.22 1.29 5.4	2.69 2.74 2.88 2.92 2.71 1.66 1.24 .57	1.35 1.43 1.54 1.67 1.63 1.23 67 	55 55 27 22 16 05 01	94 95 91 .95 .78 .61 .78 .61 .22 .02	.126 .168 .274 .401 .549 .717 .929 1.245 1.667 2.501 <b>3.556</b> 4.856 6.100	7.77 7.52 7.42 7.21 6.87 5.62 5.47 2.47 2.47 2.47 2.47 2.47 2.47 2.47 2	2.48 2.50 2.74 2.77 2.74 2.77 2.43 2.43 2.43 2.43 2.43 2.43 2.43 2.43	1.42 1.47 1.52 1.56 1.60 1.47 1.35 .91	25 26 28 30 20 12 12 07 04	1.06 1.04 1.03 1.01 .93 .70 .40 .27 .11 .02	.126 .169 .274 .445 .549 .717 .929 1.245 1.667 2.301 5.356 4.836 6.100	8.16 7.61 6.95 6.62 6.51 5.67 4.96 5.67 2.47 1.57 .68	2.23 2.34 2.39 2.11 2.00 1.94 1.72 1.12 .51	1.63 1.94 2.11 2.02 1.82 1.78 1.56 .92 .21	27 25 21 18 14 10 08 05 03	1.04 1.09 1.17 1.08 1.03 .77 .44 .28 .10

Table 5.4b Turbulence data at x = 6 ft,  $U_{\infty}$  = 30 ft/sec

Y-0.0 (1N)	Y=1.0 (IN)	V-2.0 (1ND
2 00 10 HA -00	GLI 2 (DD (VT LLI - UV - ULI PSI <sup>1</sup> (IN) (FPSI <sup>1</sup> (FPSI <sup>1</sup> (FPSI <sup>1</sup> )	
.030 14.53 .075 11.91 .131 9.95 .244 8.93 .500 8.64 2.52 1.1939 1 .661 7.63 2.67 1.2543 1 .661 7.63 2.67 1.2543 1 .660 7.12 2.63 1.2846 1 .964 5.71 2.51 1.2946 1 .964 5.71 2.51 1.1940 1.107 5.32 2.46 1.0836 1.536 4.79 2.39 .9932 1.668 3.69 1.92 .9927 1.668 3.69 1.92 .9927 1.664 3.69 1.92 .9927 1.645 .72 1.92 1.28 .5715 5.445 .26 .12 .0202	.025 15.10 .061 12.91 .127 10.61 .239 9.21 .06 220 8.07 5.70 1.2771 1.13 .06 340 8.50 3.70 1.5173 1.15 .05 .413 8.15 3.71 1.5472 1.16 .01 .522 7.75 3.59 1.5764 1.18 .00 .651 7.35 3.47 1.5564 1.18 .01 .596 6.07 5.18 1.2155 1.05 .01 .596 6.07 5.18 1.2155 1.05 .59 1.607 4.40 2.19 .9756 .50 .59 1.607 4.40 2.19 .9756 .50 .51 2.44 2.88 1.25 .42 .24 .56 .17 5.542 1.85 .73 .5618 .24 .04 4.426 .75 .51 .1500 .05 .01 5.511 .24 .12 .0502 .02	$\begin{array}{cccccccccccccccccccccccccccccccccccc$
Y=2.5 (IN)	T=3.0 (D)	7=3.5 (IN)
Z ឈ ឃ អា -ឈ -វ អេស តាមនាំពេមនាំពេមនាំពោមនាំពោ	CHI Z CO W HA -CW -CHI PSI <sup>1</sup> (INI (FPSI <sup>1</sup> (FPSI <sup>1</sup> (FPSI <sup>1</sup> )	2 00 10 HH -00 -00 (1N) (FPS) <sup>3</sup> (FPS) <sup>2</sup> (FPS) <sup>2</sup> (FPS) <sup>2</sup>
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	.000 11.40 .125 11.90 .105 12.59 .292 12.54 .40 .415 11.96 4.77 2.62 -1.61 1.82 .50 .522 11.22 4.75 2.67 -1.59 1.82 .51 .651 10.32 4.65 2.71 -1.51 1.81 .51 .759 9.59 4.59 2.54 -1.44 1.74 .44 .956 8.22 3.61 2.52 -1.16 1.54 .59 1.173 6.66 3.41 2.1995 1.37 .24 1.607 5.10 2.55 1.8271 .95 .11 2.474 3.06 1.72 1.5743 .45 .11 2.474 3.06 1.72 1.5743 .45 .11 2.474 3.0675 .22 .2107 .07 .22 5.511 .24 .15 .0603 .02 .02	$\begin{array}{cccccccccccccccccccccccccccccccccccc$

7-4.0 (IN)				T-5.0 (IN)					Y-6.0 (IND								
Z (IN)	W IFPSP	W (FPS)	(FPSť	-W IFPSI	-UH (FPSt	Z (IN)	UU (FPST	(FPS/2	(FPS)2	-W (FPS)	-UH (FPS)		(FPS)*	W IFPSI	(FPS) <sup>2</sup>	-W (FPS) <sup>2</sup>	-UN IFPST
.155 .176 .242 .571 .455	15.02 13.22 13.45 15.33 15.06	6.20	3.52 3.57	72	2.02	.135 .182 .248 .571 .435	13.89 14.37 14.03 15.73 13.55	6.30	5.54 5.61	75	2.25	.135 .180 .242 .371 .435	11.40 12.01 12.40 12.57 12.76	6.58 6.67	3.81 3.86	45	2.43
.544 .652 .761	12.22	5.97	2.98	69	1.98	.652	15.51	5.73	2.83	75	2.23	.652	12.53	5.87	5.41	54	2.45
.978 1.195 1.629	10.17 8.43 6.01	4.81 4.51 5.45	1.85	55	1.56	1,195	9.55	5.57	1.57	50	1.65	1,195	9.40	5.87	2.24	51	1.81
2.496 3.364 4.448 5.553	5.29 1.89 .75	2.80	.18	25 11 06 01	.41 .25 .07	2,496 3,564 4,448 5,533	5.66 1.98 .76 .25	1,61 ,82 ,49	.18	25	.48	2,496 3,364 4,448 5,533	1.92	.56	.58	10	.27

Table 5.4b Turbulence data at x = 10 ft,  $U_{\infty}$  = 30 ft/sec

1-0.0 (IN)			Y=1.0 (IN)					T=2.0 (IN)									
Z (IN)	(FPS)*	(FPS) <sup>2</sup>	(FPS)2	-W (FPS) <sup>3</sup>	-ON (FPS) <sup>2</sup>	2 (1N)	(FPS)2	W (FPS/	IFPST	-W	-UH	Z (IN)	UU (FPS/	W (FPS) <sup>1</sup>	IFPS:	-W	-UH
.035 .084 .170 .311 .385 .535 .701 1.229 1.229 1.251 2.294 3.540 4.817 6.084	13.20 9.44 8.55 8.29 8.22 8.16 7.79 7.15 6.36 5.54 4.08 2.14 .97 .46	4.68 4.87 5.20 4.92 4.85 4.50 3.92 2.85 2.37 1.04 .49	2.56 2.58 2.60 2.66 2.72 2.55 2.20 1.76 1.48 .52 .21	39 40 59 37 35 23 19 14 06 02	1,15 1,17 1,23 1,21 1,29 1,21 1,29 1,02 .69 .30 .11 .02	.035 .084 .169 .511 .533 .701 .912 1.651 2.284 5.340 4.017 6.094	15.69 10.87 9.45 9.22 9.13 8.69 8.22 7.90 7.45 6.18 4.42 2.64 1.25 .51	4.99 4.97 5.12 5.18 4.66 5.64 5.27 2.67 1.70 .78 .50	1.96 2.26 2.45 2.54 1.88 1.75 1.66 1.21 .47 .17	55 56 51 42 36 31 16 10 03	1,10 1,12 1,14 1,19 1,17 1,15 1,02 .75 .51 .09 .02	.035 .084 .170 .543 .587 .533 .704 .912 1.229 1.651 2.284 4.817 6.084	16.26 11.80 9.64 9.20 8.84 8.57 8.20 7.59 6.23 4.41 1.20 .50	5.76 5.95 6.10 6.38 5.51 4.70 5.81 3.11 2.55 .93 .35	2.08 2.94 3.05 3.12 2.58 1.87 1.71 1.50 1.19 .43 .20	75 77 61 45 17 02	1.26 1.26 1.25 1.25 1.25 1.15 .93 .64 .58 .12
		¥=2.5	5 (1N)					Y=3.(	) (IN)					¥=5.9	5 (IN)		
Z	(FPS)	(FPS) <sup>2</sup>	IFPS ?	IFPS P	-UH	Z (1N)	w (FPST	(FPST	(FPS)2	-W IFPSI	-ON	Z	UU (FPS/	(FPS) <sup>2</sup>	IFPST	-OV (FPS)*	-UN (FPS) <sup>2</sup>
.055 .094 .168 .565 .401 .549 .717 .929 1.245 1.667 2.501 5.556 4.034 6.100	14,56 10,89 9,05 8,72 8,68 8,50 8,50 8,50 8,50 8,50 8,50 8,50 8,5	6.06 6.11 6.57 6.42 5.28 4.29 3.40 2.20 1.56 .98 .41	5.62 5.67 5.78 5.75 5.15 2.41 2.03 1.58 1.07 .44 .21	89 91 95 71 46 51 25 17 07 02	1.34 1.35 1.37 1.35 1.25 1.25 1.09 .62 .57 .11	.105 .169 .274 .401 .549 .717 .929 1.245 1.667 2.501 5.356 4.834 6.100	16.52 13.56 11.72 11.30 10.92 10.58 9.42 7.53 6.50 4.41 2.69 1.04 .57	5.49 5.89 5.52 5.29 4.77 5.73 2.41 1.58 .70 .36	5.89 5.94 5.74 5.56 5.34 2.95 1.88 1.11 .13	84 86 87 79 64 45 18 01	1,74 1,82 1,78 1,64 1,46 1,25 .78 .46 .26 .02	.126 .168 .276 .443 .549 .717 .929 1.245 1.667 2.501 5.556 4.856 6.100	13.42 14.05 13.96 15.15 12.86 12.28 10.28 5.27 2.73 1.04	5.58 5.76 5.37 4.95 4.26 3.47 2.06 .76 .33	5.18 5.24 3.16 2.92 2.54 2.12 1.65 .93 .11	89 97 95 78 36 36 16 01	1.92 1.96 1.97 1.72 1.45 1.23 .77 .44 .21
		۲-4.	0 (IN)					1-5.	0 (11)					Y-6.(			
Z (IN)	(FPS) <sup>2</sup>	(FPS)	(FPS)	-W (FPS)	-OH (FPS)*	Z (1N)	(FPS)*	(FPSI	(FPS)*	-W	-UN (FPS)*	Z	00 (FPS)*	VV (FPST	IFPS)*	-W	-UN
.126 .170 .276 .443 .549 .717 .929 1.245 1.667 2.301 3.356 4.836 6.100	16.46 17.92 15.89 14.90 14.37 13.92 12.87 10.69 8.12 5.14 2.74 1.04 .42	4.86 4.94 5.13 4.84 4.51 3.85 3.01 1.95 .85 .41	2.83 5.07 5.50 2.87 2.31 1.92 1.52 1.19 .49 .15	82 86 81 76 71 38 22 12 08 01	2.09 2.12 2.16 2.05 1.71 1.55 .05 .39 .17 .05	.126 .168 .274 .401 .549 .717 .929 1.245 1.667 2.301 3.3556 4.836 6.100	15.95 16.31 15.63 14.67 14.51 15.65 12.54 10.28 8.06 5.03 2.48 .99 .41	4.83 4.92 5.08 4.69 4.52 5.64 2.94 1.86 .77 .51	3.06 3.15 5.54 2.52 2.11 1.75 1.44 1.07 .39 .15	76 80 72 54 54 21 12 08 01	2.06 2.07 2.09 2.04 1.72 1.40 .59 .16 .03	.126 .168 .274 .401 .549 .717 .929 1.245 1.667 2.301 5.356 4.936 6.100	15.28 15.93 14.87 13.90 13.49 15.17 12.25 10.59 5.24 2.26 .87 .32	4.76 4.84 4.89 4.46 4.14 5.47 2.78 1.63 .75 .26	3.84 5.96 4.17 3.75 5.250 1.92 1.04 .50 .13	45 47 58 35 27 20 11 06 01	2.18 2.23 2.26 2.15 1.51 1.51 1.51 .88 .46 .16

Table 5.5 Data on one-dimensional spectra at x = 6 ft,  $U_{\infty}$  = 30 ft/sec

Y=0.	0 IN	Y=0.4	0 IN	Y=0.	0 1N	Y=0.	0 111	Y=0.1	0 [1]	r=0.	0 [14
	075 IN	Z= .	131 IN	Z=	244 1N	2= .4	602 111	Z=1.	162 [1]	Z=2.	2 <b>39</b> [14
K	E(K).00	к	E (K) . 00	K	E (K) / 100	(1 FT)	E (K) .700	K	E(K)/100	K	E (K) . JU
(1.FT)	(FT)	(1. FT)	(FT)	(1.FT)	(FT)		(FT)	(1/FT)	(FT)	(1.FT)	(FT)
1.33E+00 1.93E+00 2.55E+00 3.42E+00 4.89E+00 7.60E+00 9.61E+00 1.59E+01 1.59E+01 1.59E+01 5.54E+01 5.54E+01 6.82E+01 1.11E+02 1.48E+02 3.70E+02 5.46E+02 6.77E+02 9.16E+02	8,55E-03 2.03E-02 1.90E-02 1.90E-02 1.55E-02 1.24E-02 1.02E-03 5.45E-03 4.43E-03 2.65E-03 2.13E-03 8.71E-04 5.84E-04 5.86E-05 7.96E-07	1,15E+00 1,49E+00 2,27E+00 3,72E+00 4,66E+00 6,15E+00 1,24E+01 1,75E+01 2,24E+01 1,75E+01 2,24E+01 1,31E+02 2,36E+01 1,31E+02 2,15E+02 2,96E+02 4,49E+02 5,50E+02 7,59E+02	1,76E-02 2,34E-02 1,81E-02 1,81E-02 1,69E-02 1,07E-02 1,07E-03 7,22E-03 4,02E-03 2,83E-03 5,75E-03 5,75E-03 5,75E-03 5,75E-03 5,75E-03 5,75E-03 5,75E-03 5,75E-04 4,96E-04 5,29E-05 4,19E-06	1.06E+00 1.49E+00 2.13E+00 4.05E+00 6.05E+00 6.59E+00 1.37E+01 2.75E+01 2.75E+01 2.75E+01 6.97E+01 6.97E+01 1.37E+02 2.30E+02 3.30E+02 4.46E+02 5.78E+02 7.45E+02	2.37E-02 2.81E-02 2.57E-02 2.50E-02 2.59E-02 1.25E-02 1.25E-02 1.25E-03 5.01E-03 1.75E-03 1.75E-03 1.75E-03 1.75E-04 4.01E-04 2.96E-04 1.49E-04 2.52E-05 2.02E-06	9,21E-01 1,34E+00 2,34E+00 3,45E+00 5,75E+00 9,26E+00 1,26E+01 1,67E+01 2,53E+01 5,49E+01 5,49E+01 5,49E+01 5,49E+01 1,07E+02 1,56E+02 2,30E+02 2,30E+02 5,20E+02 5,20E+02 5,20E+02 6,32E+02	5,80E-02 5,22E-02 4,16E-02 2,58E-02 2,58E-02 1,79E-02 1,55E-02 7,54E-03 4,66E-03 5,18E-03 1,97E-03 1,97E-03 1,97E-03 1,97E-03 1,97E-03 1,97E-04 4,72E-04 4,72E-04 2,72E-04 1,51E-04 2,64E-05 2,12T-06	8,26E-01 1,16E+00 2,14E+00 2,93E+00 5,09E+00 9,92E+00 1,63E+01 2,22E+01 3,13E+01 4,11E+01 1,00E+02 1,26E+02 1,26E+02 2,22E+02 5,23E+02 4,26E+02 5,51E+02	4.24E-02 4.92E-02 3.99E-02 3.72E-02 2.99E-02 2.37E-02 8.39E-03 3.43E-03 2.47E-03 1.35E-03 1.35E-03 1.35E-03 1.35E-04 3.95E-04 3.95E-04 3.95E-04 1.21E-04 1.21E-04 2.35E-05 2.29E-06	6.97E-01 1.10E+00 1.55E+00 2.76E+00 4.20E+00 9.92E+00 1.24E+01 2.04E+01 2.74E+01 4.04E+01 5.61E+01 1.56E+01 1.52E+01 1.52E+02 2.77E+02 2.74E+02 3.96E+02 5.10E+02	6.95E-02 5.79E-02 5.51E-02 5.51E-02 3.93E-02 1.96E-02 1.96E-02 1.96E-02 1.96E-03 2.42E-03 1.34E-03 1.34E-03 2.42E-03 1.34E-04 5.96E-04 2.64E-04 2.64E-04 2.61E-05 1.30E-06
Y=1.	0 IN	7=1;	0 IN	Y=1.	0 IN	7=1.	0 IN	Y=1.	0 IN	Y=1.	0 IN
Z= .	061 IN	2= ;	127 IN	Z= .	258 IN	Z= .	598 IN	Z=1.	179 IN	Z=2.	158 IN
K	E (K) .700	K	E (K) . W	K	E (K) . UU	K	E (K) /W	K	E(K)/W	к	E(K)/W
(1.FT)	(FT)	(1.FT)		(1.FT)	(FT)	(1.FT)	(FT)	(1/FT)	(FT)	(1/FT)	(FT)
1.31E+00 1.80E+00 2.93E+00 4.22E+00 5.81E+00 1.04E+01 1.52E+01 2.06E+01 2.06E+01 3.87E+01 5.25E+01 6.29E+01 1.19E+02 2.32E+02 5.16E+02 5.16E+02 5.36E+02 5.36E+02 5.36E+02 5.36E+02 5.36E+02	8.74E-03 1.34E-02 1.70E-02 1.81E-02 1.81E-02 1.34E-02 1.21E-02 8.55E-03 5.06E-03 5.00E-03 5.00E-03 5.00E-03 2.45E-03 1.15E-03 6.96E-04 2.39E-04 2.39E-04 2.59E-05 4.17E-07	1.12E+00 1.65E+00 2.22E+00 3.76E+00 1.01E+01 1.77E+01 1.77E+01 4.39E+01 6.83E+01 9.51E+01 1.29E+02 1.87E+02 2.59E+02 2.59E+02 2.59E+02 5.89E+02 5.89E+02 5.89E+02 8.21E+02	1.51E-02 2.21E-02 1.09E-02 2.29E-02 1.95E-02 1.34E-02 1.34E-02 5.75E-05 2.74E-05 2.74E-05 1.82E-03 1.82E-03 1.82E-03 9.67E-04 5.65E-04 1.56E-04 1.56E-04 1.56E-04 1.56E-04 1.50E-04 1.50E-04	1.06E+00 1.45E+00 5.07E+00 5.02E+00 6.97E+00 8.99E+00 1.15E+01 1.66E+01 2.47E+01 8.62E+01 1.35E+02 2.81E+02 2.81E+02 2.81E+02 2.81E+02 2.91E+02 3.75E+02 4.96E+02 6.01E+02 7.45E+02	2.47E-02 2.31E-02 3.27E-02 3.25E-02 2.09E-02 1.61E-02 1.32E-02 1.32E-02 1.32E-03 4.73E-05 2.96E-03 3.473E-05 2.96E-03 8.55E-04 2.51E-04 2.51E-04 2.51E-04 2.51E-04 2.51E-04 2.51E-05 1.31E-05 1.98E-06	8.39E-01 1.22E+00 2.97E+00 3.60E+00 7.35E+00 9.29E+00 1.29E+01 1.29E+01 1.29E+01 1.29E+01 1.25E+01 3.555E+01 3.555E+02 2.72E+02 3.66E+02 4.39E+02 4.39E+02 6.34E+02	5.00E-02 2.37E-02 4.15E-02 5.84E-02 5.17E-02 2.42E-02 1.57E-02 1.24E-02 1.00E-03 5.10E-03 5.10E-03 5.10E-03 5.10E-03 6.72E-04 1.35E-04 4.72E-06 1.72E-06	7.47E-01 1.10E+00 1.72E+00 2.77E+00 4.90E+00 8.74E+00 1.21E+01 1.62E+01 2.40E+01 3.30E+01 4.45E+01 9.94E+01 9.94E+01 9.94E+01 1.54E+02 3.42E+02 3.42E+02 3.42E+02 5.87E+02 5.87E+02	2.33E-02 3.56E-02 5.62E-02 2.84E-02 2.30E-02 2.30E-02 1.07E-02 7.43E-03 3.26E-03 2.22E-03 7.98E-04 4.50E-04 4.50E-04 1.22E-03 7.98E-04 1.22E-03 7.98E-04 1.22E-05 1.22E-05 1.22E-05 1.22E-05 1.22E-04 1.14E-04 2.64E-06 1.06E-06	6.97E-01 1.00E+00 1.72E+00 2.26E+00 3.20E+00 4.27E+00 7.06E+00 1.10E+01 1.92E+01 1.92E+01 3.27E+01 4.56E+01 6.25E+01 6.25E+01 1.15E+02 2.16E+02 2.66E+02 2.66E+02 5.06E+02	5.68E-02 5.54E-02 5.54E-02 3.554E-02 3.52E-02 3.52E-02 2.84E-02 2.84E-02 2.84E-02 2.84E-02 2.84E-02 7.72E-03 5.22E-03 8.25E-04 4.76E-04 2.41E-04 1.40E-04 1.42E-06
Y=2.	0 IN	Y=2.	0 IN	Y=2.	0 IN	7=2	0 IN	Y=2.	0 IN	Y=2.	0 IN
	071 IN	Z=	157 IN	Z=	244 IN	Z=	609 IN	Z=1.	191 IN	Z=2.	259 IN
K	E(K)/W	K	E (K) / W	K	E (K) /W	K	E(K)/W	K	E(K) /00	(1/FT)	E (K) /00
(1/FT)	(FT)	(1/FT)	(FT)	(1/FT)	(FT)	(1/FT)	(FT)	(1/FT)	(FT)		(FT)
1.24E+00 2.74E+00 3.90E+00 5.97E+00 8.75E+00 1.27E+01 1.60E+01 5.20E+01 5.20E+01 7.97E+01 1.15E+02 2.34E+02 5.15E+02 2.34E+02 5.92E+02 7.41E+02 9.50E+02	7.98E-03 2.51E-02 2.30E-02 2.40E-02 1.49E-02 1.45E-02 8.16E-03 4.90E-02 4.90E-03 4.90E-03 4.90E-03 4.90E-03 4.90E-03 4.90E-03 4.90E-03 4.90E-03 8.99E-04 4.52E-04 2.52E-04 2.52E-04 3.95E-05 6.87E-06 1.37E-06	1.152+00 1.642+00 2.452+00 3.462+00 4.962+00 6.672+00 9.912+00 1.322+01 1.742+01 2.642+01 3.552+01 4.752+01 6.822+01 9.182+02 2.212+02 2.212+02 2.122+02 4.612+02 6.852+02 6.852+02 6.852+02	1.04E-02 2.92E-02 2.92E-02 2.14E-02 1.65E-02 1.65E-02 9.22E-03 7.17E-05 4.70E-05 2.96E-03 1.96E-03 7.96E-03 7.96E-03 7.96E-03 7.96E-03 7.95E-04 4.51E-04 1.10E-04 1.10E-04 1.46E-05	9.62E-01 1.40E+00 2.52E+00 3.25E+00 3.25E+00 9.45E+00 1.25E+01 1.66E+01 2.20E+01 4.64E+01 6.55E+01 9.65E+01 9.65E+02 2.02E+02 2.02E+02 4.21E+02 5.99E+02 4.21E+02 5.99E+02 4.21E+02 5.99E+02 4.21E+02 5.99E+02 4.21E+02 5.99E+02 4.21E+02 5.99E+02 4.21E+02 5.99E+02 4.21E+02 5.99E+02 4.21E+02 5.99E+02 4.21E+02 5.99E+02 4.21E+02 5.99E+02 4.21E+02 5.99E+02 4.21E+02 5.99E+02 4.21E+02 5.99E+	1.56C-02 3.57C-02 3.55C-02 3.55C-02 3.57C-02 2.41C-02 1.42C-02 9.00C-03 5.90C-03 5.90C-03 3.62C-03 1.24C-03 1.24C-03 1.24C-03 7.40E-03 3.62C-03 1.24C-03 7.75C-04 4.95C-04 3.57C-04 1.60C-06	9.32E-01 1.13E+00 2.00E+00 5.09E+00 6.99E+00 9.99E+00 1.47E+01 1.95E+01 2.37E+01 5.37E+01 5.37E+01 5.37E+01 1.4E+02 1.70E+02 2.41E+02 3.55E+02 4.97E+02 5.55E+02 6.39E+02	2.61E-02 4.66E-02 3.20E-02 3.20E-02 3.05E-02 9.57E-02 9.57E-02 9.57E-03 5.20E-05 5.21E-03 1.95E-03 7.09E-04 4.10E-04 3.14E-04 1.42E-04 1.42E-04 1.42E-04 1.42E-04	7.60E-01 1.00E+00 3.50E+00 3.50E+00 3.55E+00 4.70E+00 9.56E+00 9.56E+00 1.20E+01 1.71E+01 1.30E+01 4.70E+01 1.00E+02 1.50E+02 3.13E+02 3.13E+02 3.56E+02 3.13E+02 3.56E+02 3.13E+02 3.56E+02 3.56E+02 3.55E+	5.20E-02 5.77E-02 5.00E-02 4.10E-02 2.25E-02 2.25E-02 2.25E-02 1.22E-02 7.62E-05 5.60E-03 5.60E-03 3.99E-03 2.27E-03 7.64E-04 4.04E-04 4.04E-04 2.02E-06 2.02E-06 2.22E-07 2.22E-	6.96E-01 9.13E-01 1.44E+00 2.20E+00 3.25E+00 5.96E+00 7.57E+00 1.05E+01 1.50E+01 1.92E+01 3.92E+01 3.92E+01 3.92E+01 1.56E+02 1.75E+02 2.94E+02 3.74E+02 5.01E+02 5.01E+02	4.96E-02 4.93E-02 5.33E-02 5.15E-02 3.94E-02 3.94E-02 3.94E-02 2.70E-02 1.61E-02 8.93E-03 5.74E-03 5.74E-03 1.96E-03 1.96E-03 1.96E-04 3.07E-04 1.49E-04 1.49E-04 1.91E-06

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295/11	05	=	°°∩	' 1I	9	=	х	18	erisade	IEnolender - Snorthered	uo	nara	C.C SIGBI

30-738 1	5 005+02	90-51 I	50+773.2	90-302 1	50+302.a	8.965-06	7.015+02	90-301 1	20+305 8	The Prove	C 065+13
50-366'1	4, 195+02	2.30E-05	4.41E+02	2.06E-05	5, 06E+02	S0-306'S	50+370.8	50-307 1	20+389*9	90-311.9	20+371"1
1.295-04	20+328.5	P0-35P'1	5, 15E+02	PO-356 1	5,61E+02	1.255-04	4. 34E+02	PO-310.1	20+3/1.8	S0-308'S	20+365'5
2.925-04	20+326.1	20-321.2	2.14E+02	3.08E-04	2.46E+02	2.94E-04	2.955+02	2.74E-04	S0+365.8	1.652-04	20+39L'S
P0-381 P	1.362+02	4.31E-04	1.515+02	V.26E-04	1. 75E+02	PO-321 P	2.06E+02	2.87E-04	2.295+02	5.946-04	Z0+359"Z
PO-395'L	10+350'6	PO-348'L	1.01E+02	1.322-04	1,165+02	PO-361'L	1.395+02	70-385'9	20+365"1	10-361 9	20+311-1
\$0-397'1	10+399'5	20-387'1	10+306 '9	50-3/5'1	10+3/111	£0-381'1	10+396'8	1.285-05	10+329'6	1.205-05	20+311":
S0-311.5	10+376'5	2.11E-03	10+316"7	\$0-326'1	10+371'5	£0-358'1	10+399'9	\$0-39/ 1	10+379"1	\$0-359"1	10+317 8
50-311.S	10+350'5	2'406-02	10+301 '5	2.91E-03	10+358'5	2.085-05	10+396'5	\$0-385'Z	10+395'5	50-311"Z	10+360 '9
\$0-305'P	10+362'2	£0-386'S	10+321 2	\$0-386'P	10+3/9'2	\$0-396'5	10+361 '5	S0-367'S	10+310"P	\$0-36G'S	10+366'6
C0-316'9	10+3/6'1	81226-02	10+326'1	50-3/2'9	10+396'7	61-315-02	10+360"7	\$0-39G'S	10+7/5"5	50-369'7	10+355'7
20-361-1	10+701 1	20-380'1	10+312-1	\$0-366'/	10+76/1	60-365'8	10+39'1	\$0-395 7	10+325-2	S0-30/ 'S	10+386 7
20-321.2	00+385'/	20-399'1	8.845+00	CO-36/ '6	10+316'1	60-382'6	10+3/2-1	\$0-351 7	10+799*1	60-369"/	10+70C*1
20-386'2	00+300'9	20-300'7	00+300.0	20-308'1	00+311*/	20-3/111	00+789'6	CO-72C'8	10+356*1	C0-399'6	10.2011
20-305'C	00+312"2	20-301 'C	00+318'9	20-3/5'2	00+307.0	20-399'1	00+325.0	20-361-1	00+316'6	20-360-1	00+369'9
20-20/10	00+387'C	20-328'C	00+3/8'6	20-399.2	00+368'C	20-320.2	00+39/ "7	20-322'1	00+389"	20-3601	004318-0
20-376 1	00+300'2	20-320.8	00+317'7	20-360'C	00+300'6	20-301-2	00+302 1	20-306'1	00+301 *2	20-308 1	00-300'5
20-302'5	00+304 2	20-307'C	00+300*1	20-300 1	00+386'1	20-301 0	00+310*2	20-381 *7	00+316-7	20-330 1	00+305 7
20-396 7	10-3616	20-396 9	10-3/9'6	20-328 8	00+300 1	20-301 2	3 012 +00	20-302 6	00+319 2	20-300 2	2 045400
20-310 2	10-300	20 30 31 7	10-329 0	20 31 717	10 301 1	20-300 2	00-31011	20 3.017	00+358 1	20 393 1	1 705+00
2 915-02	10-322 9	CV-38C C	10-367 2	2 215-02	10-301 0	C0-38.9 1	00+310 1	20-368 C	00+361 1	1 205-02	1 206.+00
(1.4)	(14/1)	(11)	(14/1)	(11)	(14/1)	(E1)	(13'1)	(14)	(14.1)	(1.1)	(14.1)
EIKIVO	x	E(K) 100	×	E (K) VOO	×	E (K) VO	×	E ((C) ' 00	×	ECCIVON	×
		100 - COLERA		1777 C 1177		0 <del>000</del> 6. V - 22				1770 C 1970 C 1970	
NI 871	*Z=Z	N1 781	1=Z	NI 129	• •Z	NI PPZ	-2	NI LSI	· •2	NI 120	- =Z
NI S	"Z=1	N1 9	. S=Y	NI S	*=X	NI S	.S=Y	NI 9	"Z=1	NI S	*Z=1

5:00E-09 1:44E-02 2:24E-04 3:25E-02 2:24E-04 2:24E-04 2:24E-04 2:25E-02 2:05E-	2:55:405 4:51:405 2:51:405 2:32:41:405 2:32:82:405 2:52:82:405 3:32:40 1:22:40 3:32:40 2:32:40 3:32:40 2:32:40 3:32:40 2:32:40 3:32:50 3:32:40	1:93E-06 5:36E-02 2:12E-04 2:11E-04 1:18E-03 1:18E-03 2:06E-02 2:06E-02 2:06E-02 2:06E-02 2:06E-03 2:0	2:132:405 4:462:405 2:132:405 1:252:405 1:252:405 1:252:405 2:132:405	1 256-09 1 3 256-09 1 586-09 1 586-09 1 576-09 1 576-09 1 576-09 1 586-02 2 386-02 2 386-02 2 986-02 2 986-02 2 986-02 2 986-05 2 98	e'.12E+05 3'5'1E+05 3'5'1E+05 3'5'1E+05 1'5'0E+05 1'5'0E+05 3'6'2E+01 3'6'2E+01 3'6'2E+01 3'6'2E+01 1'5'2E+01 1'5'2E+00 1'6'2E+00 1'6'2E+00 1'5'2E+00 3'5'4'E+00 3'5'5'5'5'5'5'5'5'5'5'5'5'5'5'5'5'5'5'5	1: 226-00 1: 286-02 1: 286-02 2: 226-04 2: 226-04 2: 266-02 2: 106-03 2: 286-02 2: 286-02 2: 286-02 2: 286-02 3: 286-02 1: 286-02 1: 186-05 1: 186-05 3: 486-02 3: 486-02	8'.48:405 9''''''''''''''''''''''''''''''''''''	8'02E-02 1'02E-02 3'48E-02 3'48E-04 2'52E-03 1'51E-03 2'48E-03 2'48E-03 2'48E-03 2'48E-03 2'48E-03 2'48E-03 2'48E-03 2'48E-03 1'02E-05 1'0	1'12E+02 9'24E+02 9'24E+05 9'24E+05 1'26E+05 9'46E+01 1'26E+05 9'46E+01 9'45E+00 9'46E+01 1'24E+01 1'24E+01 9'46E+00 1'24E+00 9'46+00 9'46+00 9'40+00000000000000000000000000000000000	2:55E-02 e:08E-02 e:08E-02 e:08E-02 1:11E-02 2:55E-03 1:11E-02 2:02E-03 2:02E-	1,296+03 1,206+03 1,006+03 1,006+02 2,206+02 2,206+05 1,126+05 2,206+01 1,126+05 2,206+01 2,206+01 2,206+01 2,206+01 1,246+00 2,206+00 2,000+000+000+000+000+000+000+000+000+00
(FT) E(K) /0U	(1.7FT) K	E(K) /00	(13/1) K	E (K) /UU	(13/1) K	E (K) . (U	(13/1) K	E (K) UU (FT)	(1/11) K	E (IC) . 100 (FT)	(13/1) K
256 IN 0 IN	Z=Z	NI 199	*1=Z	NI 999		NI 262 NI 0		NI SZI		NI 080	-=2 -=2

90-351 8	20+300'9	90-366'1	70+3C1 'C	00-2017	70+201 '9	90-366'1	20+32/ 6	10-366'9	\$0+35Z'.
60-307'C	20+329*0	CO-3/8'7	20-310'5	50-320 2	20+312.0	C0-318'1	20+386"/	90-39/ 1	20+3G/ '6
30-30C E	20.305 2	30-327 5	20.20.00	30-368 6	20+316 9	30-30011	20+318 6	60-311'1	20+396'9
1 032-04	2 805 403	VU-35V 1	1 10L+02	1 485-04	20+3CL 1	70-390 1	20.300'C	30 306'1	20+301 **
PO-JCL 1	2.165+02	5.00F-04	20+371.5	5. 14E-04	20+77.5	P0-369 C	20+303 F	1 005-04	CVYJEL V
A.01E-04	1.295+02	4.425-04	S0+362.1	\$0-366 °E	20+367.1	20-327.2	2.605+02	2 905-04	20+375 E
P0-320'8	10+327'8	P0-362'L	20+320.1	PO-361 '9	1, 196+02	70-319.8	20+3E1.1	P0-30E-9	20+3/1.1
1.295-03	10+396'9	1.216-05	10+317'9	£0-315'1	10+365'1	20-302 1	1.11E+02	P0-365 6	20+324.1
50-38/ 1	10+315'7	\$0-38/ 1	10+3/9'P	50-380'2	10+390'5	50-321-1	10+329'8	E0-3:2"1	20+32C-1
£0-30G'S	10+368'2	\$0-761 'S	10+3/1'5	2.632-05	10+326'6	50-369'2	10+365'5	2.27E-05	10+317"2
60-371 9	10+316'1	C0-309'C	10+3/0.7	C0-70C'C	10+366.7	SO-315.63	10+305 *7	2'CIE-02	10+319"7
C0-306'8	10+306'1	C0-3/6'8	10+7/6'1	CO-36C'8	10+388'1	6. 362-03	10+36/ 7	\$0-310"P	10+37/ "5
20-382'1	10-301 1	20-30111	10-787'1	20-302 1	10+39911	CO-396'/	10+791 7	CO-366'P	10+358-7
20-306 1	00+308'8	20-3/01	00+320'6	20-336 1	10+390.1	C0-30C'/	10+799"1	CO-398'/	10+39/11
20 30 1	00+300 0	20 309 1	00-322-00	20-335 1	00.306'1	CO-305'	10+318"1	CO-368'	10+362"
20-362 2	00+322 9	1 802-03	00+329 9	CV-3C8 1	1 014 310 1	20-300 2	10-301-0	CO. 311 '0	10-317-1
50-768.2	00+361.5	20-3/1.5	00+35/ P	20-386 2	00+390 %	20-398 1	00+371 8	20-311 0	104312 1
50-358.b	5.14E+00	5. 305-02	2.945+00	5. 445-02	00+382 5	20-389 1	00+308 5	20-302 0	8 245+00
4.93E-02	2.225+00	3, 996-02	2.25E+00	3.296-02	2.62E+00	20-311-1	00+302.5	1 345-02	5.265+00
4.255-02	00+385'1	5.50E-02	1.662+00	2,465-02	1.765+00	20-365.1	2.59E+00	E0-317.0	00+359°S
Z0-340'P	00+360'1	20-366°Z	10-305'6	20-357'5	1.285+00	1.61E-02	1.845+00	S0-320'8	2.41E+00
4.96E-02	10-327 9	20-3/1.5	10-379'L	1.945-02	10-372.8	20-351.1	1.275+00	20-388.7	00+359.1
11-11	11 4/11	() ()	114/11	11.43	11 4/11	(1.4)	(14/1)	(1.4)	(14/1)
1131	113/11	00(11)3	123/11	0011313	123.11	EIK1100	14.4	M	
TTP/ (A) 3		TTA TAT 3	~	110 1213	<u> </u>	(1) (M) 1	~	110 100 3	-
NI SSZ	"7=7	NI 0/2	"1=Z	NI 265	• =Z	NI ZPZ	* =Z	NI SEI	• =Z
NI	"Cal	NI	'Cel	NI G	"S=1	NI 9	"S=1	N1 9	*S=1

Table 5.5 Data on one-dimensional spectra at x = 6 ft,  $\rm U_{\infty}$  = 30 ft/sec -- Continued

Y=4	0 IN	Y=4.	0 IN	Y=4.	0 IN	Y=4.1	0 IN	1=4.0	IN
Z=	135 IN	Z= .	242 IN	Z= .	692 IN	Z=1.2	268 IN	Z=2.2	52 IN
K	E (K) /00	к	E(K)/00	к	E (K) /00	K	E(K) /W	K	E(K)/W
(1,FT)	(FT)	(1.FT)	(FT)	(1/FT)	(FT)	(1/FT)		(1/FT)	(FT)
1.69E+00 2.45E+00 4.52E+00 5.99E+00 7.81E+00 1.01E+01 1.49E+01 1.49E+01 2.44E+01 3.83E+01 7.70E+01 1.11E+02 2.26E+02 3.40E+02 2.26E+02 3.40E+02 7.09E+02 9.91E+02 1.27E+05	8.51E-03 1.29E-02 7.04E-03 8.45E-03 8.45E-03 6.16E-03 5.39E-03 4.55E-03 3.56E-03 2.51E-03 1.92E-03 1.92E-03 1.92E-03 5.57E-04 2.87E-04 2.87E-04 1.82E-04 5.75E-05 5.12E-06 4.77E-07	1.41E+00 1.96E+00 5.95E+00 5.92E+00 9.04E+00 1.55E+01 2.12E+01 3.96E+01 5.55E+01 9.56E+01 1.12E+02 1.75E+02 2.56E+02 5.40E+02 7.57E+02 9.73E+02	6.76E-03 1.99E-02 1.64E-02 1.04E-02 1.01E-02 1.01E-02 1.01E-02 1.01E-02 1.01E-02 1.01E-03 5.94E-03 5.94E-03 1.65E-03 1.15E-03 1.15E-03 6.27E-04 3.59E-04 9.25E-05 9.40E-06 7.50E-07	8.90E-01 1.26E+00 2.69E+00 2.69E+00 5.79E+00 8.09E+00 1.04E+01 1.36E+01 2.35E+01 4.25E+01 6.05E+01 1.21E+02 1.21E+02 1.21E+02 3.77E+02 5.20E+02 6.79E+02	3. 72E-02 4. 04E-02 4. 14E-02 4. 14E-02 2. 27E-02 2. 27E-02 1. 40E-02 1. 40E-02 1. 40E-02 5. 39E-03 5. 74E-03 2. 99E-03 1. 26E-03 6. 87E-04 4. 33E-04 4. 33E-04 1. 26E-04 1. 26E-04 1. 34E-06	7.99E-01 1.10E+00 2.39E+00 3.01E+00 9.17E+00 9.17E+00 1.14E+01 1.67E+01 5.21E+01 5.21E+01 5.21E+01 5.21E+01 1.14E+02 1.52E+02 2.15E+02 2.15E+02 3.16E+02 4.43E+02 5.69E+02	5.70E-02 5.27E-02 5.69E-02 3.69E-02 1.80E-02 1.80E-02 8.61E-03 5.31E-03 5.31E-03 3.17E-03 2.40E-03 1.64E-03 1.64E-03 5.73E-04 4.05E-04 1.40E-04 1.40E-04 1.59E-06	6.66E-01 9.60E-01 1.51E+00 3.52E+00 3.52E+00 5.84E+00 9.60E+00 1.21E+01 1.66E+01 1.91E+01 2.93E+01 4.10E+01 9.64E+01 1.40E+02 2.20E+02 2.96E+02 3.76E+02 5.12E+02	3.14E-02 4.73E-02 6.42E-02 4.11E-02 3.59E-02 3.59E-02 2.70E-02 1.24E-02 1.24E-03 6.41E-03 5.60E-03 5.60E-03 5.60E-03 2.22E-03 1.24E-03 7.26E-04 9.65E-05 2.41E-05 2.41E-05
Y=5	.0 IN	Y=5.	0 IN	7=5.	0 IN	ĭ=5.	0 IN	Y=5.0	IN
Z=	.135 IN	Z= .	248 IN	Z= .	692 IN	Z=1.,	266 IN	Z=2.2	29 IN
K	E(K)/W	K	E (K) /W	K	E (K) /W	(1/FT)	E (K) /W	K	E(K)/W
(1/FT)	(FT)	(1/FT)	(FT)	(1/FT)	(FT)		(FT)	(1/FT)	(FT)
1.76E+00 2.80E+00 5.96E+00 5.56E+00 1.15E+00 1.15E+01 1.57E+01 2.99E+01 4.49E+01 5.70E+01 1.16E+02 1.51E+02 2.28E+02 3.42E+02 4.85E+02 7.15E+02 9.98E+02 1.29E+03	6.04E-03 8.84E-03 1.01E-02 1.02E-02 9.30E-03 9.71E-03 9.71E-03 9.31E-03 7.42E-03 5.13E-03 2.33E-03 1.38E-03 1.38E-03 1.38E-03 1.38E-03 1.38E-03 1.38E-03 1.0E-05 5.50E-04 4.95E-05 4.58E-06 4.22E-07	1.51E+00 1.75E+00 2.80E+00 4.26E+00 6.45E+00 1.10E+01 2.01E+01 2.01E+01 2.02E+01 3.92E+01 3.92E+01 1.15E+02 2.60E+02 3.60E+02 3.60E+02 9.74E+02 9.74E+02	1.62E-02 2.43E-02 2.15E-02 1.69E-02 1.64E-02 1.21E-02 1.21E-02 1.21E-02 1.24E-02 2.30E-03 1.48E-03 1.48E-03 1.48E-03 1.48E-03 1.48E-03 1.59E-04 2.36E-04 2.56E-04 2.56E-04 7.81E-05 7.62E-06 7.11E-07	9.29E-01 1.33E+00 2.03E+00 4.39E+00 4.39E+00 8.00E+00 1.42E+01 1.42E+01 1.42E+01 1.49E+01 3.66E+01 8.27E+01 9.27E+01 9.27E+01 1.39E+02 1.39E+02 2.56E+02 5.26E+02 6.76E+02	2.10E-02 4.10E-02 3.40E-02 3.62E-02 2.59E-02 2.59E-02 2.59E-02 1.23E-02 1.23E-02 9.98E-03 5.23E-03 5.23E-03 5.25E-03 5.25E-03 7.78E-04 4.10E-04 4.10E-04 4.10E-04 1.20E-04 1.20E-04 1.20E-04	7.49E-01 1.09E+00 2.67E+00 3.69E+00 6.99E+00 6.99E+00 9.76E+00 1.19E+01 1.69E+01 2.49E+01 3.27E+01 4.89E+01 9.19E+01 9.19E+01 2.72E+02 4.44E+02 5.71E+02	5.15E-02 5.75E-02 5.65E-02 4.34E-02 4.34E-02 2.58E-02 1.66E-02 2.58E-02 1.66E-02 8.06E-03 5.65E-03 5.65E-03 5.65E-03 5.05E-03 5.05E-03 8.75E-04 5.85E-04 1.53E-06 1.53E-06	6.44E-01 9.30E-01 1.30E+00 2.20E+00 3.56E+00 6.24E+00 6.24E+00 7.89E+01 1.50E+01 1.50E+01 1.50E+01 4.27E+01 4.27E+01 4.27E+01 1.00E+02 1.77E+02 2.59E+02 2.59E+02 3.69E+02 4.92E+02	5.36E-02 6.10E-02 5.32E-02 4.49E-02 5.32E-02 2.93E-02 2.93E-02 2.93E-02 2.22E-02 1.64E-02 8.57E-03 5.31E-03 5.31E-03 5.31E-03 5.31E-03 5.31E-03 5.34E-04 2.98E-04 2.98E-04 1.17E-04 1.44E-05 2.13E-06
Y=6.	0 IN	Y=6.	0 IN	Y=6	0 1N	Y=6	.0 IN	Y=6	.0 IN
Z= .	135 IN	Z= .	242 IN	Z=	692 1N	Z=1	.274 IN	Z=2	2.244 IN
(1/FT)	E (K) . 00	K	E(K)/00	K	E(K).00	K	E (K) /00	K	E(K)/00
	(FT)	(1/FT)	(FT)	(1.FT)	(FT)	(1/FT)	(FT)	(1/FT)	(FT)
1.71E+00 2.61E+00 3.80E+00 5.67E+00 7.85E+00 1.05E+01 1.55E+01 2.71E+01 4.01E+01 4.01E+01 4.82E+01 8.02E+01 1.10E+02 1.55E+02 2.31E+02 2.31E+02 2.31E+02 2.31E+02 1.55E+02 4.91E+02 7.22E+02 1.01E+03 1.30E+03	1.00E-02 9.40E-03 9.97E-03 9.39E-02 9.39E-03 6.02E-03 3.75E-03 3.62E-03 2.50E-03 2.50E-03 1.27E-03 1.26E-03 6.28E-04 5.99E-04 5.99E-04 5.99E-05 5.47E-07	1.34E+00 1.93E+00 5.10E+00 4.32E+00 6.41E+00 8.59E+00 8.59E+01 1.66E+01 1.56E+01 1.56E+01 1.46E+01 1.46E+01 1.46E+01 1.46E+01 1.46E+01 1.46E+01 1.78E+02 3.77E+02 5.57E+02 5.57E+02 1.00E+03	9.62E-03 2.96E-02 2.55E-02 2.55E-02 1.27E-02 1.20E-02 8.50E-05 6.01E-05 5.50E-05 2.88E-05 2.88E-05 2.88E-05 1.55E-05 1.55E-05 1.55E-05 1.15E-05 8.40E-04 9.15E-05 8.90E-07	8.852-01 1.562+00 2.112+00 4.562+00 5.772+00 8.042+00 1.102+01 1.552+01 2.022+01 2.022+01 7.922+01 1.252+02 2.652+02 3.052+02 5.452+02 7.012+02	1.802-02 2.922-02 3.742-02 3.812-02 3.812-02 2.662-02 1.192-02 8.132-03 3.552-03 2.662-03 2.662-03 3.552-03 1.512-03 1.512-03 1.512-03 1.512-03 1.512-04 1.812-04 1.812-05 1.532-06	8.03E-01 1.18E+00 1.79E+00 3.66E+00 5.02E+00 1.26E+01 1.85E+01 2.37E+01 3.45E+01 1.65E+02 2.22E+02 3.27E+02 4.57E+02 4.57E+02 4.57E+02 4.57E+02 4.57E+02 4.57E+02 4.57E+02 4.57E+02 4.57E+02 4.57E+02 4.57E+02 4.57E+02 4.57E+02 4.57E+02 4.57E+02 4.57E+02 5.88E+	1.66E-02 2.49E-02 3.00E-02 2.23E-02 2.97E-02 2.97E-02 2.91E-02 1.91E-02 1.91E-02 1.91E-02 1.91E-02 1.91E-02 1.95E-03 1.95E-03 1.95E-04 4.55E-04 1.55E-04 1.55E-04 2.45E-04 2.45E-04 2.45E-04	7.16E-01 1.10E+00 2.39E+00 3.55E+00 5.76E+00 8.29E+00 1.15E+01 1.42E+01 1.42E+01 1.42E+01 1.42E+01 1.62E+02 2.05E+02 2.05E+02 2.05E+02 5.12E+02	3.56E-02 4.202-02 4.502-02 5.56E-02 5.16E-02 2.76E-02 2.76E-02 2.76E-02 2.76E-02 2.75E-02 9.0TE-03 5.57E-05 9.61E-04 6.32E-03 9.61E-04 6.32E-04 4.2.97E-04 1.45E-04 1.45E-04

Table 5.5 Data on one-dimensional spectra at x = 10 ft,  $U_{\infty}$  = 30 ft/sec

Y=0.1	0 IN	Y=0.0	) IN	Y=0.0	IN	Y=0.	0 IN	Y=0.	0 IN
Z=	0355 IN	Z= .2	274 IN	Z= .5	191 IN	Z=1.	119 IN	Z=2.	175 IN
K	E (K) . 00	K	E (K) . W	K	E (K) . W	K	E (K) /00	K	E(K)/00
(1. FT)	(FT)	(1./FT)	(FT)	(1.FT)	(FT)	(1/FT)	(FT)	(1/FT)	(FT)
1.60E+00 2.53E+00 3.46E+00 4.93E+00 7.42E+00 1.01E+01 1.53E+01 1.53E+01 4.53E+01 4.53E+01 4.53E+01 1.01E+02 1.40E+02 1.89E+02 3.41E+02 4.67E+02 6.73E+02 9.33E+02 9.33E+02 1.17E+03	8.12E-03 1.66E-02 1.66E-02 1.46E-02 1.46E-02 8.90E-03 7.39E-03 8.70E-03 4.91E-03 5.40E-03 5.40E-03 7.24E-04 5.71E-04 5.71E-04 5.41E-05 1.36E-04 5.41E-05 1.75E-06 1.40E-06	1.01E+00 1.52E+00 2.25E+00 3.31E+00 6.89E+00 9.72E+00 1.27E+01 1.61E+01 2.34E+01 3.14E+01 4.96E+01 6.31E+01 8.87E+01 1.42E+02 2.13E+02 3.02E+02 4.43E+02 7.99E+02	2.74E-02 3.63E-02 3.64E-02 2.62E-02 1.70E-02 1.70E-02 1.44E-02 8.52E-03 3.23E-03 1.50E-03 1.50E-03 1.52E-04 2.92E-04 1.99E-04 2.30E-05 1.81E-06	9.29E-01 1.39E*00 2.21E*00 3.06E*00 5.97E*00 8.12E*00 1.16E*01 1.42E*01 2.02E*01 2.02E*01 2.02E*01 3.35E*01 7.60E*01 1.95E*02 2.62E*02 3.92E*02 6.92E*02	5.04E-02 4.45E-02 5.02E-02 5.02E-02 1.99E-02 1.99E-02 1.96E-02 1.96E-03 5.12E-03 5.12E-03 1.60E-03 1.60E-03 1.96E-03 2.90E-04 2.90E-04 2.29E-04 2.29E-05 2.05E-06	7.85E-01 1.25E+00 2.70E+00 5.85E+00 5.26E+00 6.79E+00 9.75E+00 1.52E+01 1.81E+01 2.54E+01 1.515E+01 7.12E+01 1.10E+02 2.33E+02 3.42E+02 6.16E+02	5.24E-02 6.60E-02 5.18E-02 5.77E-02 2.82E-02 2.82E-02 1.32E-02 1.32E-02 1.32E-02 2.82E-03 4.04E-03 9.94E-03 9.94E-04 5.22E-04 5.22E-04 1.09E-04 1.09E-04 1.09E-04 1.76E-06	6.44E-01 9.97E-01 1.43E+00 2.39E+00 5.43E+00 5.91E+00 7.69E+01 1.60E+01 2.96E+01 4.13E+01 6.30E+01 9.01E+01 1.26E+02 1.72E+02 2.50E+02 2.50E+02 5.57E+02	6.10E-02 5.90E-02 4.94E-02 4.20E-02 2.68E-02 1.93E-02 1.93E-02 7.62E-03 5.14E-03 1.67E-03 1.67E-03 1.67E-03 1.63E-03 6.32E-04 2.23E-04 1.52E-04 1.52E-04 2.29E-06
Y=1.	0 IN	Y=1.	0 1N	r=1.(	0 IN	Υ=1.	0 IN	Y=1.	0 IN
Z= .	035 IN	Z=	276 1N	2= .9	591 IN	Z=1.	119 IN	Z=2.	177 IN
K	E(K).00	K	E (K), W	к	E (K) /W	K	E (K) /W	K	E(K)/W
(1.FT)	(FT)	(1.FT)	(FT)	(1.FT)	(FT)	(1/FT)	(FT)	(1/FT)	(FT)
1.50E+00 2.15E+00 5.60E+00 5.60E+00 1.09E+01 1.39E+01 1.39E+01 1.81E+01 2.55E+01 2.55E+01 2.55E+01 2.55E+01 4.81E+01 1.51E+02 2.04E+02 6.46E+02 6.46E+02 5.68E+02 5.68E+02	1.36E-02 2.10E-02 2.25E-02 1.93E-02 1.93E-02 1.01E-02 9.24E-03 5.52E-03 5.52E-03 5.52E-03 5.52E-03 9.49E-04 4.04E-03 9.49E-04 4.04E-04 1.30E-04 4.47E-05 1.92E-06 1.92E-06 1.92E-06	1.12E+00 1.60E+00 5.69E+00 5.19E+00 9.64E+00 1.33E+01 2.51E+01 3.25E+01 4.71E+01 6.79E+01 9.50E+01 1.46E+02 2.19E+02 3.10E+02 4.56E+02 9.22E+02 9.22E+02	5.08E-02 3.95E-02 1.92E-02 1.92E-02 1.21E-02 8.25E-03 5.28E-03 5.28E-03 5.28E-03 5.28E-03 5.28E-03 5.28E-03 5.28E-03 9.77E-03 9.77E-04 1.97E-04 9.41E-05 1.87E-05 1.45E-06	9.71E-01 1.395E+00 3.05E+00 4.12E+00 9.09E+00 1.32E+01 2.16E+01 2.16E+01 2.16E+01 2.16E+01 2.16E+01 1.39E+02 1.26E+02 1.99E+02 2.67E+02 3.93E+02 3.93E+02 2.55EE+02 7.07E+02	5.45C-02 3.31E-02 6.01E-02 4.59E-02 3.35E-02 2.27E-02 1.46E-02 9.60E-03 5.05E-03 3.76E-03 3.76E-03 1.42E-03 1.42E-03 1.42E-03 1.42E-03 1.42E-03 1.93E-04 2.91E-04 1.93E-04 9.74E-05 1.73E-06	9.66E-01 1.61E+00 2.21E+00 3.35E+00 4.21E+00 6.17E+00 7.56E+00 9.99E+00 1.42E+01 1.42E+01 3.27E+01 3.27E+01 3.27E+01 3.27E+01 3.27E+02 4.84E+02 6.23E+02	4,71E-02 5,36E-02 4,06E-02 5,24E-02 2,35E-02 1,95E-02 1,95E-02 1,95E-02 8,56E-03 4,67E-03 5,02E-03 1,63E-03 1,63E-03 1,06E-03 2,95E-04 2,95E-04 2,95E-04 2,01E-04 1,12E-04 2,06E-05 2,06E-06	7.43E-01 1.00E+00 1.63E+00 2.28E+00 3.02E+00 5.97E+00 8.35E+00 1.25E+01 1.25E+01 1.25E+01 1.25E+01 1.25E+01 1.62E+01 1.05E+02 1.67E+02 2.45E+02 3.52E+02 3.52E+02 5.35E+02	2.94E-02 6.57E-02 5.65E-02 5.65E-02 5.70E-02 5.77E-02 2.64E-02 1.02E-02 1.02E-02 1.02E-03 2.16E-03 5.75E-03 2.16E-03 8.85E-04 4.46E-04 7.09E-04 1.72E-04 7.05E-05 2.19E-06
Y=2	.0 IN	Y=2	.0 IN	Y=2	.0 IN	Y=2	0 IN	Y=2.	0 IN
2=	.035 IN	Z=	.276 IN	Z=	591 IN	Z=1	119 IN	Z=2.	175 IN
(1.FT)	E (K) . 100 (FT)	(1.FT)	E(K). W (FT)	(1.FT)	E (K) . W (FT)	(1./FT)	E(K)/00 (FT)	K (1/FT)	E(K)/W (FT)
1.69E+00 2.20E+00 5.19E+00 6.82E+00 1.50E+01 1.95E+01 2.41E+00 3.45E+0 4.76E+0 8.25E+00 1.21E+02 1.77E+02 2.45E+02 1.77E+02 2.45E+02 4.47E+02 6.41E+02 8.92E+00 1.14E+02	1.71E-02 1.77E-02 1.77E-02 1.77E-02 1.76E-02 1.76E-02 1.95E-02 1.95E-02 1.95E-02 1.95E-03 1.22E-02 1.22E-02 1.22E-02 1.23E-03 2.101E-03 2.101E-03 2.101E-03 2.101E-04 2.101E-05 2.25E-06 5.225E-06 5.225E-06	1,04E+00 1,59E+00 2,62E+00 2,62E+00 5,67E+00 1,04E+01 1,37E+01 2,83E+01 5,349E+01 5,349E+01 1,47E+02 2,21E+02 5,349E+01 1,47E+02 2,21E+02 5,313E+02 5,313E+02 5,48E+02 5,313E+02 5,312E+02 5,312E+01 5,312E+02 5,	1.51E-02 4.54E-02 4.54E-02 2.85E-02 2.25E-02 1.85E-02 9.60E-03 5.62E-03 5.56E-03 5.56E-03 5.56E-03 5.56E-03 1.59E-03 1.59E-03 1.59E-03 1.59E-04 2.25E-04 2.25E-04 2.25E-04 2.25E-04 1.53E-05 1.53E-05 1.24E-06	9.45E-01 1.56E+00 2.23E+00 3.24E+00 5.96E+00 8.96E+00 1.11E+01 1.56E+01 2.75E+01 4.29E+01 1.26E+02 2.69E+02 3.94E+02 5.51E+02 5.51E+02 5.94E+02	3.27E-02 5.31E-02 4.19E-02 3.17E-02 2.44E-02 1.49E-02 1.49E-02 9.24E-03 5.99E-03 4.34E-03 2.40E-03 1.51E-03 9.71E-04 5.79E-04 5.22E-04 2.19E-04 1.99E-04 1.99E-04 1.99E-04	8.24E-01 1.21E+00 1.77E+00 2.91E+00 4.54E+00 5.65E+00 1.29E+01 2.47E+01 5.29E+01 1.29E+01 1.29E+01 6.87E+01 1.0E+02 2.54E+02 3.45E+0	6.75E-02 6.46E-02 5.46E-02 5.36E-02 3.36E-02 2.76E-02 2.20E-02 1.46E-02 9.76E-03 4.06E-03 1.61E-03 1.61E-03 1.61E-03 5.82E-04 3.56E-04 2.29E-04 1.10E-04 2.29E-04 1.55E-06	6.92E-01 9.00E-01 1.55E+00 3.54E+00 6.14E+00 8.66E+00 1.10E+01 1.47E+01 2.00E+01 1.47E+01 6.16E+01 1.02E+02 1.59E+02 3.59E+02 3.59E+02 4.12E+02	5.90E-02 7.50E-02 6.32E-02 4.95E-02 4.95E-02 2.91E-02 1.66E-02 1.66E-02 1.66E-03 5.52E-03 5.52E-03 5.52E-03 5.52E-03 1.75E-04 2.87E-04 2.87E-04 2.56E-04 7.55E-04 7.55E-05 2.17E-05 2.17E-05

Table 5.5 Data on one-dimensional spectra at x = 10 ft,  $\rm U_{\infty}$  = 30 ft/sec  $\rm --$  Continued

Y=2.5	5 IN	Y=2.	5 IN	r=2.	5 IN	Y=2.	5 IN	Y=2.	5 IN
Z= .0	0575 IN	2=	276 IN	Z= .	591 IN	Z=1.	119 IN	Z=2.	175 IN
K	E(K).00	K	E (K) . W	K	E (K) / 00	к	E(K)/00	K	E(K) /W
(1.FT)	(FT)	(1.FT)	(FT)	(1.FT)	(FT)	(1/FT)	(FT)	(1/FT)	(FT)
1.58E+00 2.55E+00 3.68E+00 7.04E+00 7.04E+00 1.01E+01 1.39E+01 2.44E+01 3.29E+01 4.59E+01 5.91E+01 9.74E+01 1.35E+02 2.09E+02 5.06E+02 6.47E+02 8.87E+02 8.87E+02	6.29E-03 1.09E-02 1.70E-02 1.72E-02 1.47E-02 1.47E-02 9.36E-03 6.70E-03 4.91E-03 6.70E-03 4.91E-03 6.70E-03 2.29E-03 2.29E-03 2.29E-04 5.04E-04 4.39E-05 2.56E-06 5.76E-07 1.55E-07	1.29E+00 1.83E+00 2.93E+00 3.96E+00 5.62E+00 7.31E+00 9.02E+00 1.38E+01 1.72E+01 2.74E+01 2.74E+01 9.39E+01 1.54E+02 2.31E+02 4.81E+02 6.73E+02 8.66E+02	2.45E-02 3.97E-02 2.97E-02 2.97E-02 2.97E-02 1.06E-02 1.06E-02 1.06E-03 8.72E-03 5.82E-03 3.57E-03 2.57E-03 1.512E-03 1.512E-03 5.48E-04 2.96E-04 2.01E-04 1.03E-04 2.03E-05 1.64E-06	9.14E-01 1.24E+00 2.52E+00 5.46E+00 4.20E+00 9.40E+00 9.40E+00 1.07E+01 1.60E+01 2.78E+01 6.49E+01 6.49E+01 1.28E+02 1.92E+02 2.73E+02 4.01E+02 5.61E+02 7.22E+02	5.84E-02 5.34E-02 4.51E-02 5.43E-02 2.99E-02 1.48E-02 1.36E-02 8.79E-03 6.36E-03 5.91E-03 2.67E-03 1.42E-03 1.42E-03 5.68E-04 3.12E-04 2.10E-04 1.09E-04 2.20E-05 1.72E-06	8.74E-01 1.17E+00 1.81E+00 2.64E+00 3.63E+00 6.65E+00 9.73E+00 1.34E+01 1.70E+01 2.52E+01 3.66E+01 4.81E+01 1.10E+02 1.65E+02 3.44E+02 4.81E+02 6.19E+02	5.01E-02 6.55E-02 4.95E-02 2.55E-02 2.09E-02 2.09E-02 2.09E-02 1.47E-02 1.47E-02 1.47E-02 1.50E-03 2.70E-03 1.76E-03 1.39E-03 6.29E-04 3.40E-04 2.26E-04 1.24E-04 2.84E-05 2.21E-06	6.96E-01 1.05E+00 2.31E+00 3.36E+00 6.42E+00 6.42E+00 1.16E+01 1.68E+01 1.68E+01 3.17E+01 4.55E+01 3.17E+01 4.55E+02 3.59E+02 4.50E+02 5.45E+02	5.49E-02 6.49E-02 7.06E-02 7.02E-02 3.57E-02 2.56E-02 1.77E-02 1.09E-02 6.95E-03 3.24E-03 2.02E-03 3.24E-05 2.02E-03 6.21E-04 2.79E-04 1.74E-04 7.17E-05 1.33E-05 3.21E-06
Y=3.	0 IN	r=3.	0 IN	Y=3.	0 IN	Y=3.	0 IN	Y=3.	0 IN
2= .	105 IN	Z= .	274 IN	2= .	591 IN	Z=1.	119 IN	Z=2.	175 IN
K (1.FT)	E(K)/W (FT)	(1/FT)	E (K), W (FT)	(1.FT)	E (K) /W (FT)	к (1./FT)	E (K) /W (FT)	K (1/FT)	E(K)/W (FT)
1.72E+00 2.52E+00 3.52E+00 7.87E+00 1.12E+01 1.52E+01 1.52E+01 1.52E+01 3.70E+01 6.35E+01 1.13E+01 1.57E+02 2.36E+02 3.46E+02 5.21E+02 7.11E+02 9.39E+02 1.22E+05	4.47E-03 9.45E-03 1.37E-02 9.05E-03 1.15E-02 9.05E-03 4.20E-03 4.20E-03 2.55E-03 1.40E-04 1.78E-04 1.78E-04 1.78E-04 3.20E-05 3.20E-05 3.20E-05 3.20E-05 3.20E-05 3.20E-07	1.20E+00 1.72E+00 2.44E+00 4.03E+00 5.82E+00 9.21E+00 1.14E+01 1.46E+01 2.17E+01 3.53E+01 5.20E+01 1.08E+02 2.48E+02 2.48E+02 3.51E+02 5.16E+02 5.16E+02 5.16E+02 9.28E+02	1.75E-02 2.33E-02 3.15E-02 2.24E-02 2.24E-02 1.26E-02 9.79E-03 6.69E-03 5.77E-03 5.77E-03 5.77E-03 5.77E-03 5.77E-03 5.52E-04 2.99E-04 2.99E-04 9.97E-05 1.76E-05 1.76E-05 1.58E-06	1.01E+00 1.42E+00 2.09E+00 5.33E+00 6.90E+00 9.61E+00 1.21E+01 1.58E+01 2.09E+01 2.09E+01 2.09E+01 3.92E+01 1.31E+02 1.31E+02 1.31E+02 4.09E+02 5.72E+02 7.56E+02	5.51E-02 5.24E-02 5.70E-02 2.68E-02 1.96E-02 1.52E-02 1.40E-02 1.40E-02 1.40E-02 1.40E-03 6.22E-03 6.22E-03 1.47E-03 1.47E-03 1.47E-03 1.47E-03 1.47E-03 1.47E-04 5.54E-04 5.09E-04 1.11E-04 2.20E-05 1.69E-06	7.87E-01 1.20E+00 1.79E+00 4.59E+00 4.59E+00 5.71E+00 9.09E+00 1.25E+01 1.65E+01 1.65E+01 5.10E+01 6.90E+01 1.11E+02 1.66E+02 3.46E+02 4.95E+02 6.25E+02	6.10E-02 8.70E-02 7.19E-02 3.06E-02 2.42E-02 2.25E-02 1.90E-02 1.90E-02 1.90E-02 1.90E-02 1.90E-03 3.34E-03 1.95E-03 1.24E-03 6.39E-04 3.57E-04 2.46E-04 1.51E-04 2.46E-05 2.08E-06	7.66E-01 1.09E+00 1.55E+00 2.36E+00 5.75E+00 4.76E+00 9.82E+00 1.34E+01 1.75E+01 1.75E+01 1.75E+01 3.02E+01 4.27E+01 1.05E+02 1.60E+02 2.13E+02 2.30E+02 4.13E+02 5.31E+02	6.50E-02 6.39E-02 7.12E-02 8.06E-02 2.45E-02 9.45E-02 9.45E-02 9.45E-03 1.97E-03 1.97E-03 7.02E-04 4.07E-04 2.26E-04 1.51E-04 2.56E-05 1.95E-06
Y=3.	5 IN	Y=3.	5 IN	Y=3.	5 IN	r=3.	5 IN	Y=3.	5 IN
Z= .	126 IN	Z= .	276 IN	Z= .	591 IN	Z=1.	119 IN	Z=2.	175 IN
K	E(K).700	K	E (K). UU	K	E (K) / W	K	E (K) /W	K	E(K)/W
(1.FT)	(FT)	(1.FT)	(FT)	(1.FT)	(FT)	(1/FT)	(FT)	(1/FT)	(FT)
2.14E+00 2.07E+00 4.69E+00 1.05E+01 1.54E+01 2.09E+01 2.44E+01 2.09E+01 2.49E+01 6.76E+01 9.05E+01 1.40E+02 1.75E+02 2.79E+02 4.79E+02 5.94E+02 8.75E+02 8.75E+02 5.75E+03	6.90E-03 8.022-03 7.632-03 6.01E-03 7.232-03 4.972-03 4.972-03 5.01E-03 5.01E-03 2.292-03 7.252-04 5.01E-03 7.252-04 9.242-05 5.402-05 5.402-05 4.972-02 4.972-03	1.24E+00 1.97E+00 2.99E+00 4.69E+00 6.76E+00 8.50E+00 1.95E+01 2.57E+01 3.04E+01 4.20E+01 4.20E+01 4.20E+01 5.42E+01 9.19E+02 1.95E+02 5.91E+	1.92E-02 2.93E-02 2.22E-02 1.66E-02 1.25E-02 1.14E-02 8.66E-03 5.69E-03 5.69E-03 5.59E-03 5.59E-03 5.35E-03 2.26E-03 1.26E-03 8.81E-04 5.00E-04 2.55E-05 1.45E-05 1.45E-05 1.45E-05	9.61E-01 1.59E+00 2.21E+00 3.36E+00 4.67E+00 6.12E+00 1.19E+01 1.59E+01 3.05E+01 3.05E+01 3.05E+01 7.85E+01 1.34E+02 2.02E+02 2.06E+02 3.96E+02 7.56E+02	5.21E-02 5.34E-02 4.61E-02 3.59E-02 2.62E-02 2.62E-02 1.74E-02 1.34E-02 9.05E-03 3.59E-03 2.61E-03 3.59E-03 1.57E-05 1.14E-05 6.12E-04 2.39E-04 1.22E-04 1.99E-06	8.60E-01 1.72E+00 1.72E+00 2.90E+00 5.09E+00 9.54E+00 9.54E+00 1.80E+01 2.40E+01 2.40E+01 3.49E+01 6.74E+01 1.11E+02 2.35E+02 3.46E+02 4.84E+02 6.22E+02	4.70E-02 4.50E-02 3.65E-02 2.52E-02 2.52E-02 2.52E-02 9.75E-05 6.55E-02 9.75E-05 4.35E-03 1.55E-03 1.56E-03 1.56E-03 1.56E-03 1.56E-04 3.35E-04 2.35E-04 1.25E-04 1.25E-04 1.90E-06	7.28E-01 1.19E+00 1.60E+00 2.27E+00 3.46E+00 4.39E+00 0.51E+00 1.66E+01 2.16E+01 2.26E+01 2.26E+01 4.60E+02 1.57E+02 2.34E+02 3.06E+02 5.06E+02 5.10E+02	8.45E-02 7.42E-02 6.30E-02 5.15E-02 4.39E-02 2.63E-02 1.25E-02 1.25E-02 1.25E-02 1.25E-02 1.25E-03 5.55E-04 3.27E-03 5.55E-04 3.27E-04 2.39E-04 1.63E-04 4.52E-05 6.31E-06

Table 5.5 Data on one-dimensional spectra at x = 10 ft, U<sub>∞</sub> = 30 ft/sec -- Continued

141 0	3.4	int v	See 4	101 1	3-4		3-1		3-4
5-42E-00 5-42E-00 2-22E-02 6-122E-02 2-22E-02 2-22E-02 2-22E-02 2-22E-02 1-22E-02 2-22E-05 2-22E	20+302*3 20+302*3 20+321*9 20+302*5 20+300*5 20+300*5 20+300 20	1:302-02 5:625-02 1:325-04 5:7415-04 5:7415-02 6:7455-02 6:7455-02 6:7455-02 6:055-05 6	e*536*05 3*46*05 5*326*05 3*46*05 1*190*1 3*326*05 3*326*05 3*326*05 1*526*05 1*526*00 1*526*05 1*526*00 2*326*05 1*526*00 2*326*00 2*326*00 2*326*00 2*326*00 2*326*00 2*326*00 2*326*00 2*326*00 2*326*00 2*36*000 2*36*00000000000000000000000000000000000	1' 326-02 1' 326-04 1' 026-04 2' 526-04 2' 526-04 2' 526-03 4' 026-02 4' 026-02 5' 026-02 1' 216-05 5' 026-05 5' 026-05	2, 4, 435, 402 2, 1, 26, 402 2, 5, 26, 402 3, 2, 26, 403 4, 1, 46, 405 4, 1, 46, 405 4, 1, 46, 405 4, 1, 46, 405 4, 1, 405	1,216-06 1,286-05 1,286-05 1,186-05 2,286-05 1,166-05 4,366-05 1,166-05 4,366-05 1,166-05 9,286-05 1,166-	20+352'5 22+352'5 22+325'5 23+325'5 25+325'7 25+325'7 25+325'7 10+325'1 10+328'7 10+328'7 10+320'7 10+320'7 10+320'7 10+325'1 10+325'7 10+352'7 100+3557 100+35	20-342 2112E-02 2112E-02 212E-02 212E-02 212E-02 212E-02 212E-02 212E-02 212E-02 2032E-02 212E-02 2	21 - 322 - 12 21 - 225 - 62 22 - 326 - 62 23 - 326 - 62 23 - 326 - 62 24 - 326 - 62 25 - 326 - 62 10 - 336 - 12 10 - 3
W/ 1313	(17FT) K		(13/1) K	(173) 3 (173)	(13/1) K	UV: (X) 3 (T3)	(11/11) K	E (K) / UU	(1.FT) K
NI SZI NI O	.Z=2 .Y=4.	NI 611 NI 0	'i=Z 'p=1	NI 169 NI 0		NI 922 NI 0		NI 9ZI NI 9ZI	:=2 . <sub>7=4</sub>

2:50E-02 1:12E-04 2:08E-04 2:22E-04 2:22E-04 2:22E-04 2:22E-02 2:08E-	2:20E+05 2:21E+05 2:21E+05 1:28E+05 1:28E+05 1:28E+05 2:14E+01 2:14E+01 2:14E+01 2:14E+01 1:28E+00 3:28E+00 3:28E+00 2:28E+00 3:28E+	1:366-00 3:15:6-02 1:46:-04 2:406-04 2:406-04 2:406-02 2:56:	e 248+05 e 248+05 e 258+05 e 258+05 e 258+05 e 258+05 e 21 288+05 e 21 288+05 e 21 288+01 e 288+0	11316-00 1308-15 1308-02 1512-00 1512-00 1512-00 1512-00 1512-00 1512-02 15	125240 1252400 125240 12524000 1252400000000000000000000000000000000000	4 36E-00 4 208E-02 1 11E-07 1 11E-07 1 208E-07 4 208E-07 4 208E-07 4 208E-02 5 368E-02 2 387E-02 8 47E-02 2 387E-02 1 288E-02 1 28	32548-05 325425 325425 325425 325425 325425 325425 325425 325426 325466 325566 325566 325566 325566 325566 325566 325566 325566 325566 3255666 3255666 3255666 3255666 32556666 32556666 32556666 32556666 325566666 32556666 32556666 325566666 325566666 325566666 325566666 325566666 3255666666 325566666 325566666666666666666666666666666666666	20-309'£ 90-368'P 50-369'P 50-369'P 50-322'i 10-350'P 10-309'8 20-351'2 20-351'2 20-351'2 20-351'2 20-396'P 20-396'P 20-390'2 20-390'2 20-390'1 20-395'1 20-325'1 20-325'1	1 - 392 - 1 20 - 320 - 2 20 - 369 - 1 20 - 369 - 7 10 - 350 - 5 10 - 350 - 5 10 - 300 - 2 10
20-377 9	10-300-2	20-380 9	10-317.8	20-398 7	10-369.0	20-312 1	00+367 1	20-369 1	00+129-1
(E1))	(1/L1) K	E (K1 / UU	(13/1) K	<u>100</u> /(30)3	(13/1) ×	UU. (N) 3	K (1.771)	E(K) /00	(11/1) ×
NI GLI'Z=Z NI O'G=X		NI 611'1=Z NI 0'5=A		NI 165' "Z NI 0'5"A		NI PLZ" =Z NI 0"S=A		NI 921' -2 NI 0'5-1	

5:16E-09 5:26-09 1*26:02 2:26:07 2:26:07 2:26:00 2:26:	2:56E*05 7:06E*05 7:06E*05 2:526*05 2:526*05 2:26*01 2:26*01 2:26*01 2:26*01 5:26*05 1:16*30 2:26*00 1:126*00 2:36E*00 2:36E*00 2:36E*00 1:25E*00 2:36E*00 2:36E*00 1:25E*00 2:36E*00 2	S0-372 6 S0-372 6 S0-301 7 S0-301 7 S0-301 7 S0-301 7 S0-302	2022 2120 2120 2120 2120 2120 2120 2120	1*20E-09 5*68E-07 1*26E-07 1*26E-07 5*88E-07 1*06E-02 5*26E-03 1*26E-03 1*26E-03 1*26E-03 1*26E-03 1*26E-03 5*26E-03 1*26E-03 5*26E-	1,252,402 2,362,402 2,362,402 2,362,402 2,362,402 2,362,402 1,222,400 2,362,400 1,222,400	1.016-02 1.026-02 1.026-02 1.026-02 1.026-03 0.025-035-03 0.025-035-03 0.025-035-035-035-035-035-035-035-035-	1'00E+02 2'28E+05 2'28E+05 2'28E+05 2'28E+05 2'28E+05 1'28E+05 1'28E+05 1'28E+05 2'708E+01 2'708E+01 1'28E+01 1'28E+01 1'28E+05 1'28E+05 2'08E+00 2'08E+05 1'28E+05 2'08E+05 1'28E+05 2'08E+05 1'28E+05 1'28E+05 2'08E+05 1'28	20-300°Z 1.726-00 20-391°T 1.908-02 2.176-07 2.176-07 2.176-02 2.176-07 2.176-02 2.176-	1'768-02 1'284-02 1'284-02 1'284-02 1'284-05 1'2
و(لا) ک <del>س</del> (۲٦)	لا (1751)		(13/1) K	E (K) VM	(1.4FT) K		(13/1) K	(FT) E(K) /W	(14/1) ×
NI 521'Z=Z NI 0'9=1		NI 611'1=Z NI 0'9=A		NI 165' =2 NI 0'9=1		NI 727 =2		NI 921 -2 NI 0'9-1	



Fig. 1.1 Schematic presentation of boundary layer flow along alternate, longitudinal strips of smooth and rough surfaces



Fig. 1.2 Coordinate system and flow features







Fig. 4.2 Plan view of flat plate










Fig. 4.5 False wind tunnel ceiling to obtain zero streamwise pressure gradient over flat plate



Fig. 4.6 Non-dimensional plot of static pressure distribution at the free stream along the centerline of the flat plate



Fig. 4.7 Non-dimensional plot of wall pressure distribution of the flat plate



Fig. 4.8 Mean velocity distribution at x = 5 ft. to check plane symmetric condition of the flow



Fig. 4.9 Typical comparison of mean velocity profiles measured by hot-wire and pitot tube at x = 6 ft, y = 0 in.







Fig. 4.11 Typical calibration curves of the single hot-wire

. .



Single Hot Wire



Crossed Wire in X-Z Plane



Crossed Wire in X-Y Plane





Fig. 5.1a Boundary layer growth over the flat plate along the centerline of a smooth strip







Fig. 5.2 Displacement thickness as a function of x



Fig. 5.3 Momentum thickness as a function of x



Fig. 5.4 Shape factor as a function of x

17.



Fig. 5.5 Spanwise variation of shear velocity



Fig. 5.6 Comparison of shear velocity for the smooth strip by various methods





Fig. 5.7b Longitudinal plots of power law profiles (y = 6 in.)-- U/U  $_{\!_{\infty}}$  vs. z/ $\!\delta$ 



Fig. 5.8 The parameter m in power law profile as a function of x



Fig. 5.9a Spanwise comparison of power law profiles at x = 6 ft.--  $U/U_{\infty}$  vs.  $z/\delta$ 



Fig. 5.9b Spanwise comparison of power law profiles at x = 10 ft.--  $U/U_{\infty}$  vs.  $z/\delta$ 



Fig. 5.10a Semi-logarithmic plots of mean velocity data at x = 6 ft.--  $U/U_{\infty}$  vs. z



Fig. 5.10b Semi-logarithmic plots of mean velocity data at x = 10 ft.--  $U/U_{\infty}$  vs. z



Fig. 5.11a Spanwise comparison of wall law profiles at x = 6 ft.--  $U/U_* vs. zU_*/v$ 



Fig. 5.11b Spanwise comparison of wall law profiles at x = 10 ft.--  $U/U_*$  vs.  $zU_*/v$ 



Fig. 5.12 Roughness effect  $\Delta U/U_{\star}$  vs.  $k_{r}U_{\star}/v$ 







Fig. 5.14a Spanwise comparison of velocity-defect profiles at x = 6 ft. --  $(U_{\infty}-U)/U_{\star}$  vs.  $z/\delta$ 

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Fig. 5.14b Spanwise comparison of velocity-defect profiles at x = 10 ft. --  $(U_{\infty}-U)/U_{*}$  vs.  $z/\delta$ 



Fig. 5.15a Spanwise comparison of velocity-defect profiles at x = 6 ft. --  $(U_{\infty}-U)/U_{*}$  vs.  $zU_{*}/\delta^{*}U_{\infty}$ 

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Fig. 5.15b Spanwise comparison of velocity-defect profiles at x = 10 ft. --  $(U_{\infty}-U)/U_{\star}$  vs.  $zU_{\star}/\delta^{\star}U_{\infty}$ 



Fig. 5.16a Spanwise comparison of velocity-defect profiles at x = 6 ft. --  $2(U_{\infty}-U)/(U_{*s}+U_{*})$  vs.  $z/\delta$ 



Fig. 5.16b Spanwise comparison of velocity-defect profiles at x = 10 ft. --  $2(U_{\infty}-U)/(U_{*s}+U_{*})$  vs.  $z/\delta$ 



Fig. 5.17a Spanwise comparison of velocity-defect profiles at x = 6 ft. --  $2(U_{\infty}-U)/(U_{*s}+U_{*})$  vs.  $zU_{*s}/\delta_{s}^{*}U_{\infty}$ 



Fig. 5.17b Spanwise comparison of velocity-defect profiles at x = 10 ft. --  $2(U_{m}-U)/(U_{*s}+U_{*})$  vs.  $zU_{*s}/\delta_{s}^{*}U_{\infty}$


Fig. 5.18a Spanwise velocity distribution of constant height at x = 6 ft. for  $U_{\infty}$  = 30 ft/sec



Fig. 5.18b Spanwise velocity distribution of constant height at x = 10 ft. for  $U_{\infty}$  = 30 ft/sec



Fig. 5.19a Curves of constant velocity at x = 6 ft. for  $U_{\infty}$ = 30 ft/sec



Fig. 5.19b Curves of constant velocity at x = 10 ft. for  $U_{\infty}$  = 30 ft/sec



Fig. 5.20a Distribution of turbulent intensities at x = 6 ft. and y = 0 in.

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Fig. 5.21a Distribution of Turbulent shear stresses at x = 6 ft. and y = 0 in.

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## Fig. 5.22e Spanwise comparison of -uw



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Fig. 5.23a Normalized spectra of streamwise turbulent velocity at x = 6 ft. and y = 0 in.



Fig. 5.23b Normalized spectra of streamwise turbulent velocity at x = 6 ft, and y = 6 in.



Fig. 5.24a Spanwise comparison of normalized spectra of streamwise turbulent velocity at x = 6 ft. and  $z/\delta \simeq 0.03$ 



Fig. 5.24b Spanwise comparison of normalized spectra of streamwise turbulent velocity at x = 6 ft. and  $z/\delta \approx 0.30$ 



Fig. 5.24c Spanwise comparison of normalized spectra of streamwise turbulent velocity at x = 6 ft. and  $z/\delta$   $\simeq$  0.54



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ABSTRACT The structure of the turbulent b alternate, longitudinal strips of smooth an allel to the direction of the flow has been velocity, wall shear stress, turbulent inte energy spectra of the streamwise turbulent The flow field can be subdivided into region along the centerline of each smooth centerline of each rough strip; and (c) the	oundary laye d rough (San investigate ensities, tur velocity wer three region strip; (b) t "intermedia	er on a fla ad paper, ( ed. Measur- bulent sho be obtained as. These the "rough" ate" region	at plate consisting of Grit 4) surfaces par- rements of mean ear stresses and d. are: (a) the "smooth" " region along the n lying between the				
smooth and the rough regions. In both smoo	oth and rough	regions,	the flow conditions				
are found to be nearly analogous to those 1	n two-dimens	ate region	ndary layers over				
smooth and rough walls, respectively. In t	dly to the 1	ace region	itions of the wall				
while going from the smooth to the rough re	gion. The v	vertical d	istribution of mean				
velocity can satisfactorily be expressed by	the two-dim	nensional	descriptions for the				
law of the wall and the velocity-defect law. However, the two laws are not universal							
out form a family of curves depending on the local wall shear stress. The spanwise							

change from smooth-wall conditions to rough-wall conditions will induce a weak crossflow directed from the rough to the smooth region near the wall. A theoretical description of the turbulent shear stress  $-\rho v w$  which is responsible for the generation of this cross-flow has been presented. The spanwise adjustment of turbulence quantities to local conditions of the wall takes place at a slower rate than does the wall shear stress. The normalized spectra of the streamwise turbulent velocity remains largely unaffected by the surface configuration investigated here.

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